

WILL DETMOLD

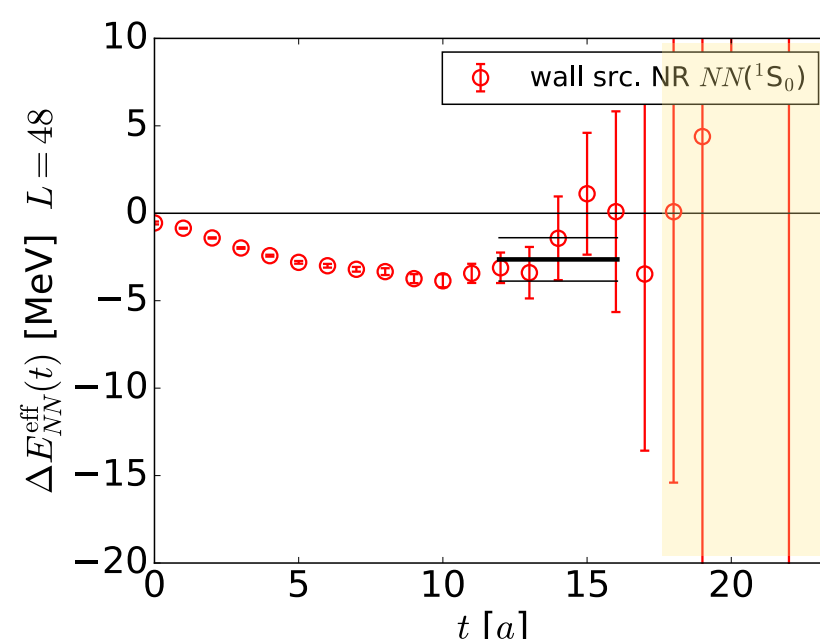
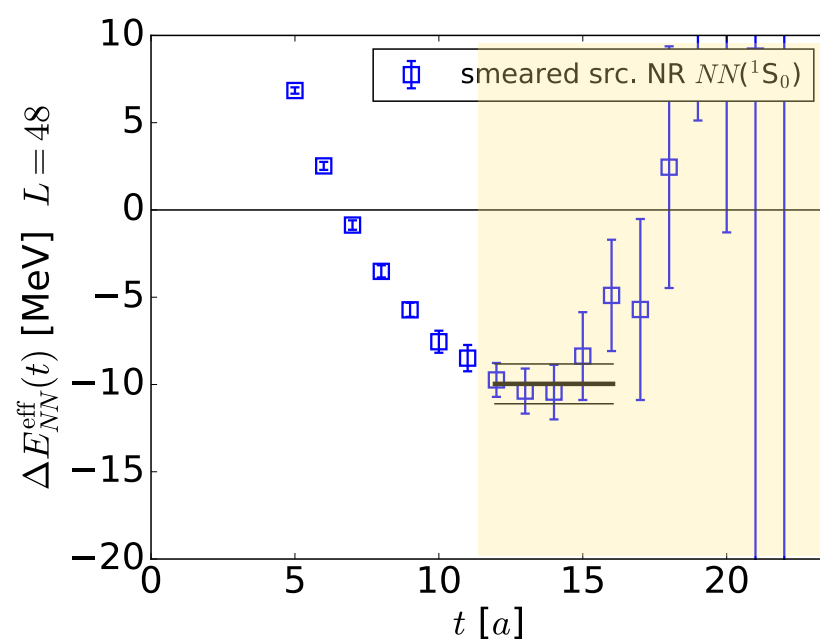
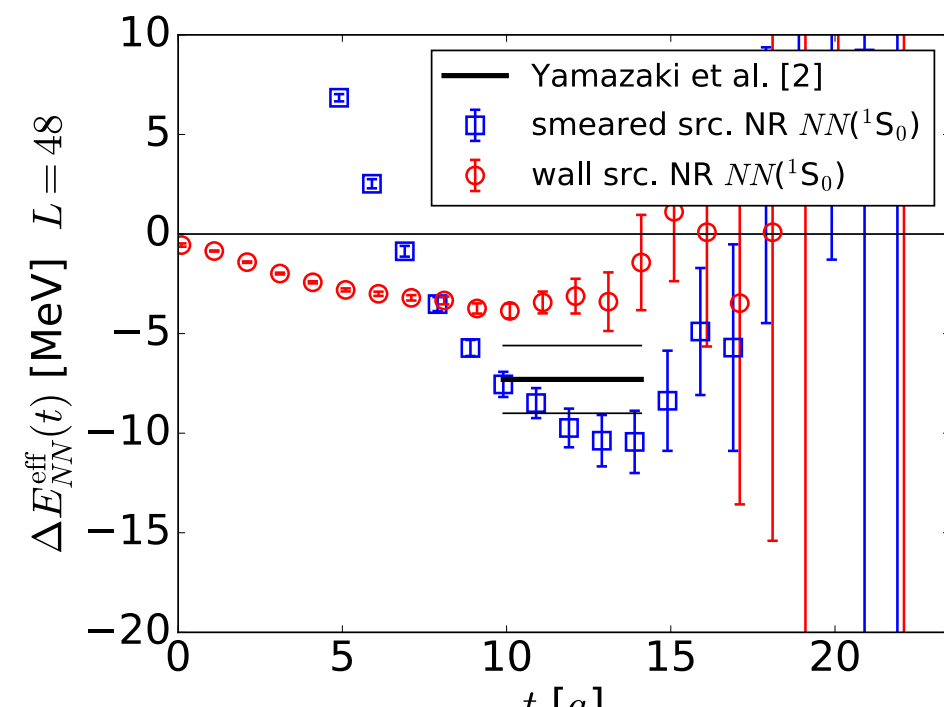
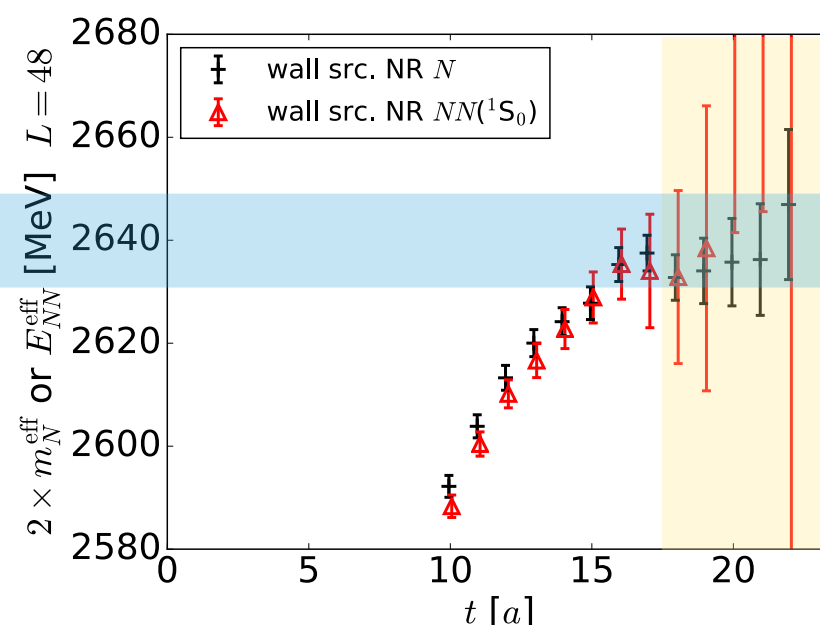
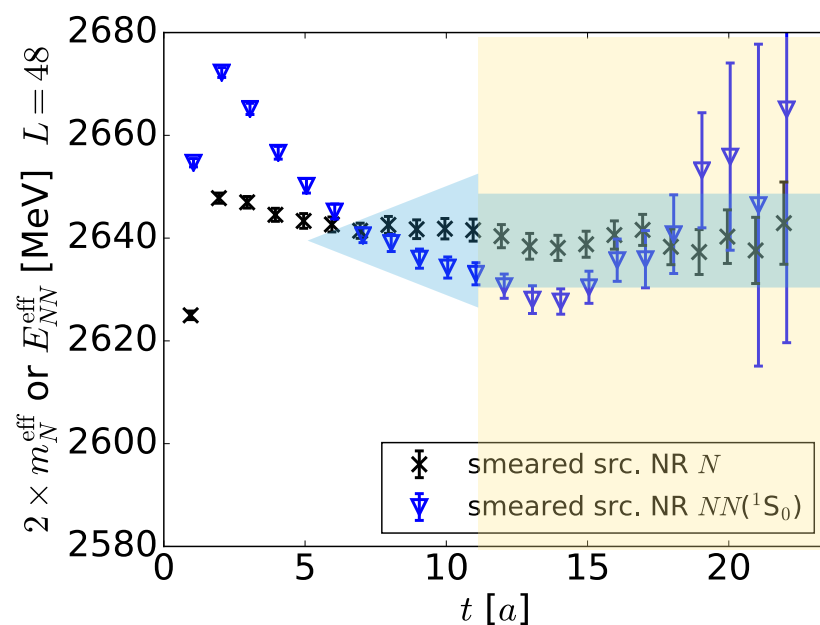
MIT

PROPERTIES OF LIGHT NUCLEI FROM LATTICE QCD

HAL2016

from Iritani et al [HALQCD], JHEP1610(2016)101

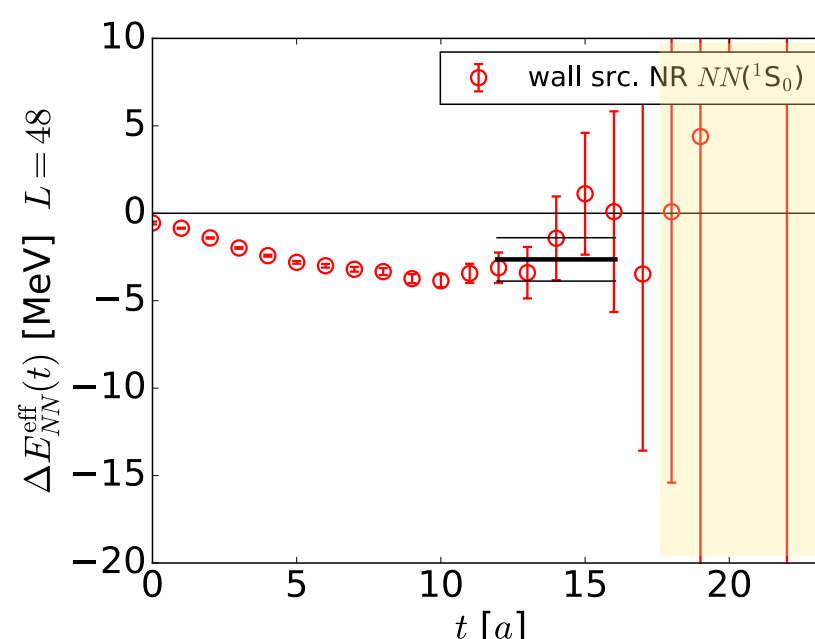
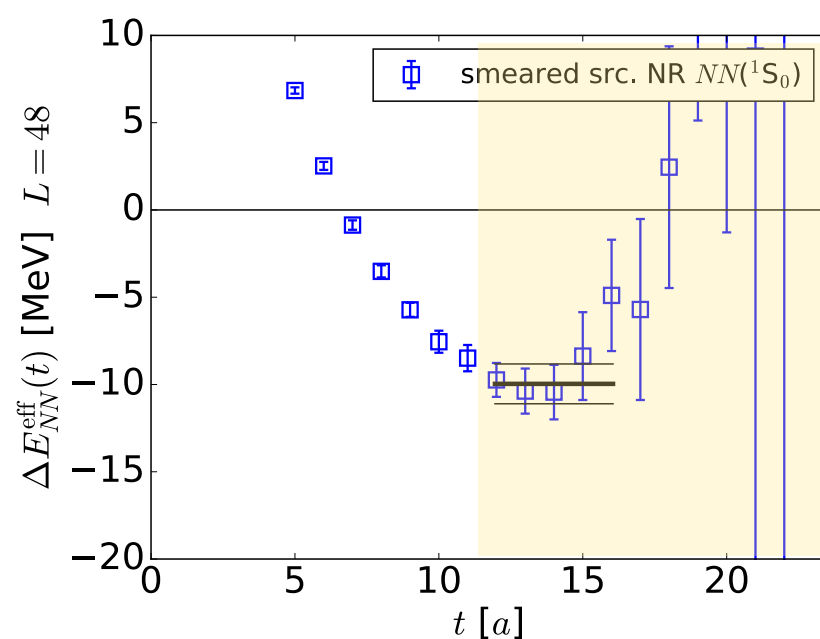
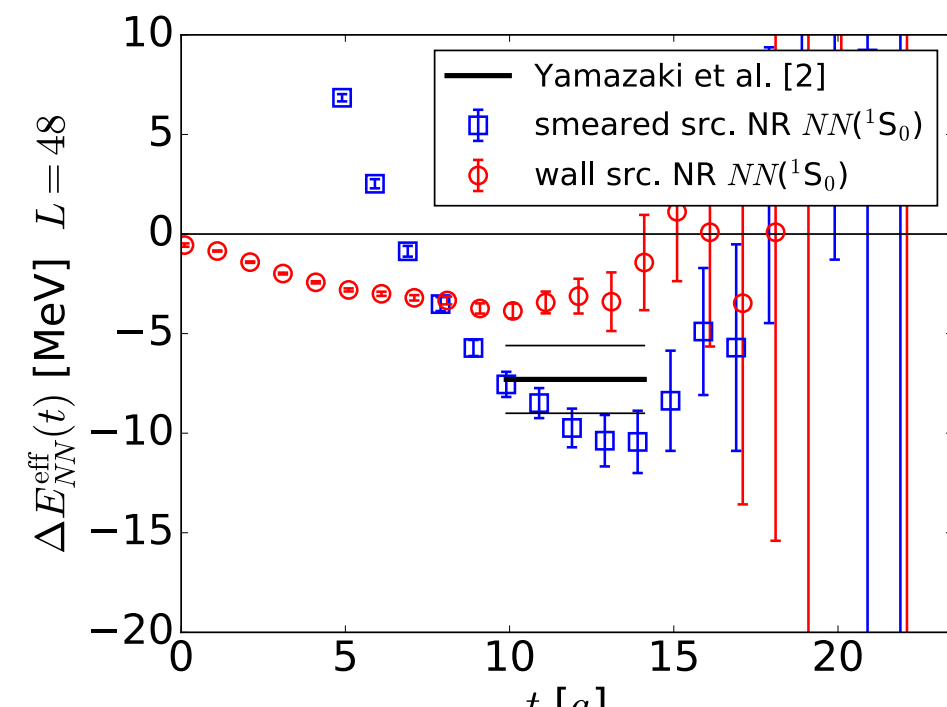
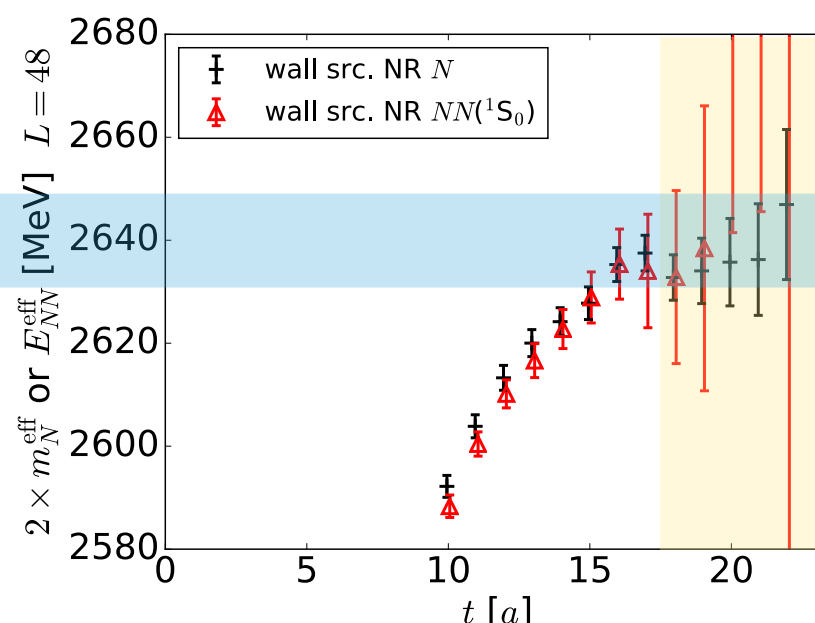
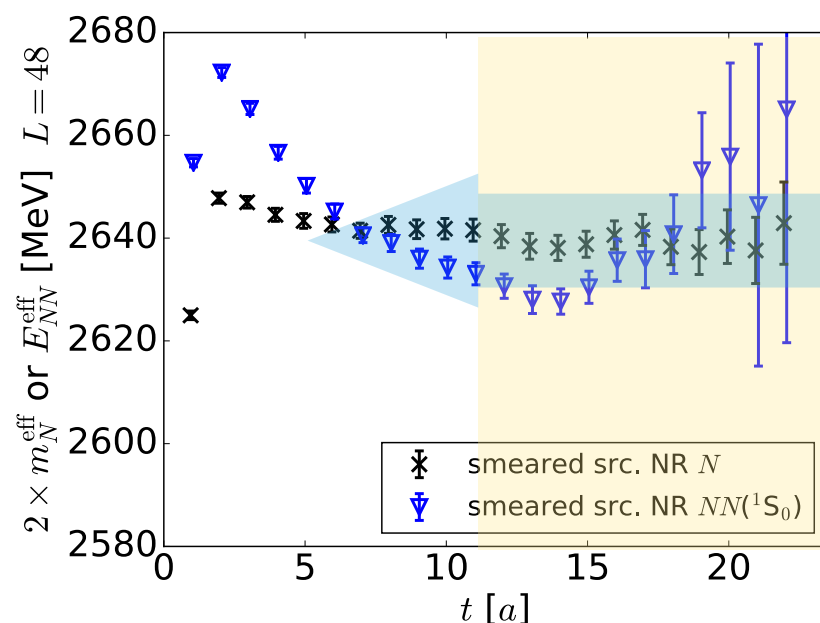
$$R(t) = \frac{C_{NN}(t)}{(C_N(t))^2}$$



