

Towards electroweak reactions with quantum computers

Alessandro Roggero
with: R. Gupta & J. Carlson (LANL)
A. Li & G. Perdue (FNAL)

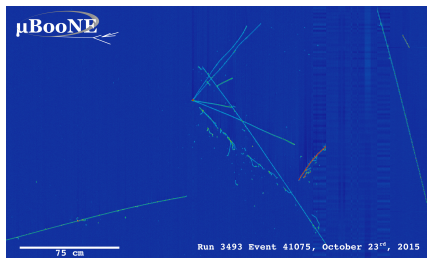


figure credit: μ BooNE collab.

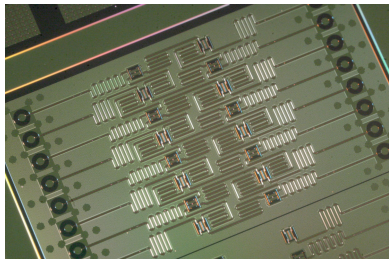


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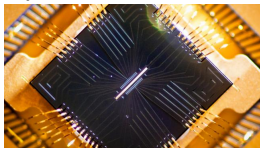


Santa Fe LQCD Meeting
Santa Fe – 30 August, 2019



What is a Quantum computer?

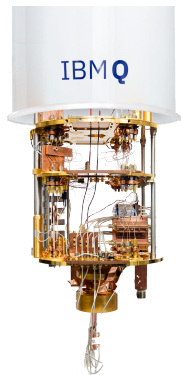
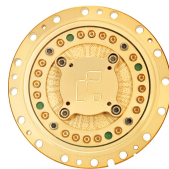
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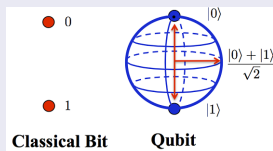


Rigetti



● Microsoft?

Bits vs Qubits



- N bits: an integer number $< 2^N$
- N qubits: a vector $|\psi\rangle$ in 2^N -dim Hilbert-space
 \implies exponentially more information available

Quantum Simulations with qubits

“Nature isn’t classical, dammit, and if you want to make a simulation of nature, you’d better make it quantum mechanical.”

— R.Feynman (1982)

- in 1996 S.Lloyd shows this intuition is correct for local interactions

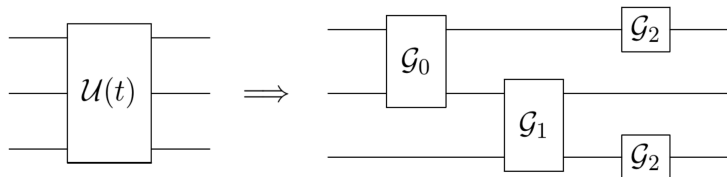
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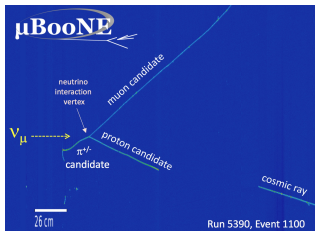
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- in 1996 S.Lloyd shows this intuition is correct for local interactions
- choose a finite basis to discretize system $\rightarrow \dim(\mathcal{H}) = \Omega \propto e^A$
- physical states can be mapped in states of $\sim \log_2(\Omega)$ qubits

$$|\Psi(t)\rangle = U(t) |\Psi(0)\rangle$$



Exclusive cross sections in neutrino oscillation experiments



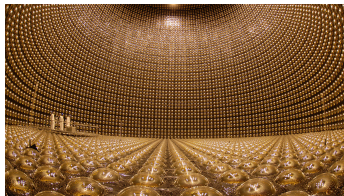
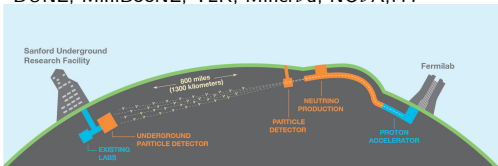
Goals for ν oscillation exp.

- neutrino masses
- accurate mixing angles
- CP violating phase

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2 L}{4E_\nu}\right)$$

- need to use measured reaction products to constrain E_ν of the event

DUNE, MiniBooNE, T2K, Minerva, NO ν A, ...