HAL2016

from Iritani et al [HALQCD], JHEP1610(2016)101

$$R(t) = \frac{C_{NN}(t)}{(C_N(t))^2}$$

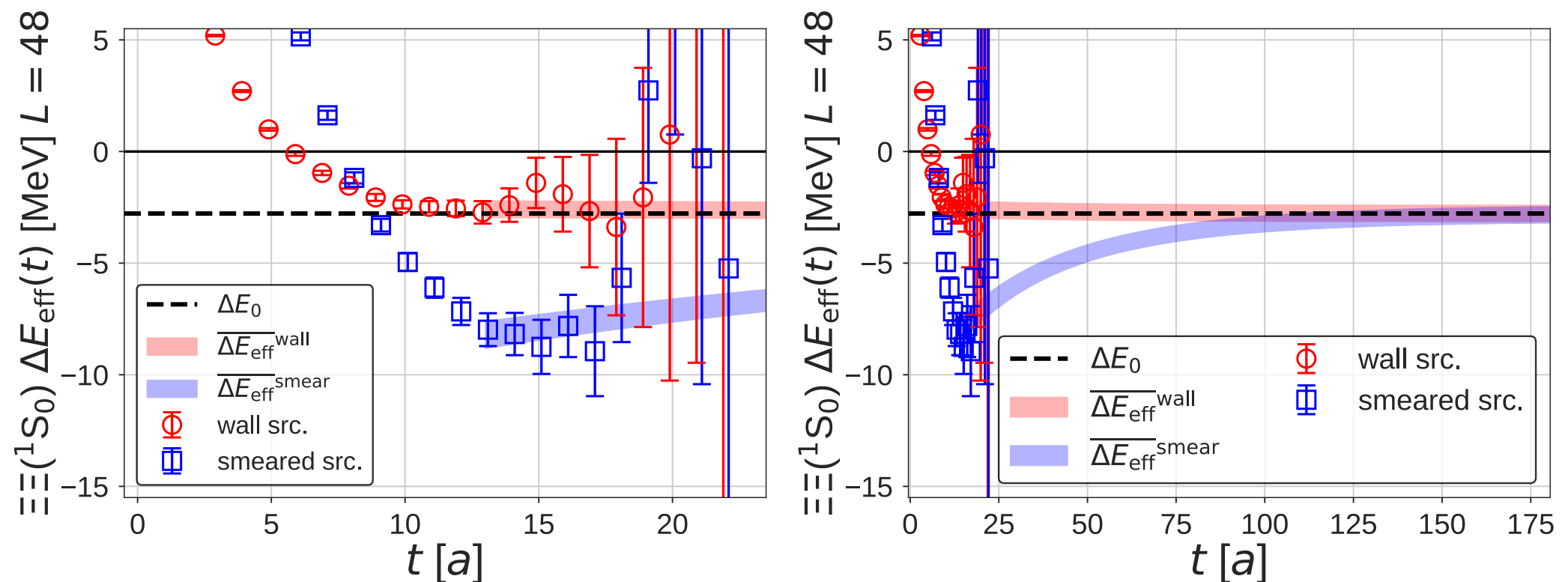


When single hadrons are in ground state there is consistency between extracted energy or energy shifts

HAL2016

from Iritani et al [HALQCD], JHEP1610(2016)101

Wall and Exp sources give different ΔE

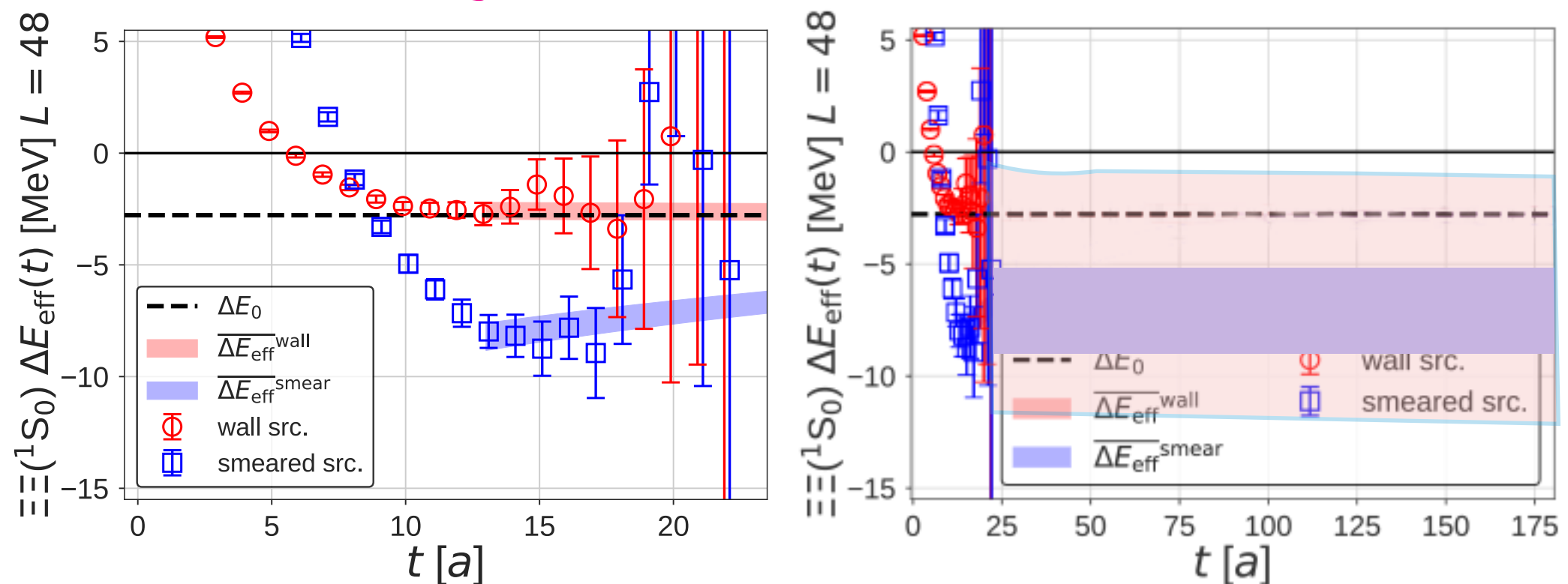


NB: this fig now for $\Xi\Xi$ rather than NN, but same behaviour seen

HAL2016

from Iritani et al [HALQCD], JHEP1610(2016)101

Wall and Exp sources give different ΔE

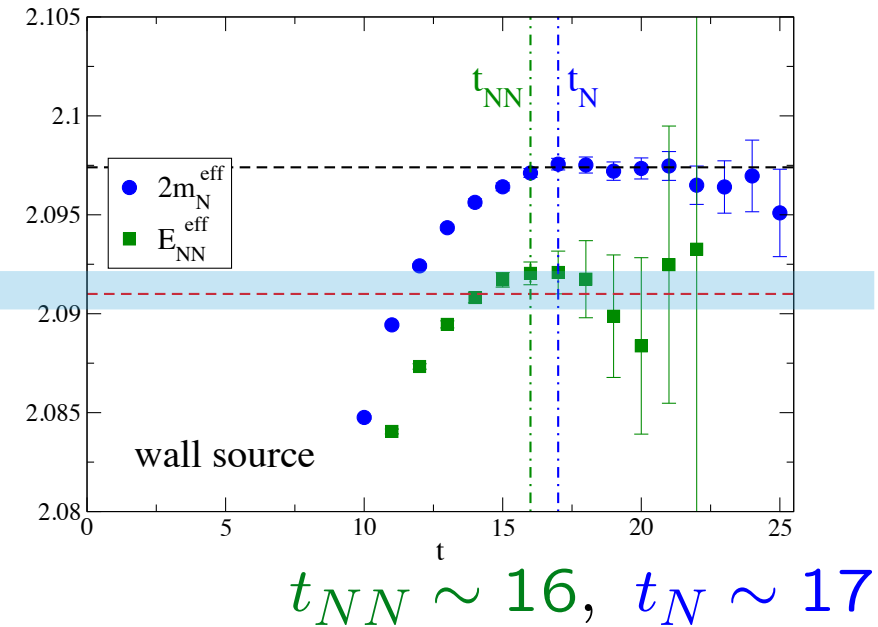
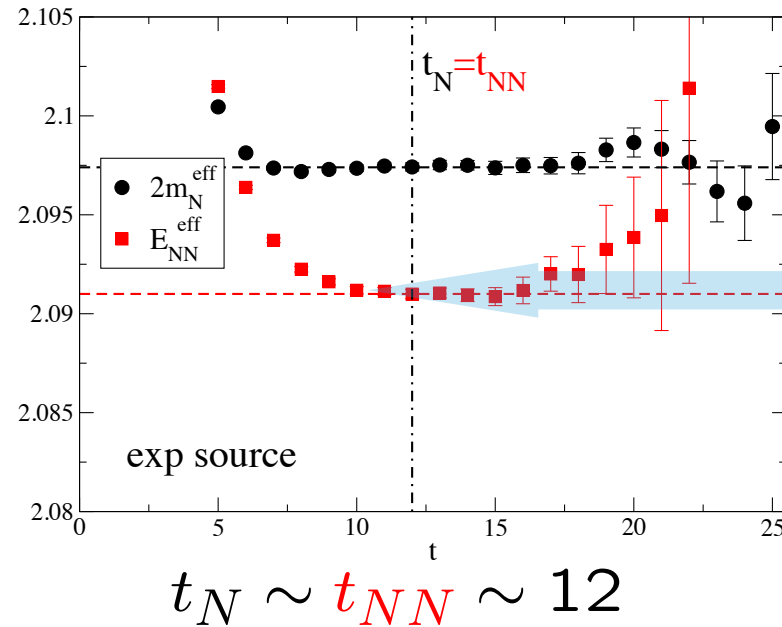


For more details, see extensive high statistics study by PACS-CS
T. Yamazaki et al [PACS collaboration], LATTICE 2017 [arXiv:1710.08066](https://arxiv.org/abs/1710.08066)
updates in YITP workshop in 2019
<http://www2.yukawa.kyoto-u.ac.jp/~flqcd2019/slides/Yamazaki.pdf>

NB: this fig now for $\Xi\Xi$ rather than NN, but same behaviour seen

PACS 2017 HIGH STATISTICS

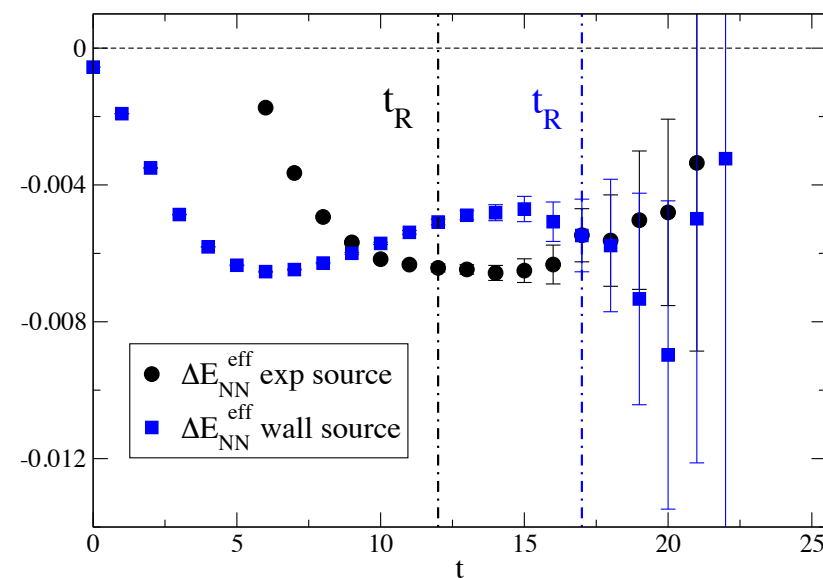
Effective $2m_N$ and E_{NN}



Completely
consistent
in plateau
regions

| L | T | source | N_{meas} |
|-----|-----|--------|-------------------|
| 16 | 64 | Exp | 15,544,000 |
| | | Wall | 8,307,200 |
| 20 | 64 | Exp | 5,504,000 |
| | | Wall | 4,480,000 |
| 32 | 64 | Exp | 10,496,000 |
| | | Wall | 8,307,200 |

Amazing statistics!



Completely
consistent
in plateau
regions once
single hadron
in ground
state

NB: have exp & wall rather than smeared & wall, quenched cfs

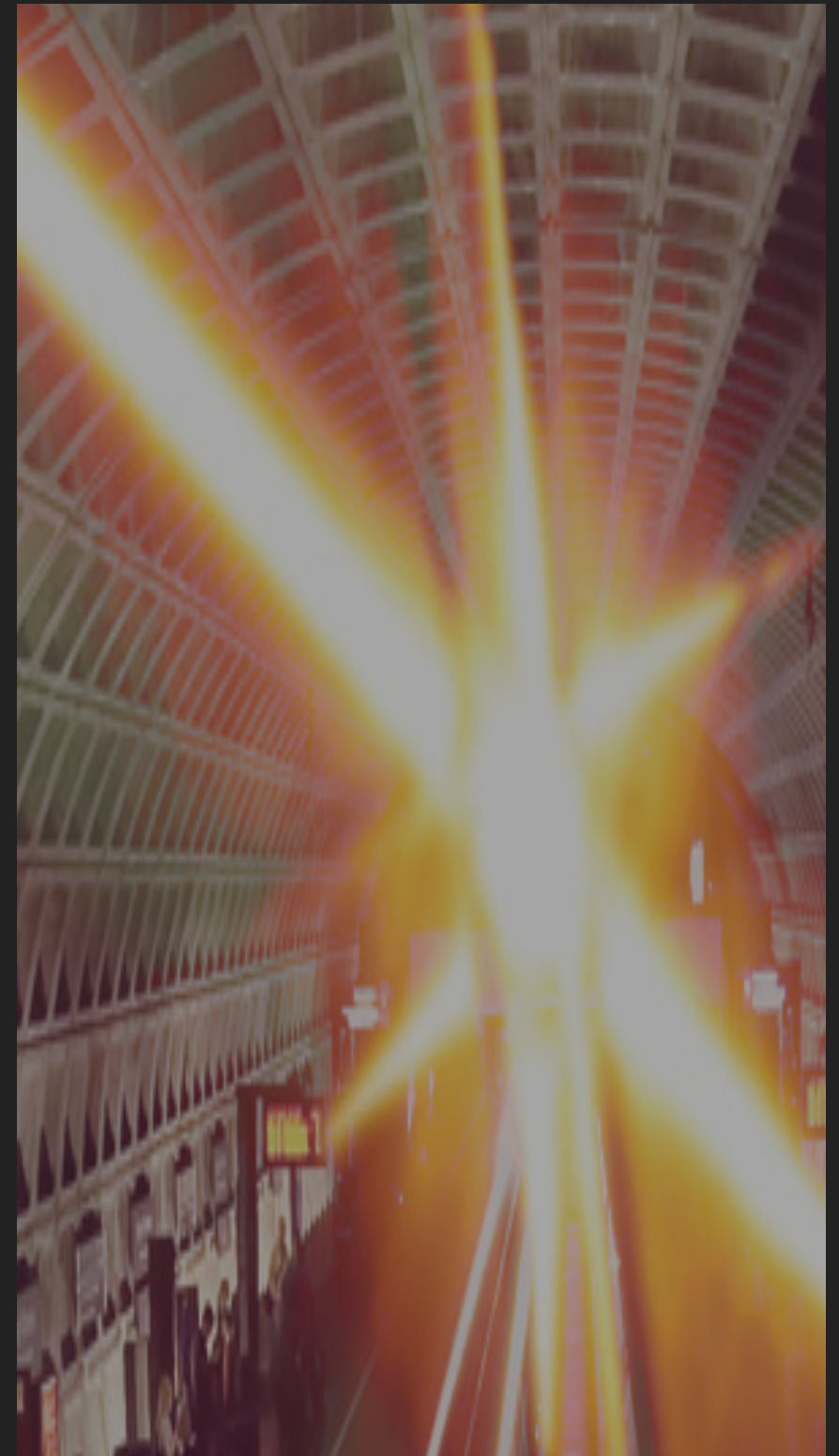
WILL DETMOLD

MIT

PROPERTIES OF LIGHT NUCLEI FROM LATTICE QCD

THE INTENSITY FRONTIER

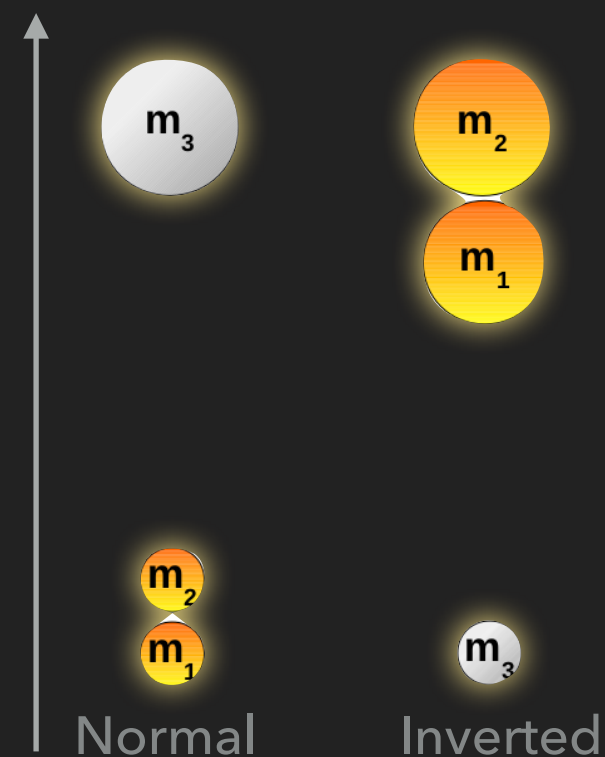
- ▶ Seek new physics through quantum effects
- ▶ Precise experiments
 - ▶ Sensitivity to probe the rarest interactions of the SM
 - ▶ Look for effects where there is no SM contribution
- ▶ Major component is nuclear targets
- ▶ Important focus of HEP/NP experimental program
 - ▶ Neutrino physics
 - ▶ Dark matter direct detection
 - ▶ Charged lepton flavour violation, EDMs, $\beta\beta$ -decay,



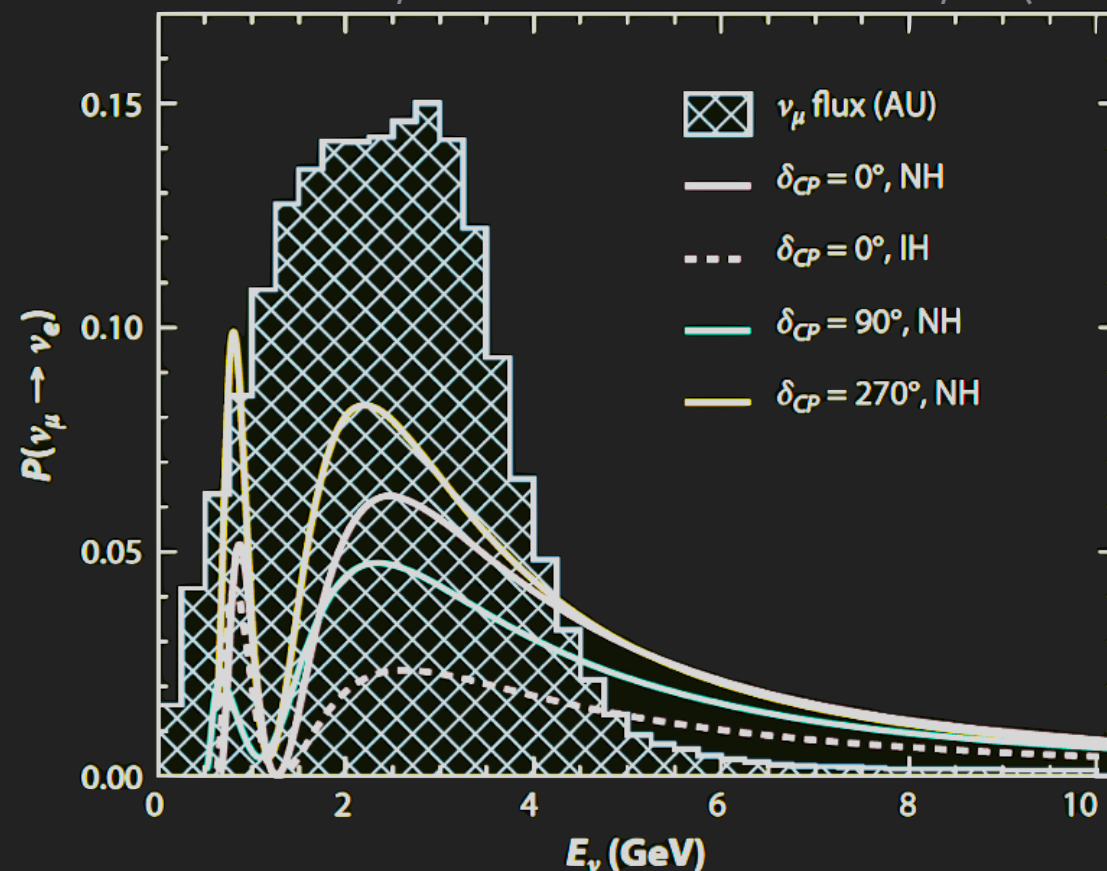
LONG BASELINE NEUTRINO EXPERIMENTS

- ▶ Deep Underground Neutrino Experiment
 - ▶ Flagship facility for US HEP for next decades
 - ▶ Determine neutrino mass hierarchy and extract mixing parameters
- ▶ Neutrino scattering on argon target
 - ▶ Need fluxes/energies to high accuracy
 - ▶ Need to know interactions with argon over a wide range of energies

Neutrino Mass²

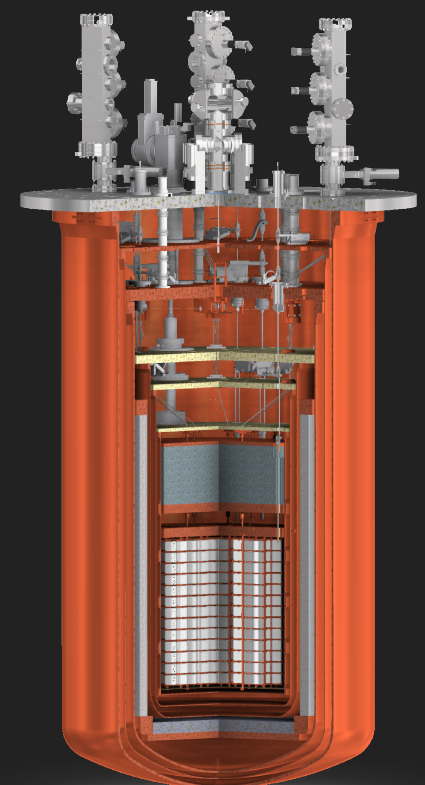
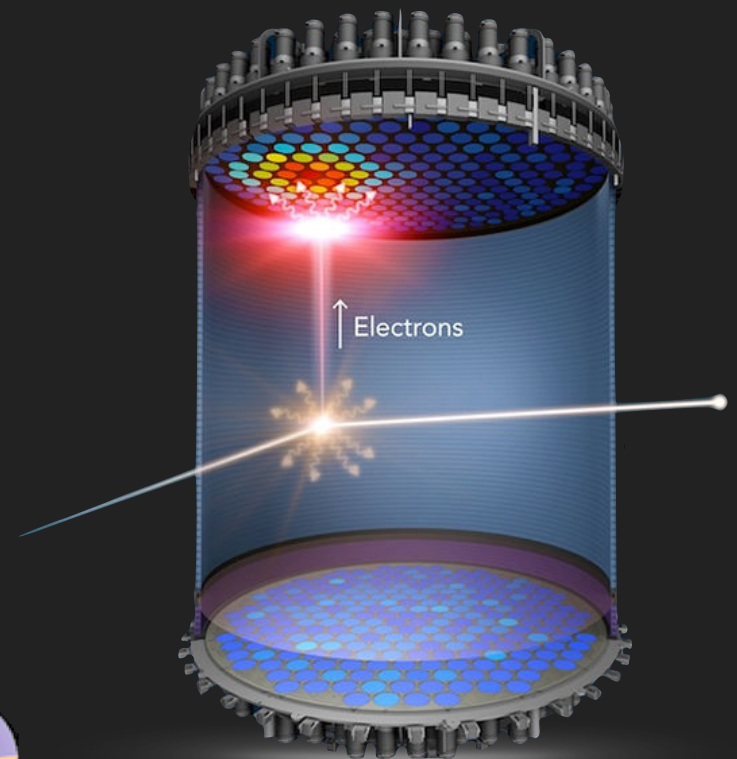
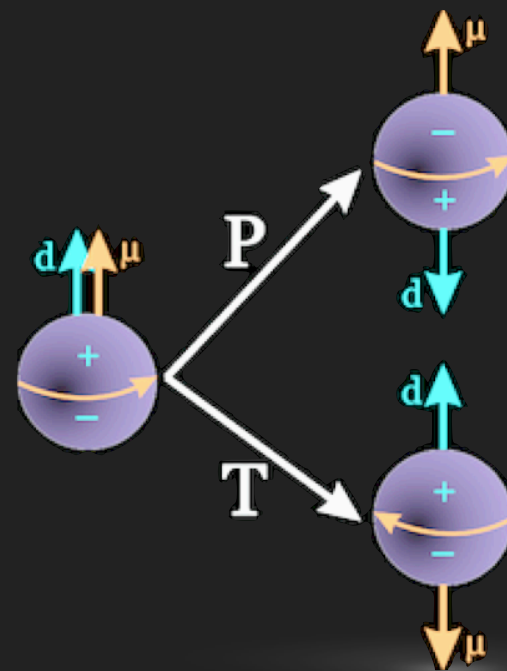


Diwan et al, Ann. Rev. Nucl. Part. Sci. 66, 47 (2016)