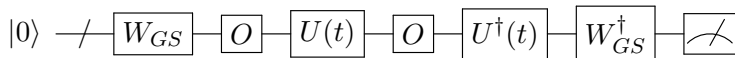


Quantum algorithms for the nuclear response

$$R(\omega) = \int dt e^{i\omega t} C(t) \quad \text{with} \quad C(t) = \langle \Psi_0 | O(t) O(0) | \Psi_0 \rangle$$

Blueprint of quantum algorithms

state preparation \rightarrow unitary evolution \rightarrow measurement

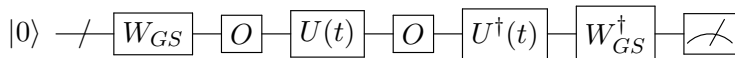


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- strategy A [Ortiz, Somma et al (2001-2003)]
 - compute $C(t)$ on quantum computer for different times
 - perform Fourier transform classically
- strategy B [Roggero & Carlson (2018)]
 - sample directly final states from approximate response function

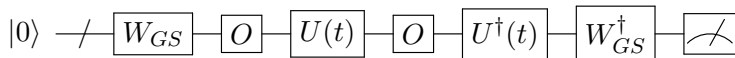
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- both algorithms are *poly* in A and target energy resolution !

Part I: baby steps on current machines

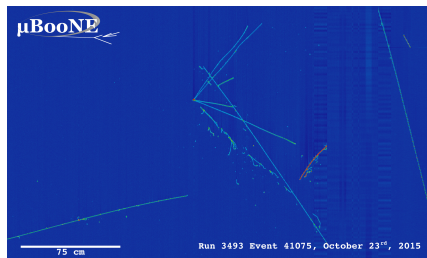


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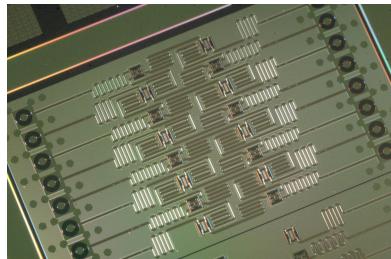


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credit: Atari Inc.

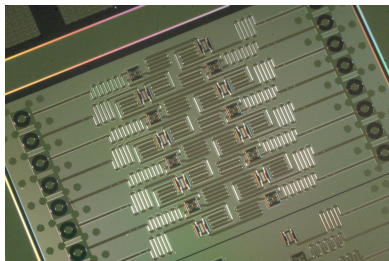


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A (very) simplified nuclear model

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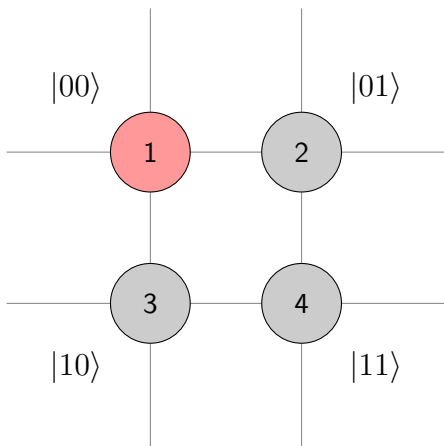
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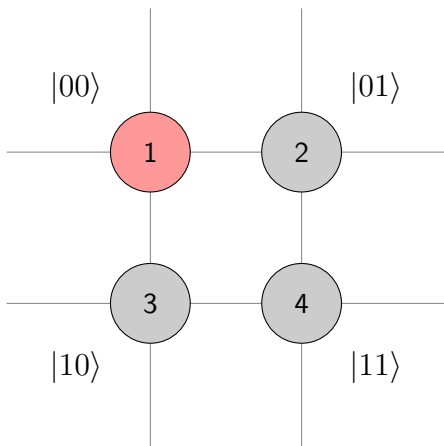
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- for every nucleon map 4 states into 2 qubits \Rightarrow 4 qubits total

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- interaction Hamiltonian from pion-less EFT becomes

$$H = -t \sum_k X_k + U_1 \sum_k Z_k + \sum_{i,j} U_{ij} Z_i Z_j + \sum_{i,j,k} V_{ijk} Z_i Z_j Z_k + \sum_{i,j,k,l} W_{ijkl} Z_i Z_j Z_k Z_l$$

Preparation of an approximate ground state

Variational Quantum Eigensolver

Perruzzo(2014), McClean(2015), ...