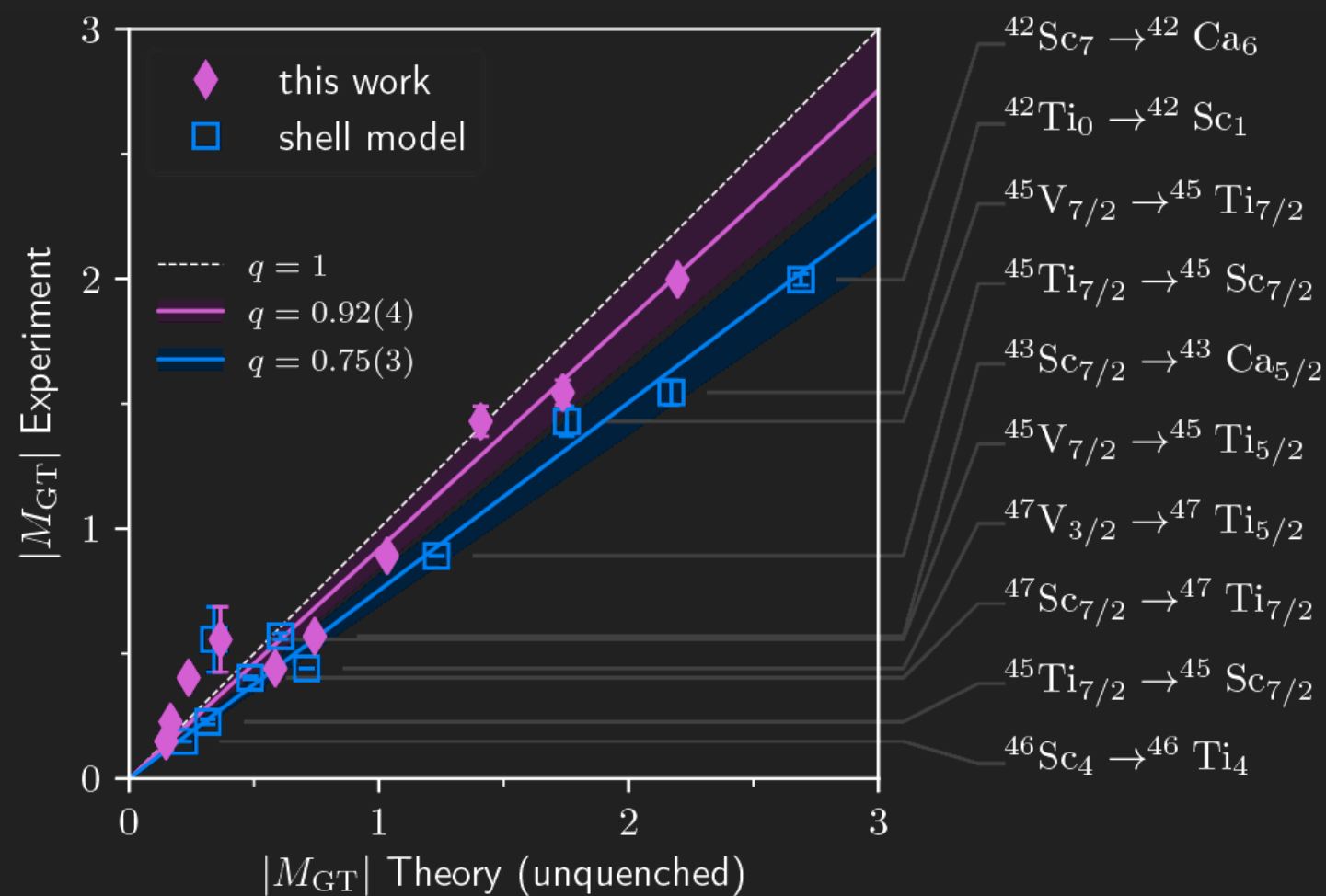
NUCLEI IN NEW PHYSICS

- ▶ Scalar currents
 - ▶ Dark matter direct detection
 - ▶ Lepton flavour violation: $\mu 2e$
 - ▶ Precision spectroscopy
- ▶ Tensor currents
 - ▶ Electric dipole moments of neutrons and nuclei
- ▶ Neutrinoless double beta decay



NUCLEAR UNCERTAINTIES

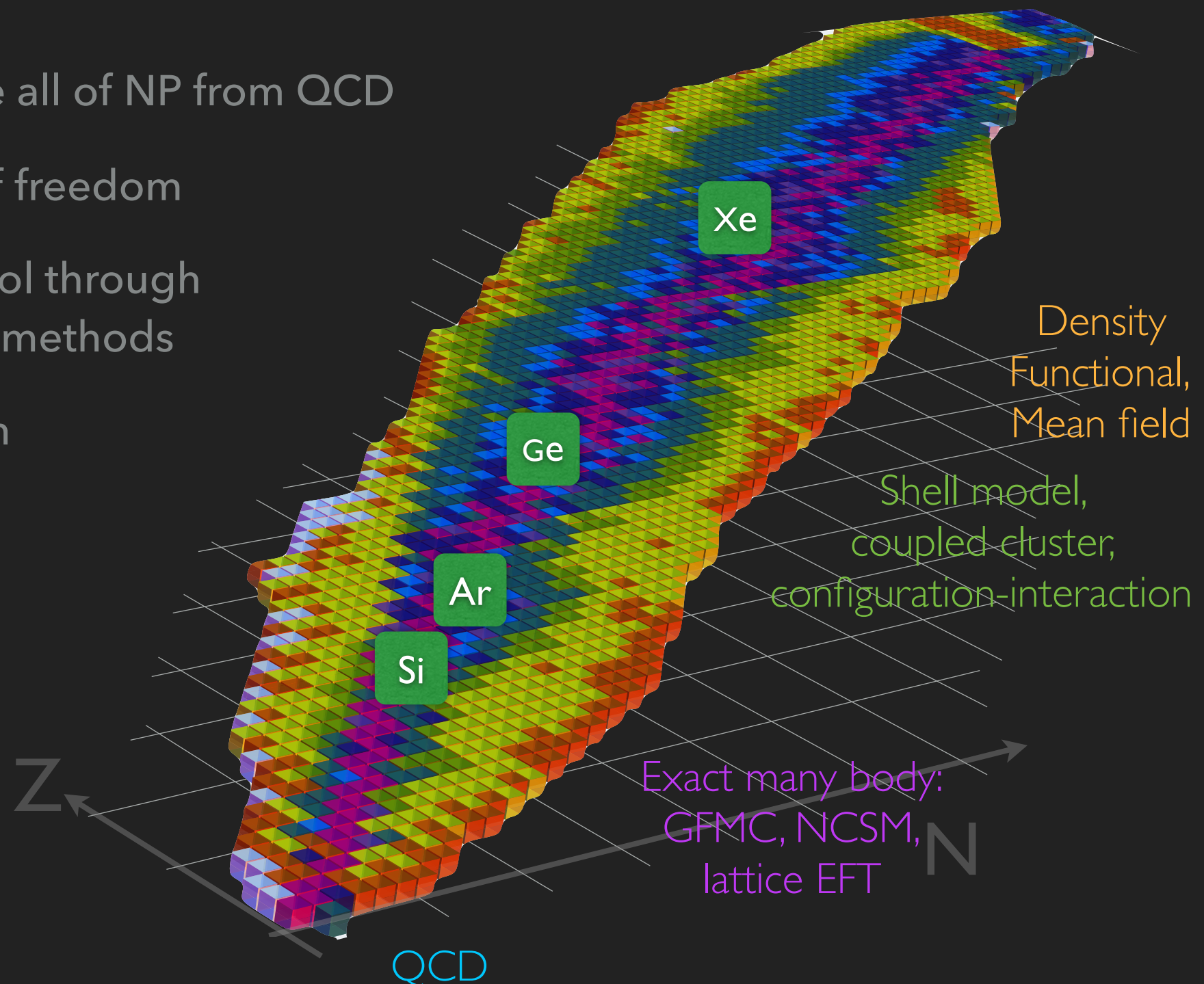
- ▶ How well do we know nuclear matrix elements?
- ▶ Gamow-Teller transitions in nuclei
 - ▶ Well measured for large range of nuclei ($30 < A < 60$)
 - ▶ Many nuclear structure calculations (shell-model,...) describe spectrum well
 - ▶ Matrix elements systematically off by 20-30%
 - ▶ Correct using 2 body currents
- ▶ Fundamental understanding from QCD



P Gysbers et al, Nature Phys. 15 (2019) no.5, 428-431

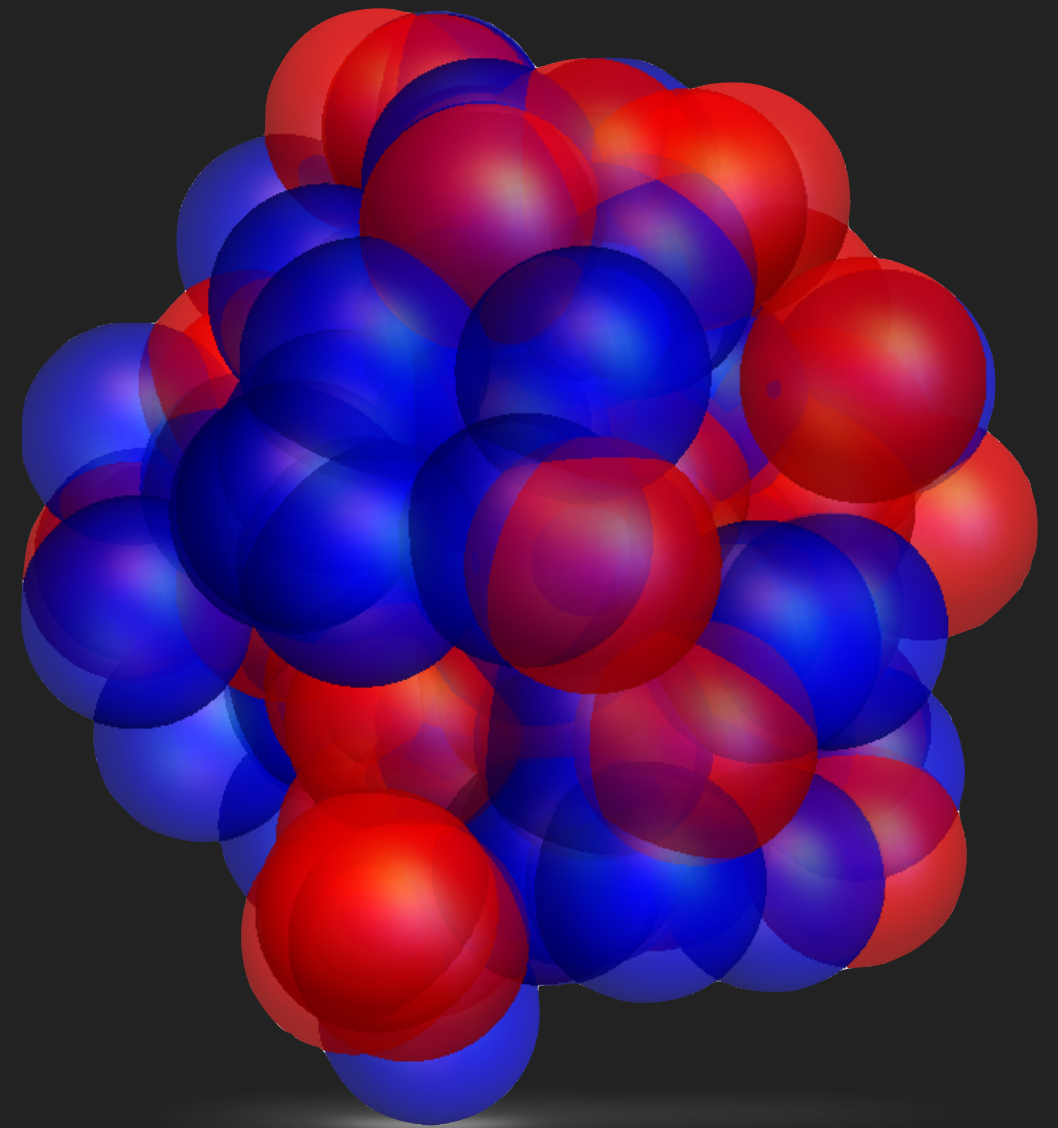
PRECISION NUCLEAR PHYSICS

- ▶ Very challenging to explore all of NP from QCD
- ▶ Exploit effective degrees of freedom
- ▶ Establish quantitative control through linkages between different methods
 - ▶ QCD forms a foundation determines few body interactions & matrix elements
 - ▶ Match existing EFT and many body techniques onto QCD



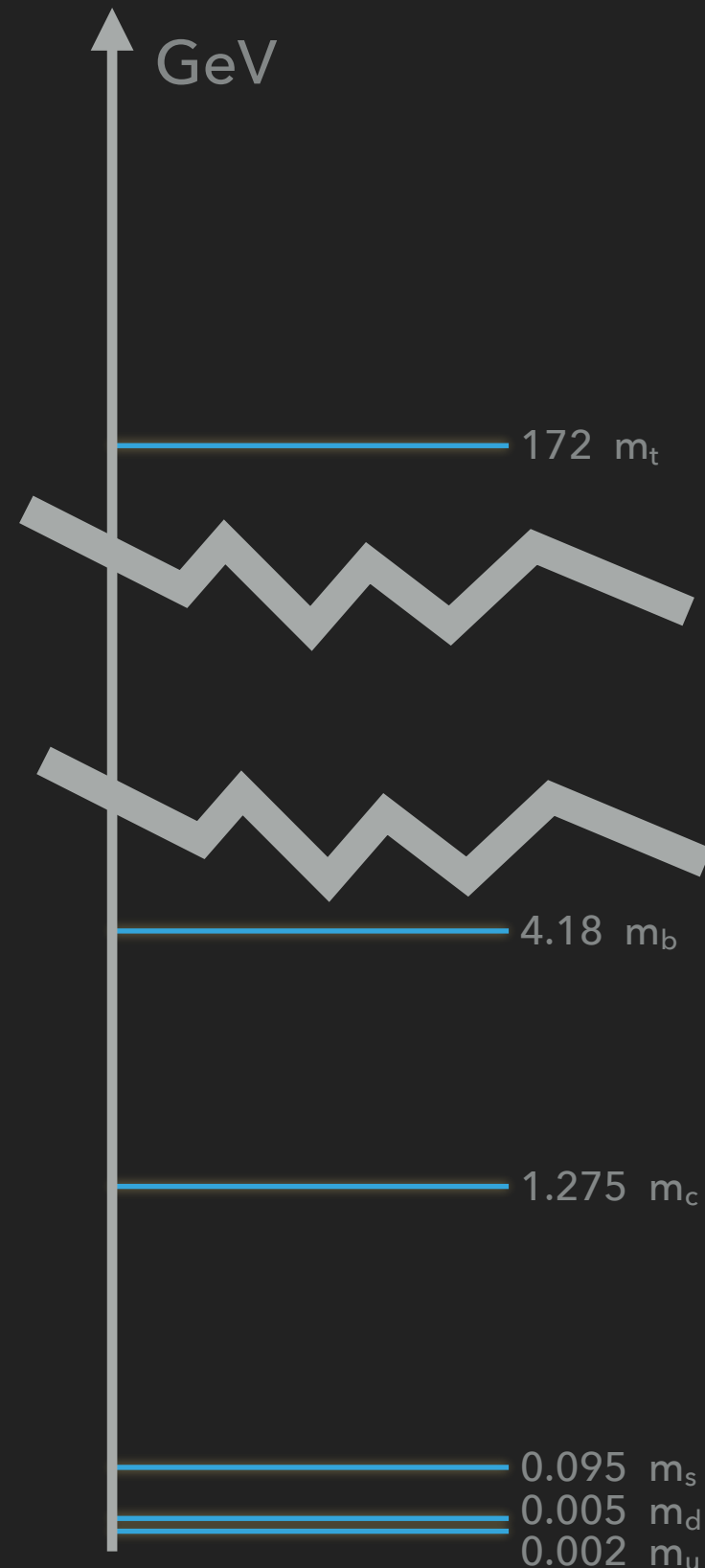
QCD FOR NUCLEAR PHYSICS

- ▶ Nuclear physics is Standard Model physics
 - ▶ Can compute the mass of lead nucleus ... in principle
- ▶ Complex physics
 - ▶ Wide range of scales
 - ▶ Closely spaced excitations
- ▶ Numerical challenges:
 - ▶ Statistical sampling
 - ▶ Contraction complexity



QCD FOR NUCLEAR PHYSICS

- ▶ Nuclear physics is Standard Model physics
 - ▶ Can compute the mass of lead nucleus ... in principle
- ▶ Complex physics
 - ▶ Wide range of scales
 - ▶ Closely spaced excitations
- ▶ Numerical challenges:
 - ▶ Statistical sampling
 - ▶ Contraction complexity



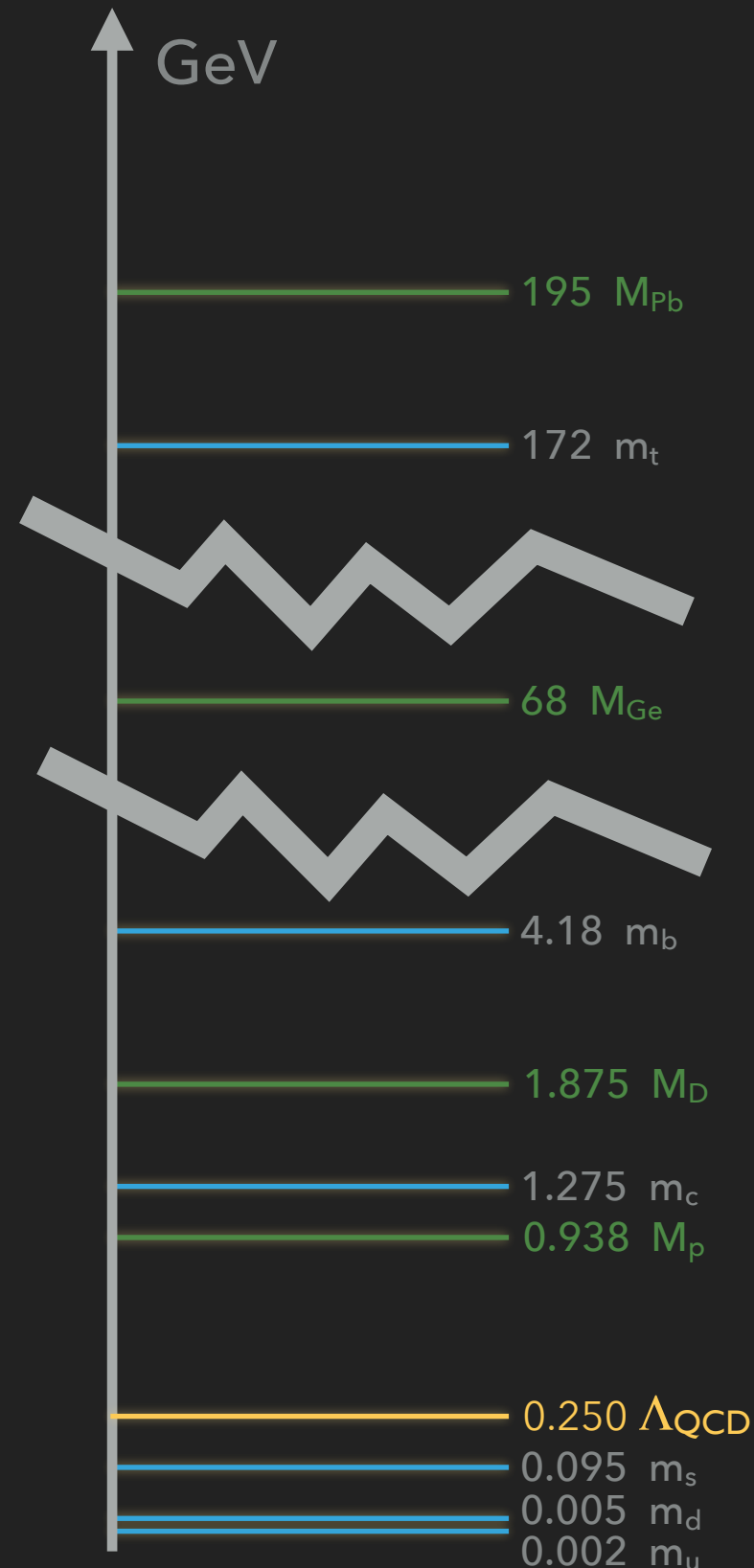
QCD FOR NUCLEAR PHYSICS

- ▶ Nuclear physics is Standard Model physics
 - ▶ Can compute the mass of lead nucleus ... in principle
- ▶ Complex physics
 - ▶ Wide range of scales
 - ▶ Closely spaced excitations
- ▶ Numerical challenges:
 - ▶ Statistical sampling
 - ▶ Contraction complexity



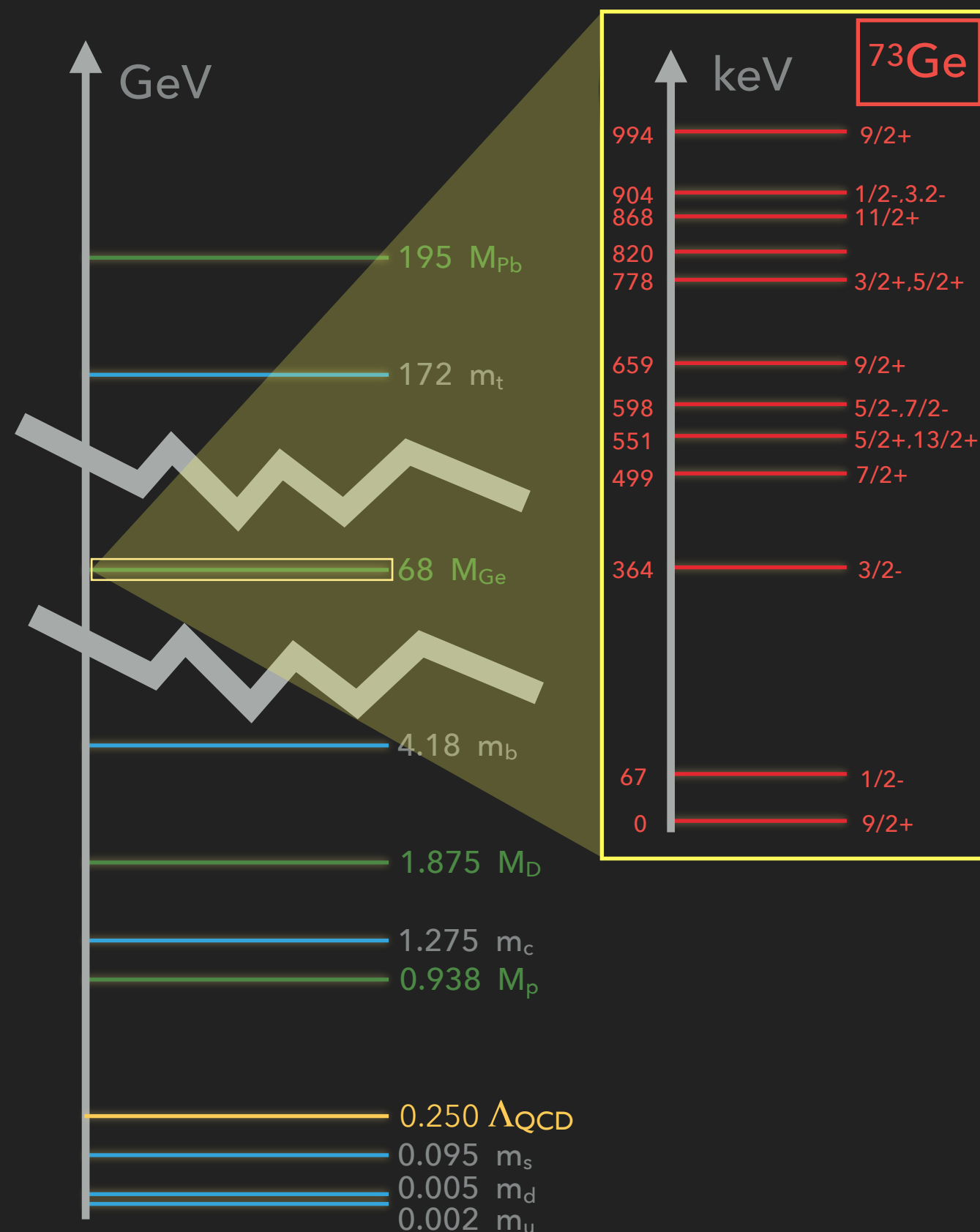
QCD FOR NUCLEAR PHYSICS

- ▶ Nuclear physics is Standard Model physics
 - ▶ Can compute the mass of lead nucleus ... in principle
- ▶ Complex physics
 - ▶ Wide range of scales
 - ▶ Closely spaced excitations
- ▶ Numerical challenges:
 - ▶ Statistical sampling
 - ▶ Contraction complexity



QCD FOR NUCLEAR PHYSICS

- ▶ Nuclear physics is Standard Model physics
 - ▶ Can compute the mass of lead nucleus ... in principle
- ▶ Complex physics
 - ▶ Wide range of scales
 - ▶ Closely spaced excitations
- ▶ Numerical challenges:
 - ▶ Statistical sampling
 - ▶ Contraction complexity



NUCLEI

- ▶ New algorithms enabling study of nuclei

- ▶ Efficient contractions

- ▶ Graph theory &

recursions [WD & Savage 2011;

WD & Orginos 2012, Doi&Endres 2012;

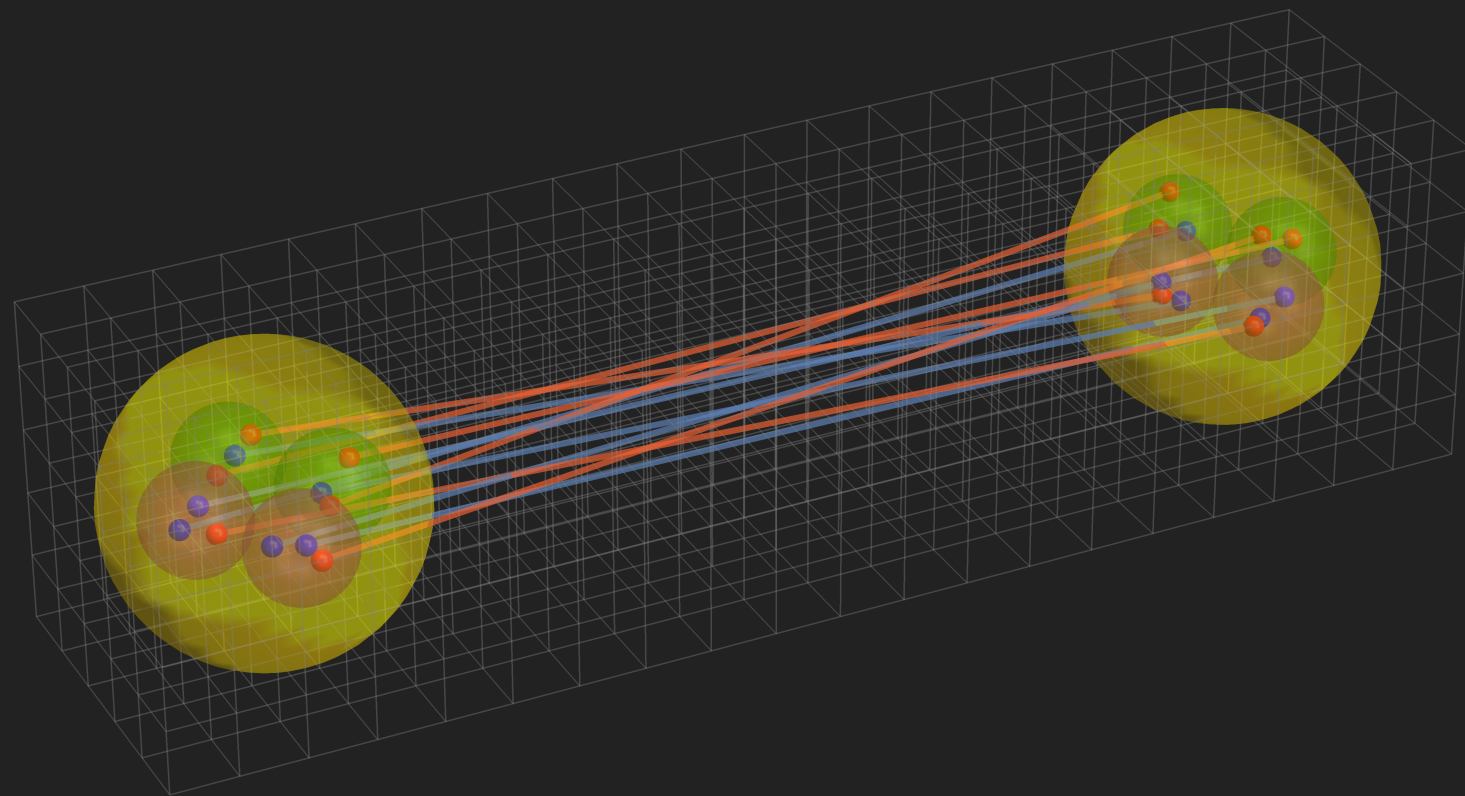
Gunther&Varnhorst 2013; WD & Vachaspati 2014]

- ▶ Better statistical estimators

[WD & Endres 2014, Wagman & Savage 2016;

WD, Kanwar & Wagman 2018, Murphy et al. 2019]

- ▶ ... and lots of computing!



$$\frac{\text{cost}({}^N A_Z)}{\text{cost}(\text{proton})} \sim (2Z + N)!(2N + Z)!/2$$

Cost Ratio

| Nucleus | Naive | Optimised |
|---------------------|-------------|-----------|
| ${}^4\text{He}$ | 250,000 | ~100 |
| ${}^8\text{Be}$ | 10^{31} | 10^7 |
| ${}^{208}\text{Pb}$ | 10^{1300} | ? |

NPLQCD

- ▶ Case study QCD with unphysical quark masses ($m_\pi \sim 800$ MeV, 450 MeV)

1. Spectrum and scattering of light nuclei ($A < 5$) [PRD 87 (2013), 034506]
2. Nuclear structure: magnetic moments, polarisabilities ($A < 5$) [PRL 113, 252001 (2014), PRL 116, 112301 (2016)]
3. Nuclear reactions: $np \rightarrow d\gamma$ [PRL 115, 132001 (2015)]
4. Gamow-Teller transitions: $pp \rightarrow d e \nu$, $g_A(^3\text{H})$ [PRL 119 062002 (2017)]
5. Double β decay: $pp \rightarrow nn$ [PRL 119, 062003 (2017)]
6. Parton structure ($A < 4$) [PRD 96 094512 (2017)]
7. Scalar/tensor currents ($A < 4$) [PRL 2018]



+ Arjun Gambhir

NPLQCD

- ▶ Case study QCD with unphysical quark masses ($m_\pi \sim 800$ MeV, 450 MeV)