Use Rayleigh-Ritz variational principle to find the lowest energy state

$$\min \langle \phi(\vec{p}) | H | \phi(\vec{p}) \rangle \quad \text{with} \quad |\phi(\vec{p})\rangle = U(\vec{p}) |0\rangle$$

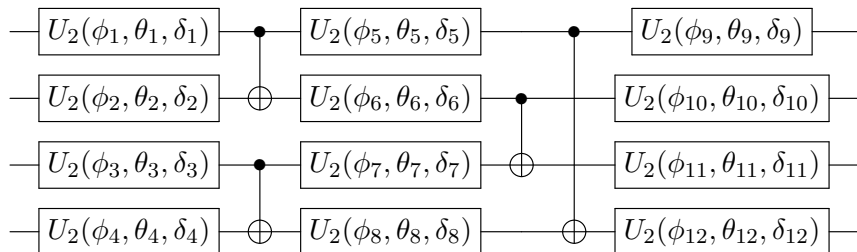
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- up to 36 parameters \rightarrow could be reduced using symmetries

Preparation of an approximate ground state II

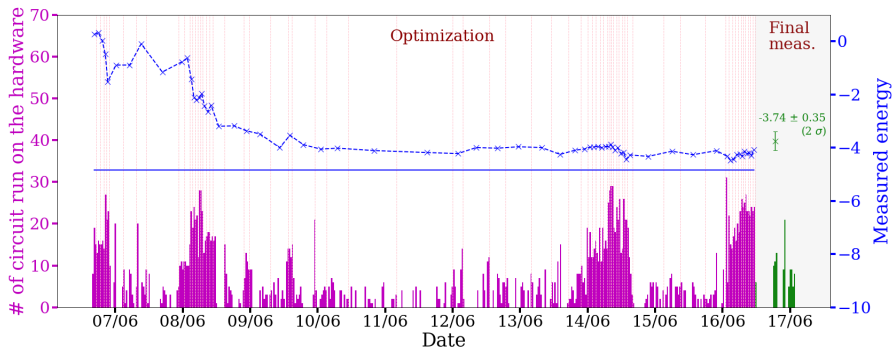
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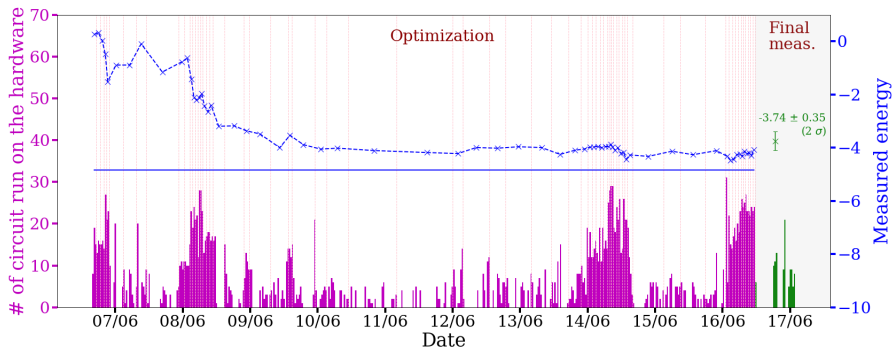
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- first problem: latency \rightarrow more compact trial states, better optimizers



Preparation of an approximate ground state III

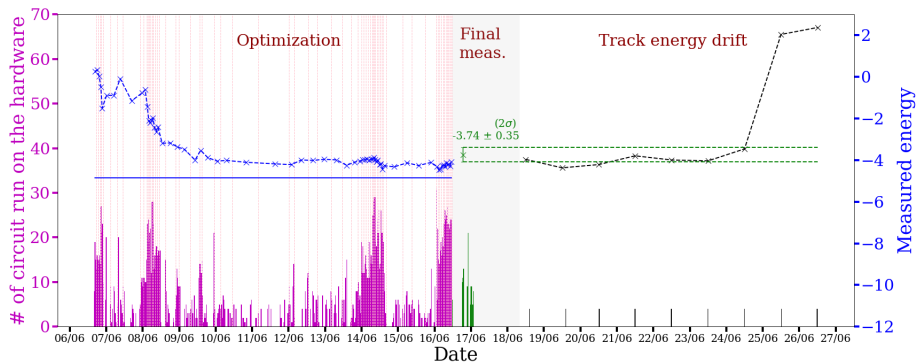
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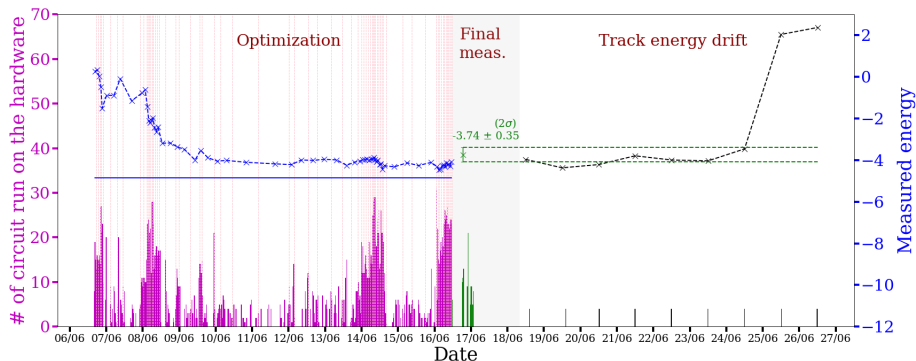
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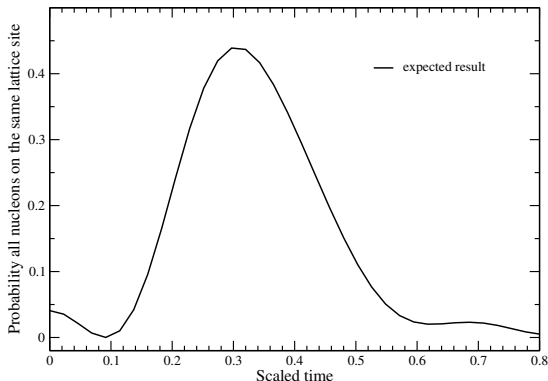
$$\min \langle \phi(\vec{p}) | H | \phi(\vec{p}) \rangle \quad \text{with} \quad |\phi(\vec{p})\rangle = U(\vec{p}) |0\rangle$$

- second problem: persistence \rightarrow track changes and reoptimize



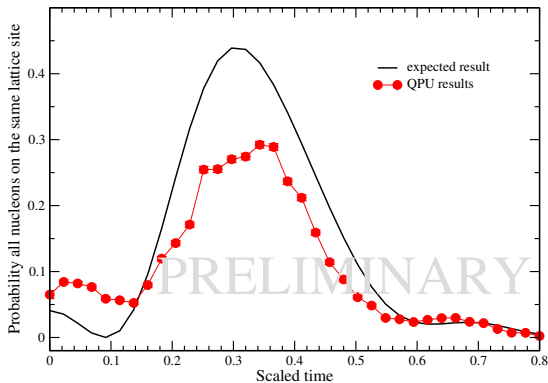
Preliminary results for real time dynamics

For now just $\langle \Psi | O(t) | \Psi \rangle \rightarrow$ move to 2pt functions in the future



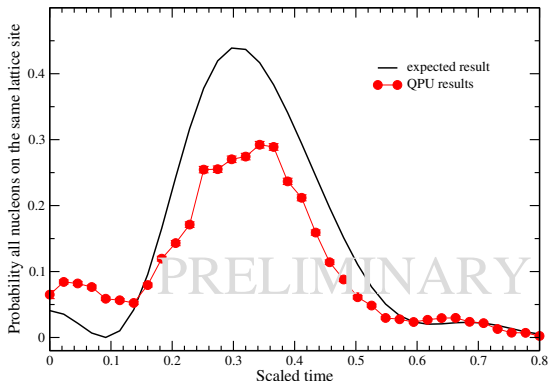
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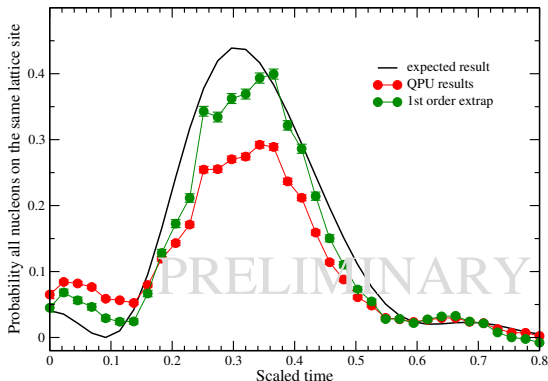