1. Spectrum and scattering of light nuclei ($A < 5$) [PRD 87 (2013), 034506]
2. Nuclear structure: magnetic moments, polarisabilities ($A < 5$) [PRL 113, 252001 (2014), PRL 116, 112301 (2016)]
3. Nuclear reactions: $np \rightarrow d\gamma$ [PRL 115, 132001 (2015)]
4. Gamow-Teller transitions: $pp \rightarrow d e \nu$, $g_A(^3\text{H})$ [PRL 119 062002 (2017)]
5. Double β decay: $pp \rightarrow nn$ [PRL 119, 062003 (2017)]
6. Parton structure ($A < 4$) [PRD 96 094512 (2017)]
7. Scalar/tensor currents ($A < 4$) [PRL 2018]

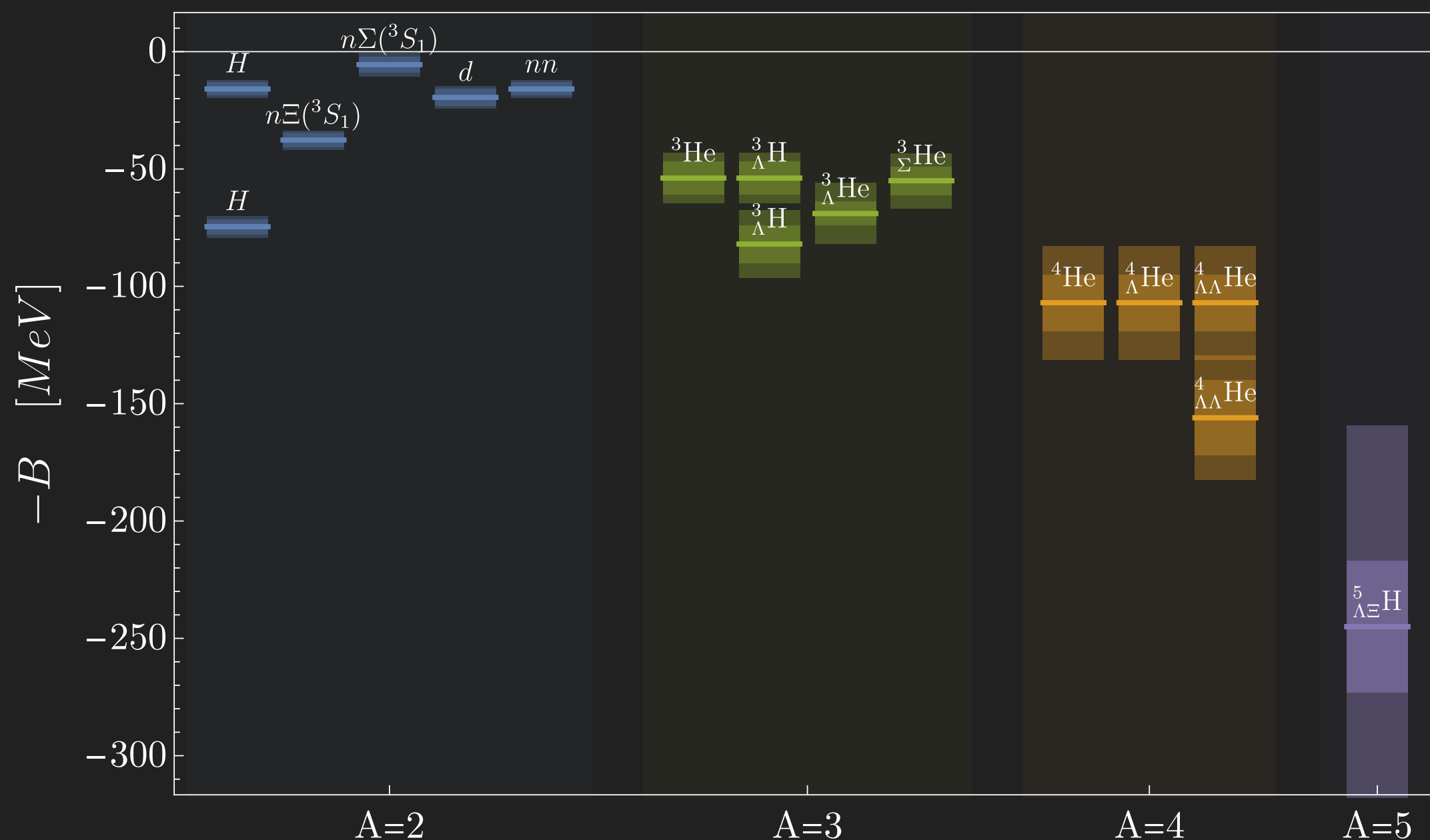
new since 2017



+ Arjun Gambhir

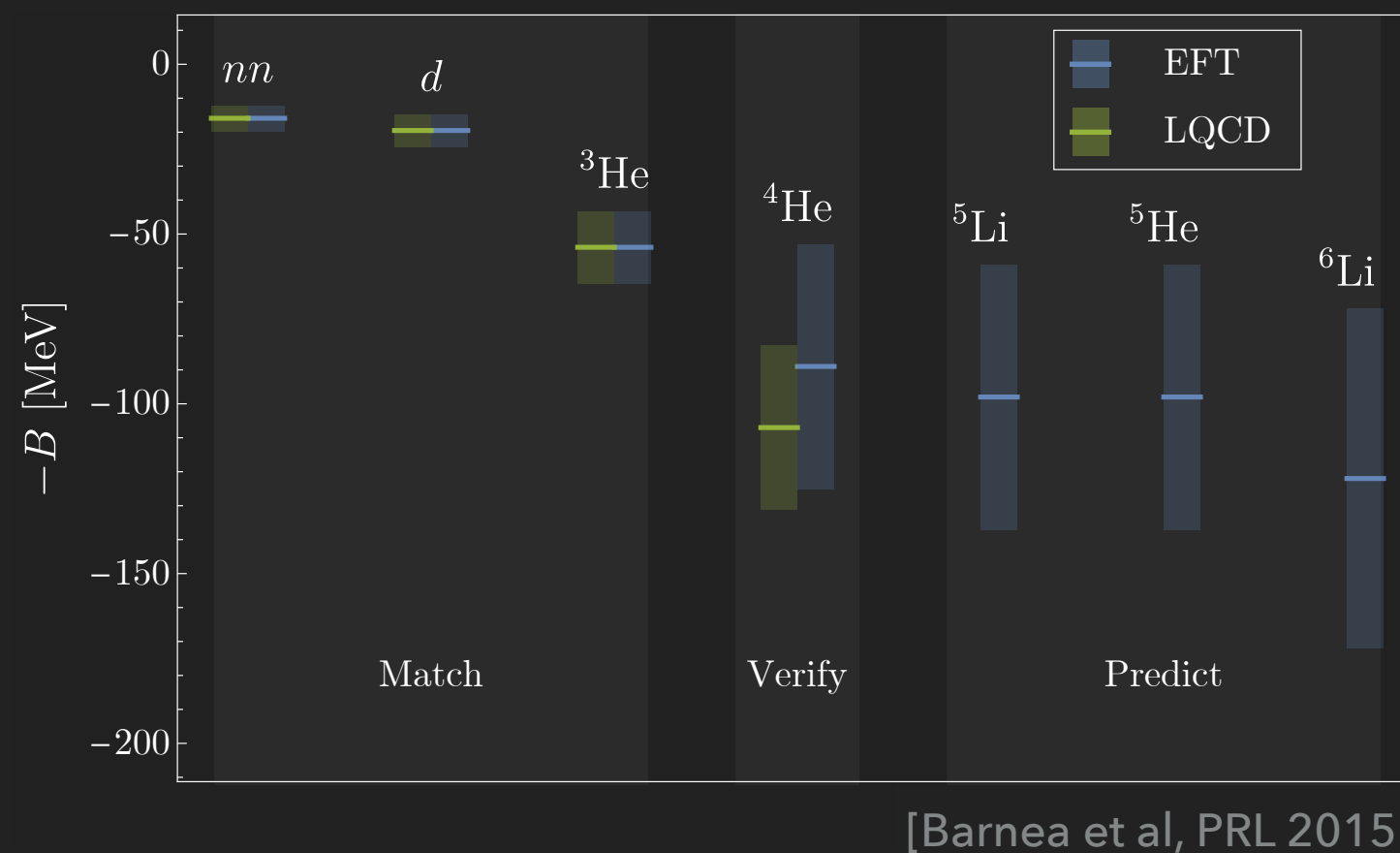
NUCLEI (IN A HEAVY QUARK UNIVERSE, $M_\pi \sim 800$ MEV)

- 2013: first QCD calculation of nuclei (heavy masses as numerically cheaper)



NUCLEI (IN A HEAVY QUARK UNIVERSE, $M_\pi \sim 800$ MEV)

- ▶ Combine LQCD and nucleon based many-body effective field theory (EFT) methods
- ▶ Matching to LQCD determines NN, NNN interactions: allows predictions for larger nuclei



**MORE EFFECTIVE:
DIRECTLY MATCH FV
ENERGIES IN LQCD AND
EFT**

- ▶ Further studies extend to complex nuclei such as ${}^{16}\text{O}$ [Contessi et al, Bansal et al.]

MAGNETIC STRUCTURE

- ▶ Hadron/nuclear energies are modified by presence of fixed external fields

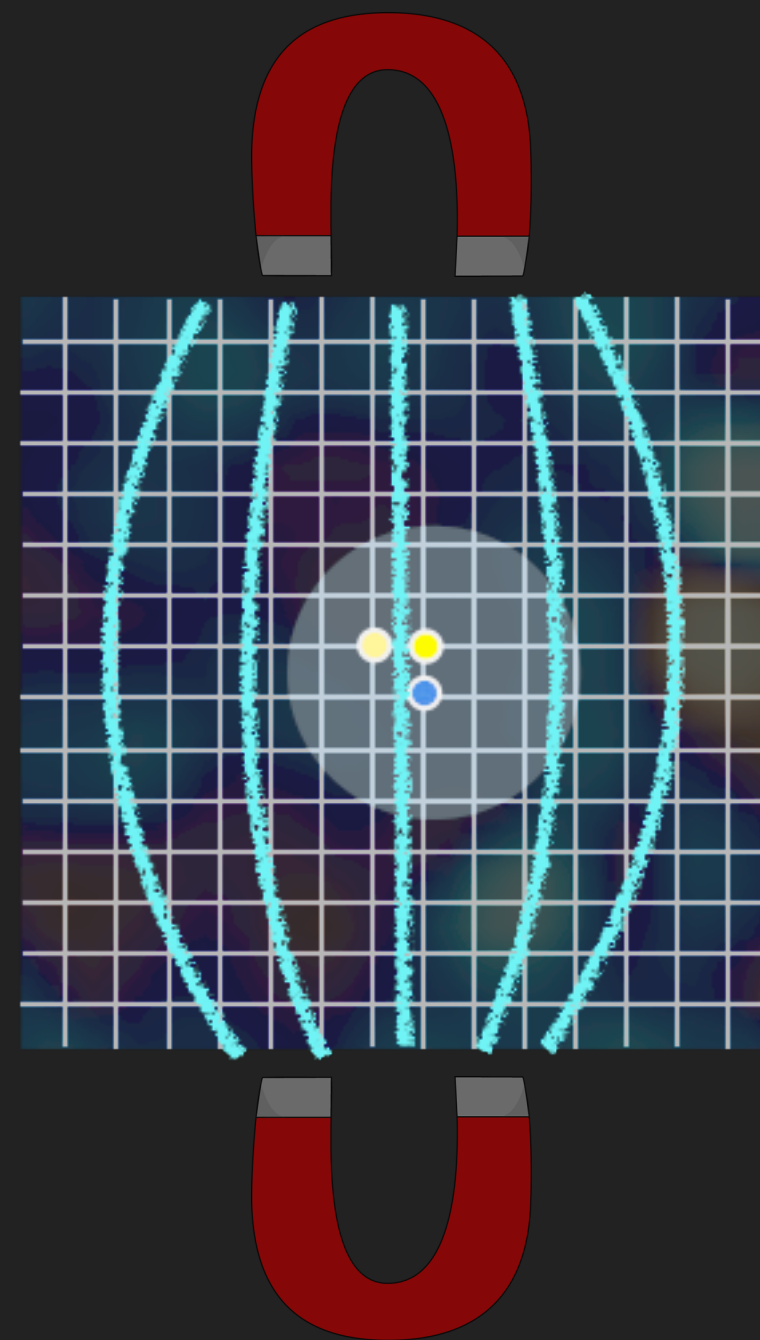
- ▶ Eg: fixed B field

$$E_{h;j_z}(\mathbf{B}) = \sqrt{M_h^2 + (2n+1)|Q_h e B|} - \boldsymbol{\mu}_h \cdot \mathbf{B} - 2\pi\beta_h^{(M0)}|\mathbf{B}|^2 + \dots$$

- ▶ QCD calculations with multiple fields enable extraction of coefficients of response

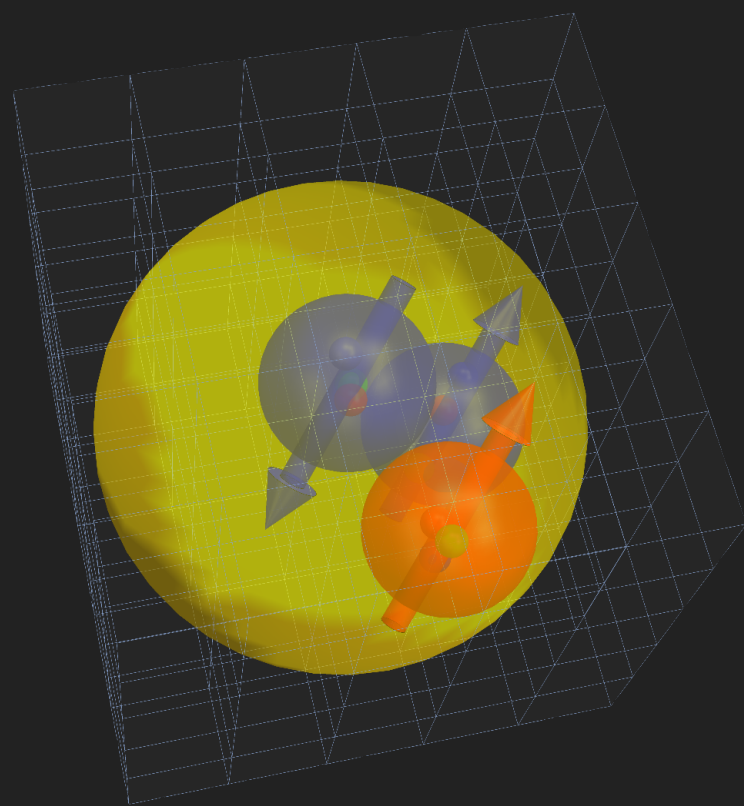
- ▶ Magnetic moments, polarisabilities, ...