PROBLEM: large systematic errors from machine noise

for perturbative noise: $R(\epsilon) = R_0 + \epsilon R_1 + \epsilon^2 R_2 + \dots$

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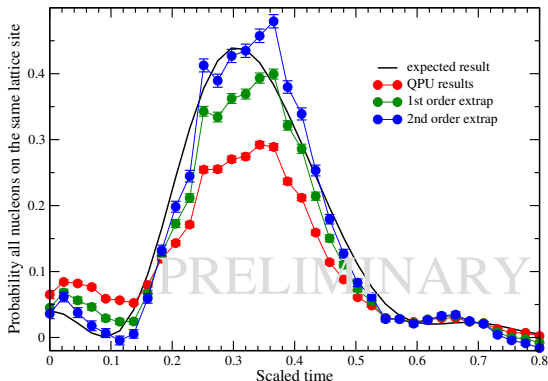


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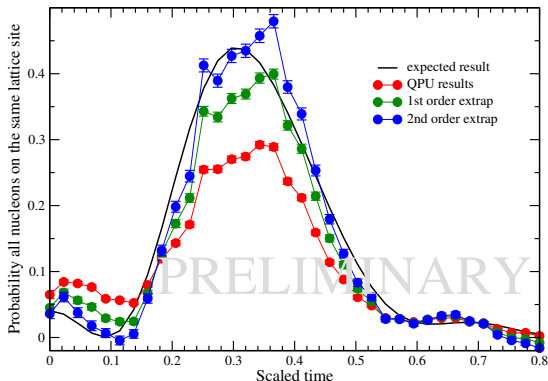


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- currently we are working hard on different mitigation techniques

Part II: back to neutrino scattering off ^{40}Ar

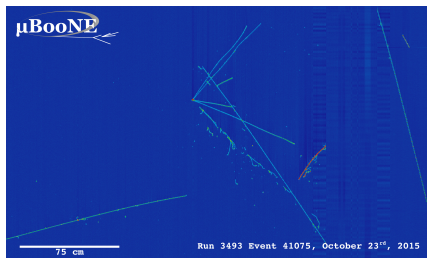


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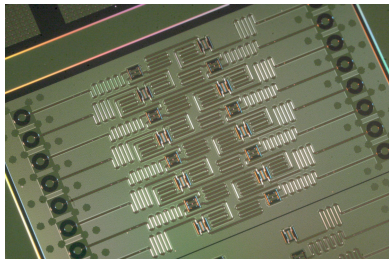
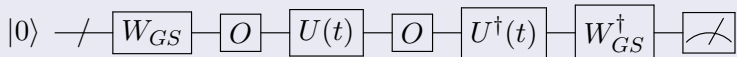


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Blueprint of quantum algorithms from beginning of talk



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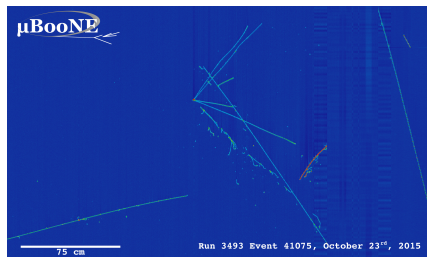


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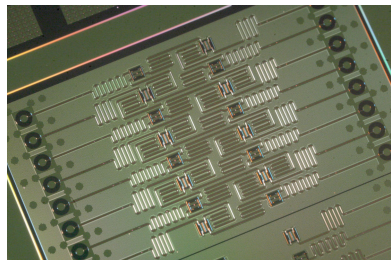
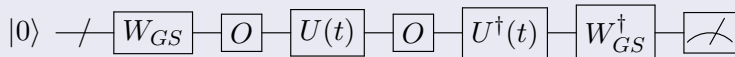


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- we can use variational ansatz to prepare initial state W_{GS}

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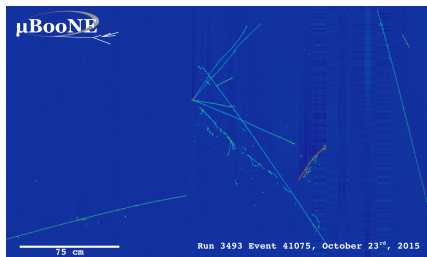


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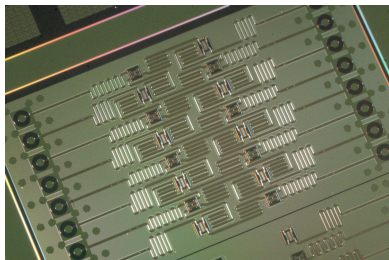
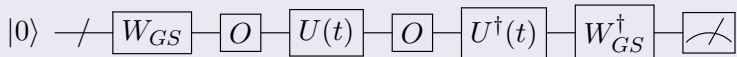


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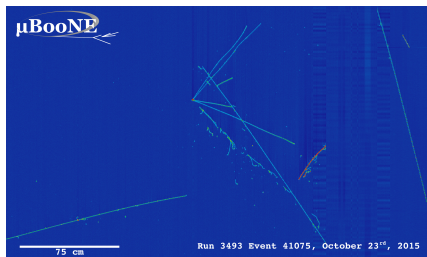


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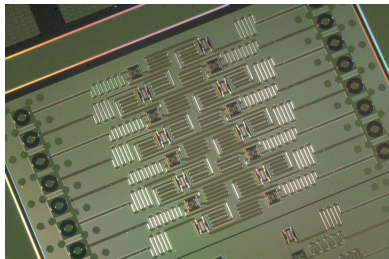
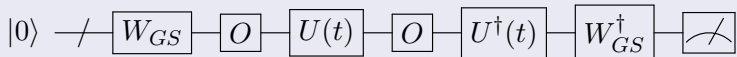


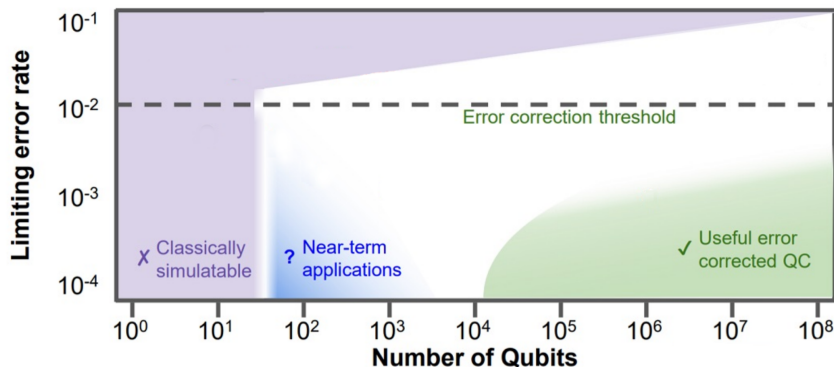
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- we can use variational ansatz to prepare initial state W_{GS}
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- need to repeat for many values of momentum transfer $O \equiv O(q)$

Need Both Quality and Quantity



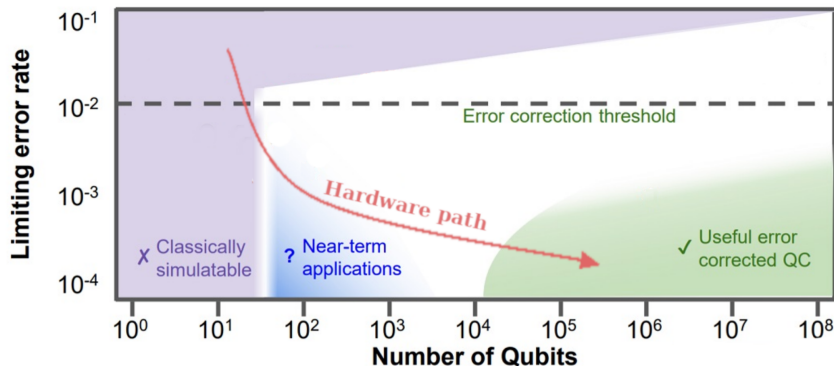
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Threshold Theorem(s)

Ben-Or, Aharonov, Kitaev, Knill, Gottesman, . . .