- ▶ Similar techniques to study EW interactions, DM interactions, twist-2 matrix elements

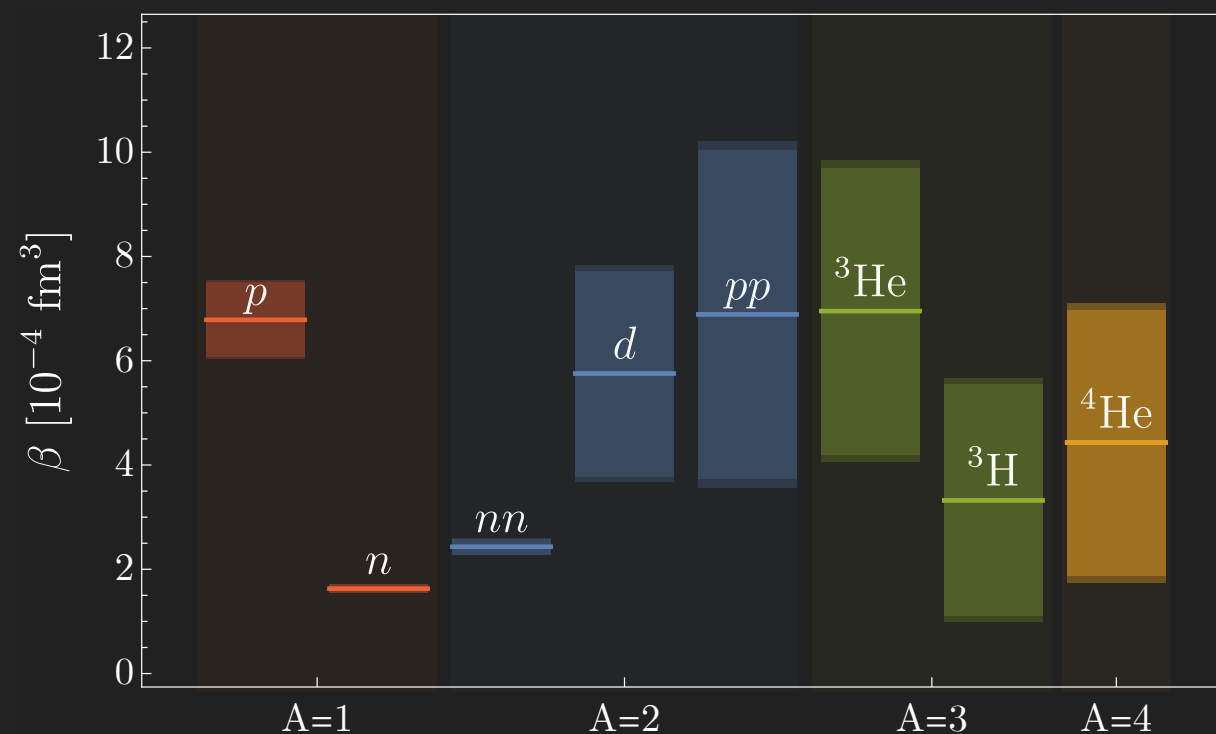
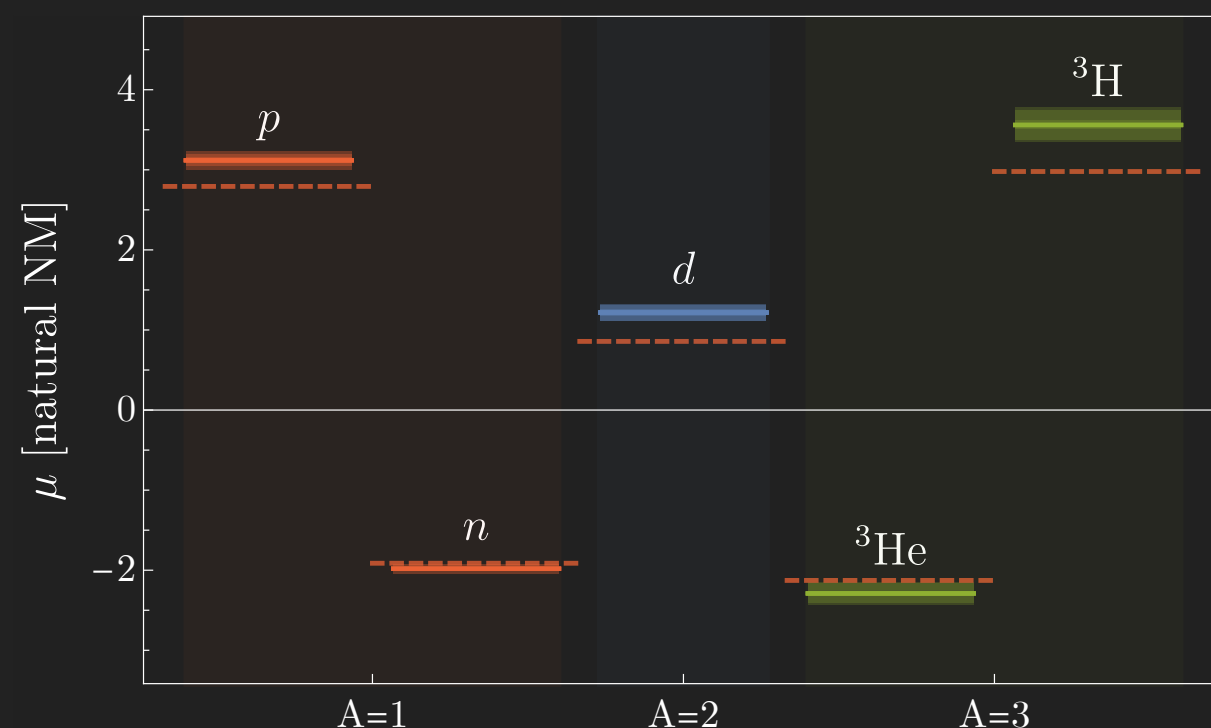


MAGNETIC STRUCTURE

- ▶ LQCD calculation of nuclear magnetic moments (μ) and magnetic polarisabilities (β , deformation in B field)
- ▶ Simple shell model expectations

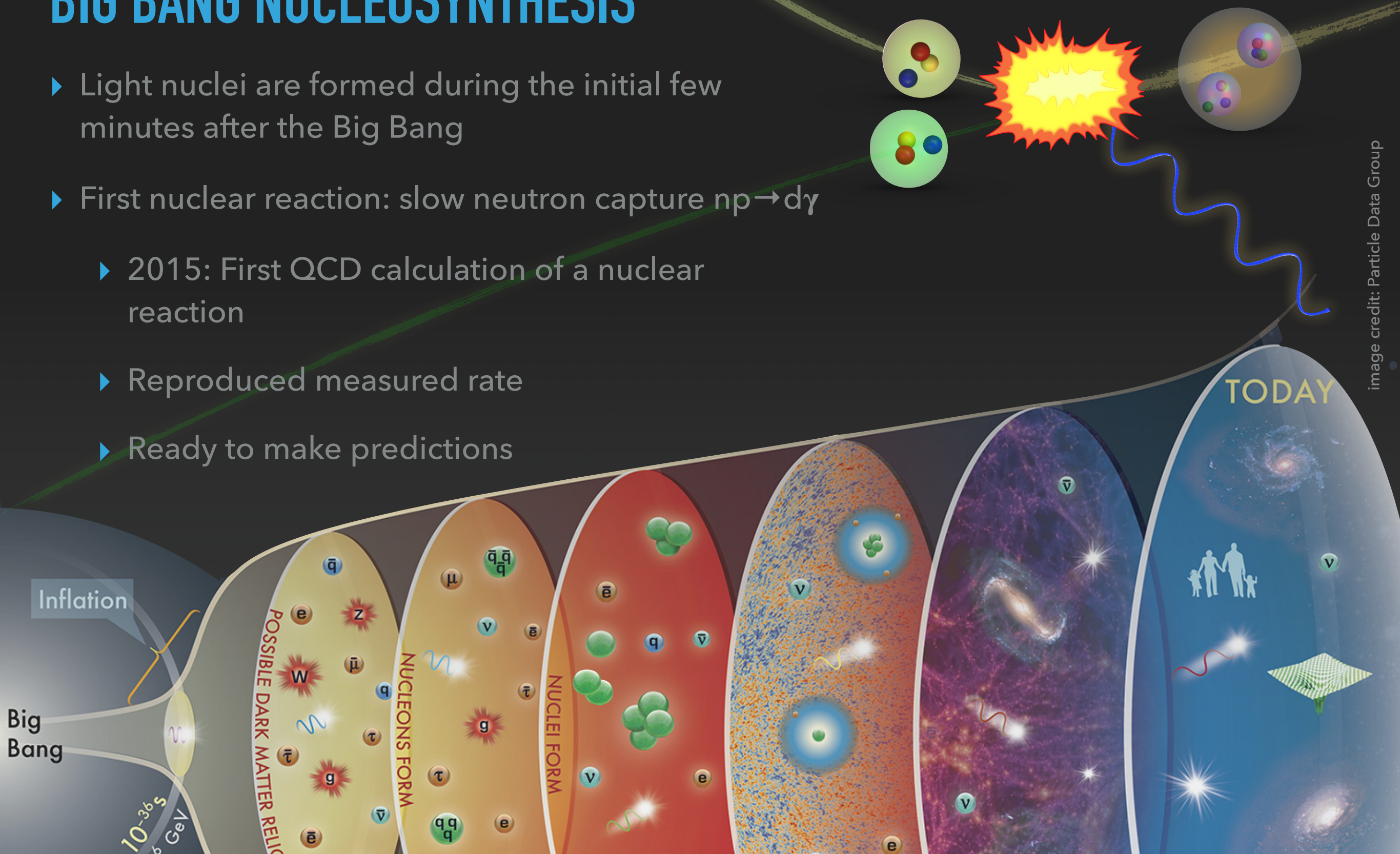


- ▶ Lattice results suggest heavy quark mass nuclei are shell-model like!



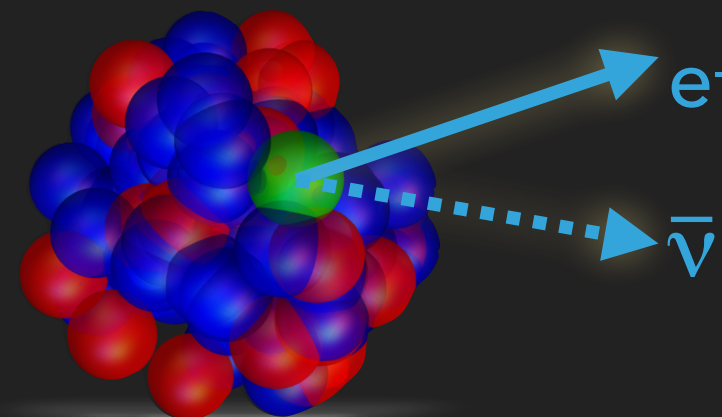
BIG BANG NUCLEOSYNTHESIS

- ▶ Light nuclei are formed during the initial few minutes after the Big Bang
- ▶ First nuclear reaction: slow neutron capture $np \rightarrow d\gamma$
 - ▶ 2015: First QCD calculation of a nuclear reaction
 - ▶ Reproduced measured rate
 - ▶ Ready to make predictions

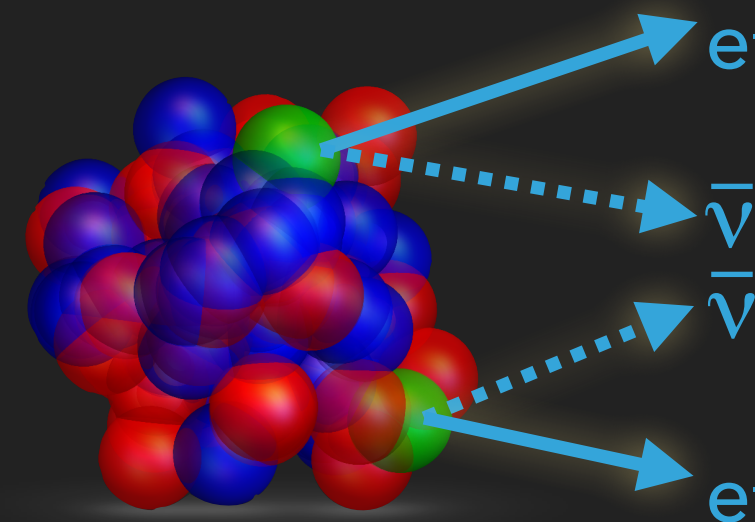
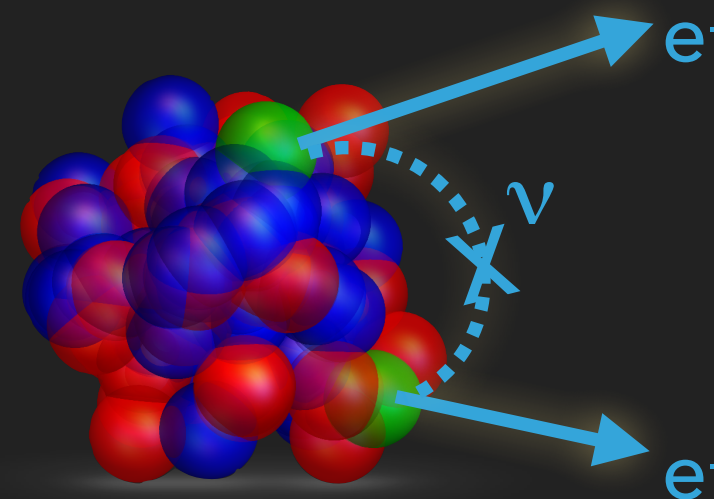


ELECTROWEAK PROCESSES

- ▶ Single β -decay
LQCD calculation of decay of tritium
- ▶ Double β -decay
 - ▶ Neutrinoless case is rarest process observed
 - ▶ Neutrinoless case
 - ▶ Majorana particles? Lepton number violation? Baryon asymmetry?
 - ▶ Rates depend on nuclear matrix elements
 - ▶ Currently quite uncertain but important for design of future DBD search experiments
- ▶ Proton-Proton fusion powering the Sun



Beta decay

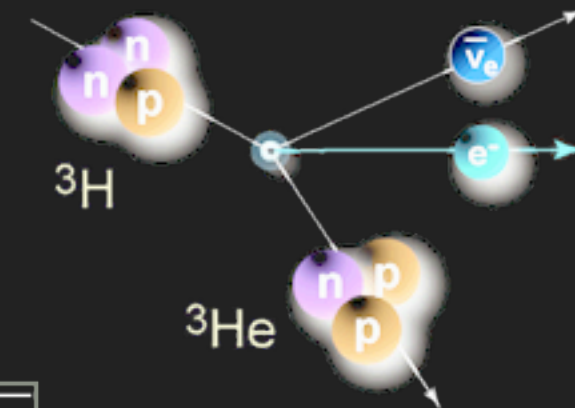
Two neutrino double β -decayNeutrinoless double β -decay

TRITIUM BETA DECAY

- ▶ Tritium decay half life

$$\frac{(1 + \delta_R) f_V}{K/G_V^2} t_{1/2}^{\text{half-life}} = \frac{1}{\langle \mathbf{F} \rangle^2 + f_A/f_V g_A^2 \langle \mathbf{GT} \rangle^2}$$

known from theory or expt.