When rate below threshold can extend $\tau_{coh}^{eff} \rightarrow \infty$ with $polylog(N)$ effort

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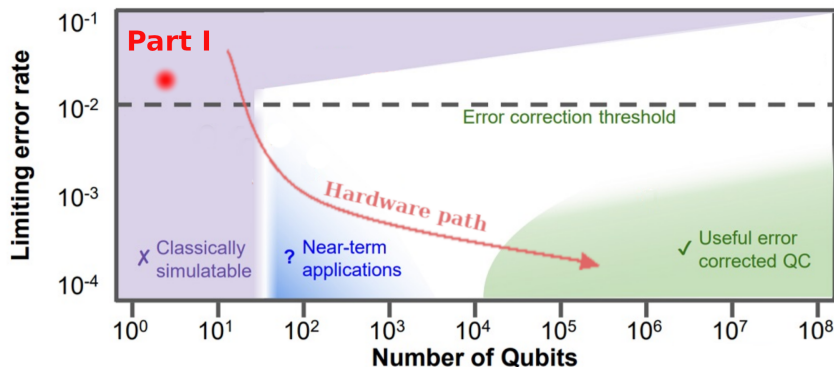
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- pionless EFT on a 10^3 lattice of size 20 fm [$a = 2.0$ fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want $R(q, \omega)$ with 20 MeV energy resolution

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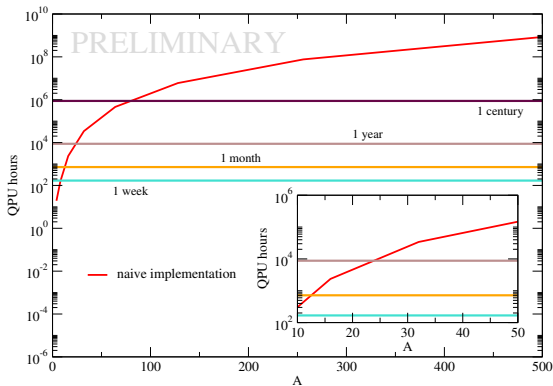
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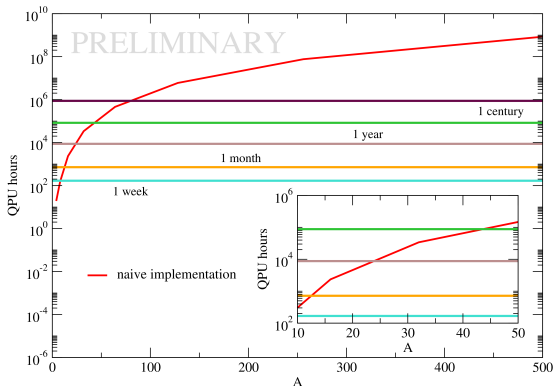
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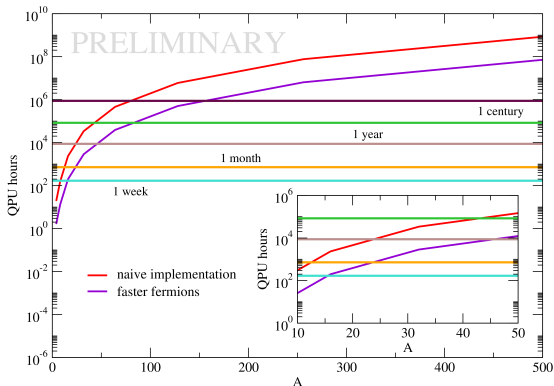
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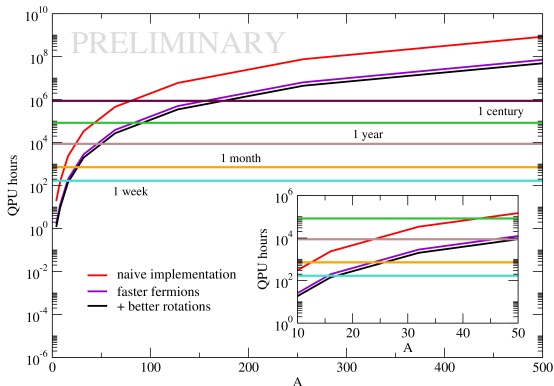
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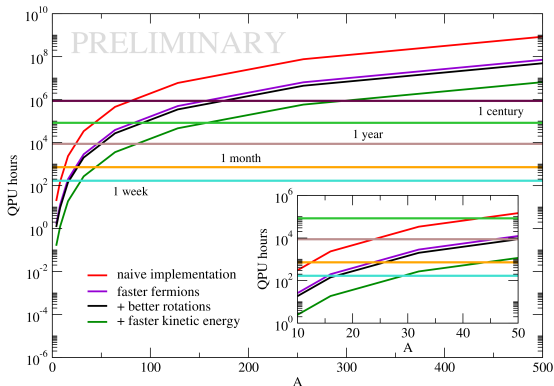
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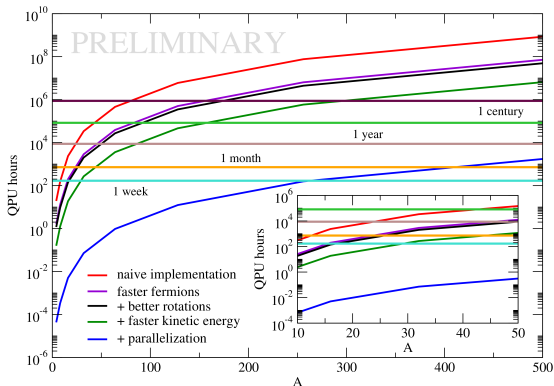
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- pionless EFT on a 10^3 lattice of size 20 fm [$a = 2.0$ fm]
- 10x faster gates and negligible error correction cost (very optimistic)
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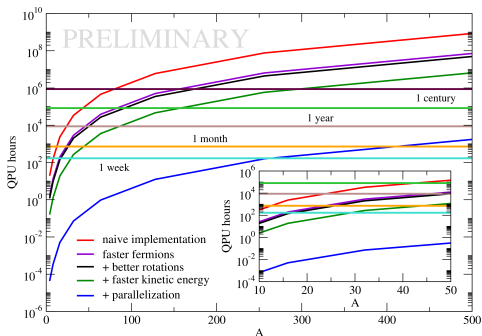
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coherence time for ^{40}Ar

naive ≈ 9 years

optimized ≈ 3 minutes

- algorithm efficiency is critical
- there is still a long way to go
- find new algorithms and/or approximations for near term

Summary

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- QC is an emerging technology with the potential of revolutionize the way theory calculations are done
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-
- J. Carlson & R. Gupta (LANL)
 - Andy Li & G. Perdue (FNAL)
 - QPU access thanks to ORNL



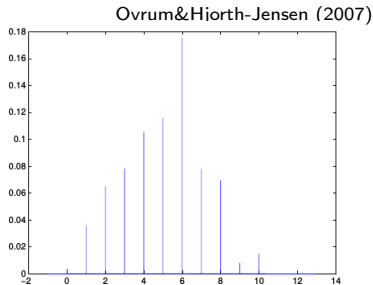
Quantum Phase Estimation

Kitaev (1996), Brassard et al. (2002), Svore et. al (2013), Weibe & Granade (2016),...

QPE is a general algorithm to estimate eigenvalues of a unitary operator

$$U|\xi_k\rangle = \lambda_k|\xi_k\rangle, \lambda_k = e^{2\pi i\phi_k} \iff U = e^{-itH}$$

- starting vector $|\psi\rangle = \sum_k c_k|\xi_k\rangle$
- store time evolution $|\psi(t)\rangle$ in auxiliary register of M qubits
- perform (Quantum) Fourier transform on the auxiliary register
- measures will return λ_n with probability $P(\lambda_n) \approx |c_n|^2$



BONUS: final state after measurement is $|\psi_{fin}\rangle \approx \sum_k \delta(\lambda_k - \lambda_n)c_k|\xi_k\rangle$