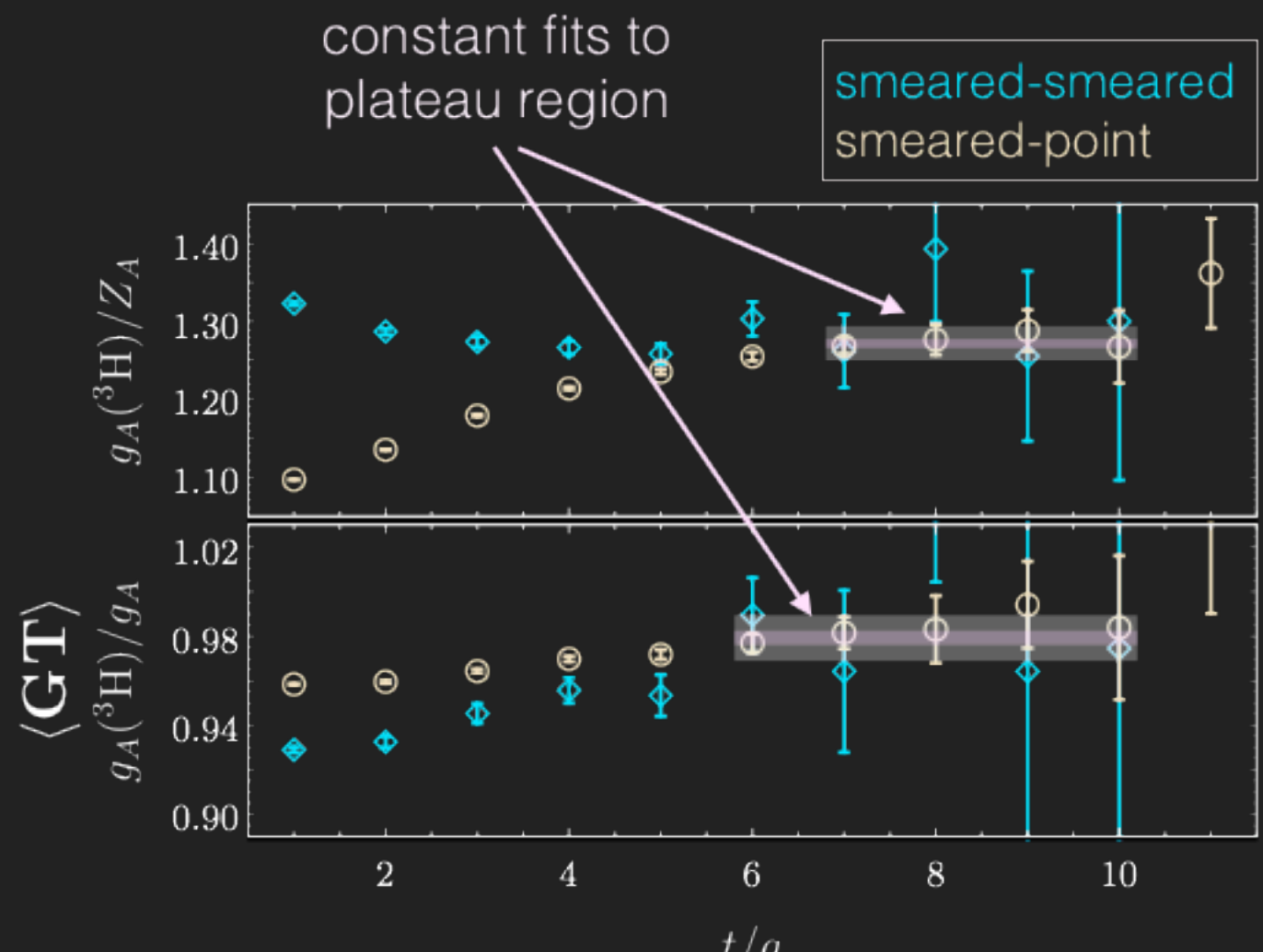


- ▶ Biggest uncertainty in

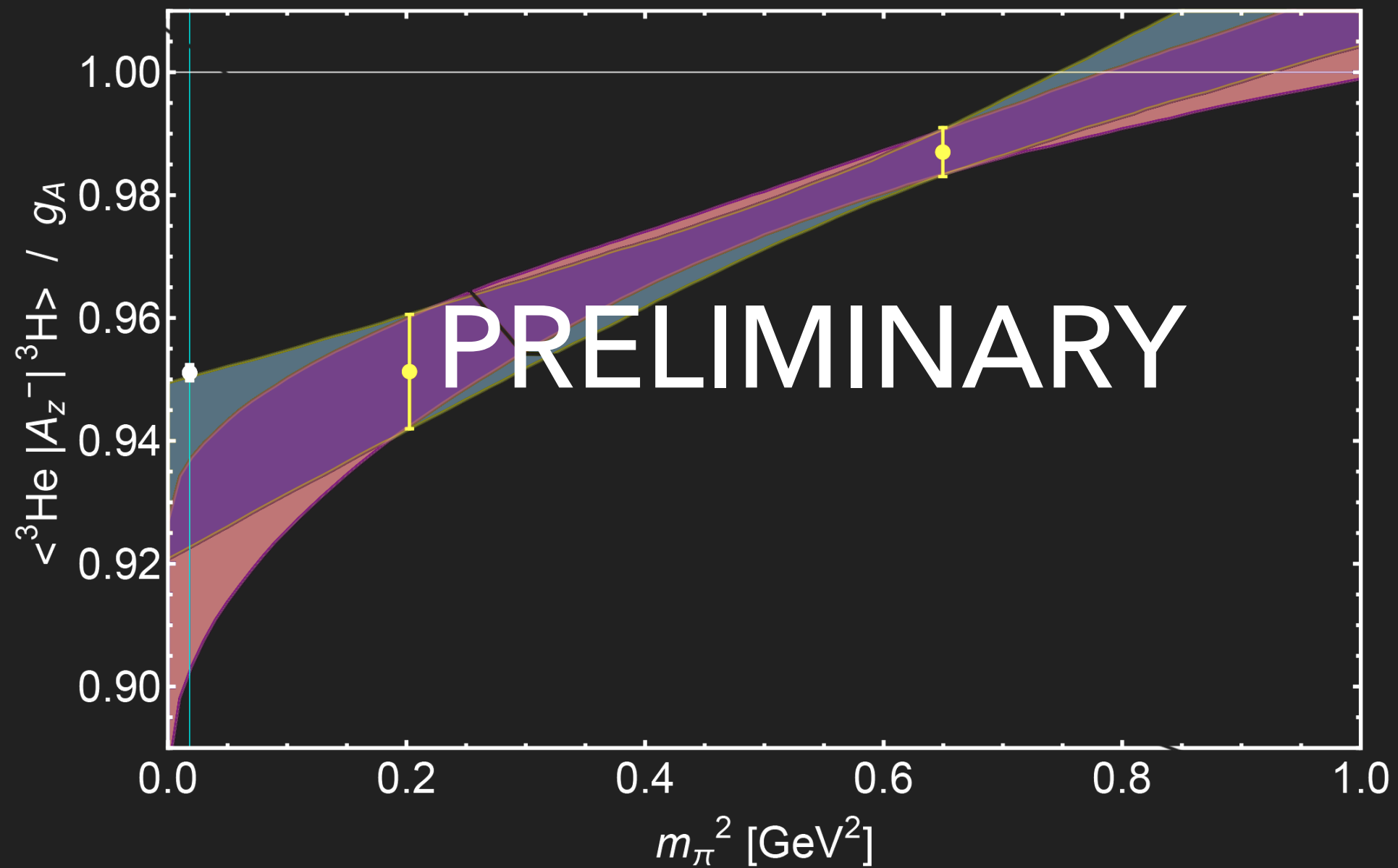
$$g_A \langle \mathbf{GT} \rangle = \langle {}^3\text{He} | \bar{\mathbf{q}} \gamma_{\mathbf{k}} \gamma_5 \tau^- \mathbf{q} | {}^3\text{H} \rangle$$

- ▶ Form ratios of correlators in axial background fields to extract QCD matrix element

$$\frac{\overline{R}_{3\text{H}}(t)}{\overline{R}_p(t)} \xrightarrow{t \rightarrow \infty} \frac{g_A({}^3\text{H})}{g_A} = \langle \mathbf{GT} \rangle$$

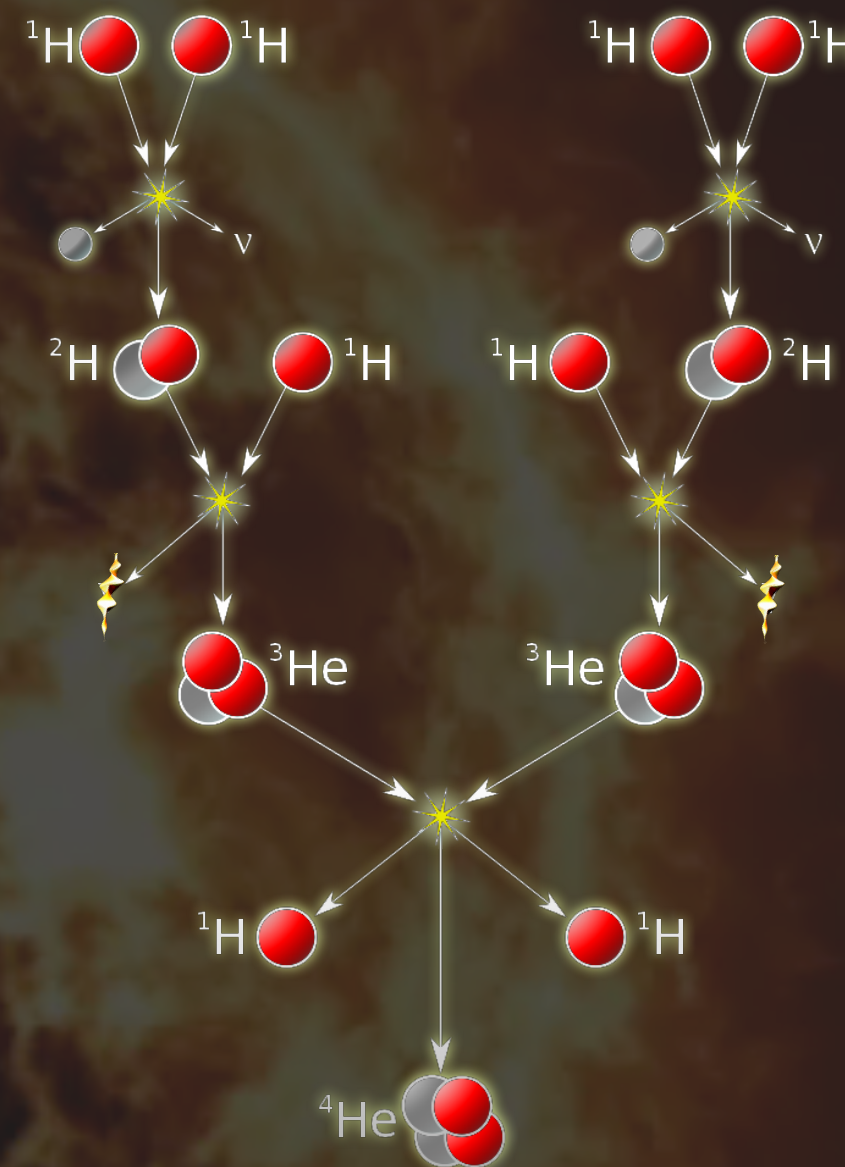


MASS DEPENDENCE OF TRITON AXIAL CHARGE



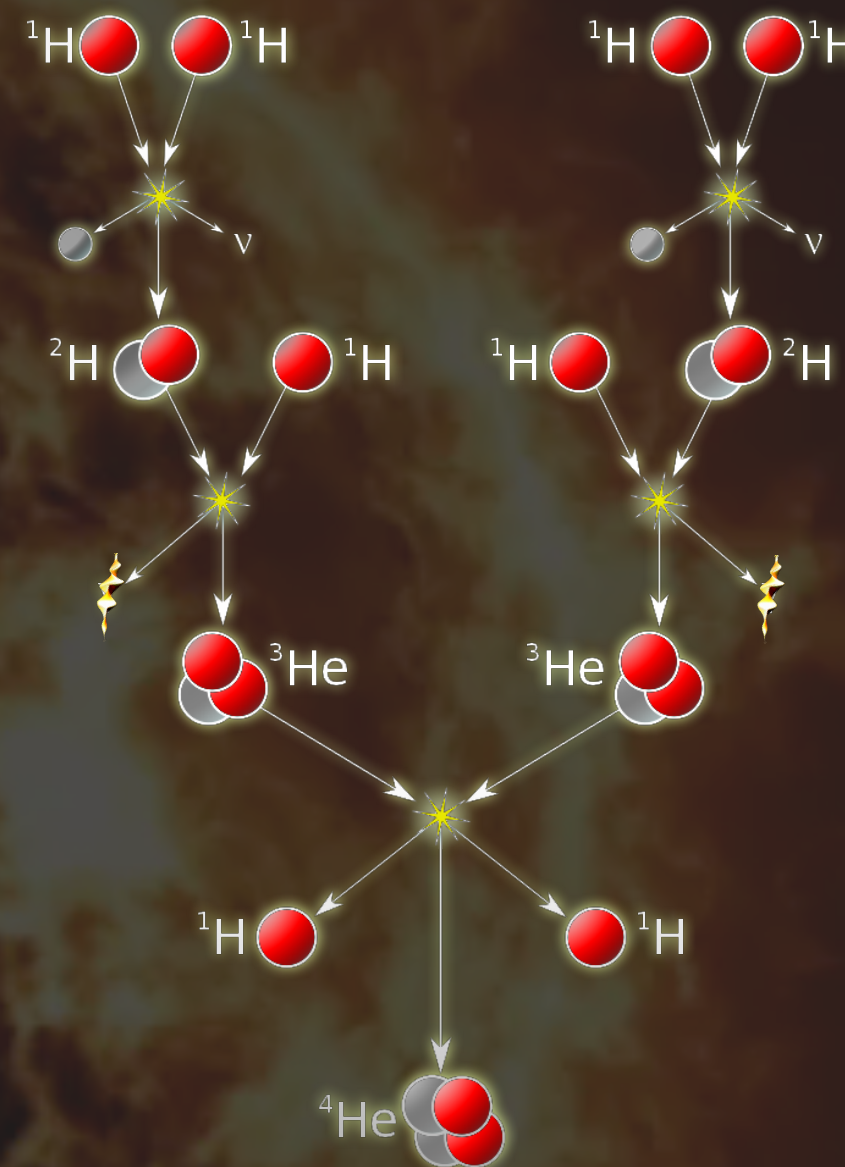
PROTON-PROTON FUSION

- ▶ First step in chain of reactions powering stars like the sun
- ▶ Intricate process involve all three SM forces
- ▶ Difficult to measure (Coulomb barrier)
- ▶ 2017: LQCD calculation of pp fusion rate
 - ▶ Uncertainties competitive with phenomenological extractions
 - ▶ Next generation calculations will improve precision
 - ▶ Improve solar modelling



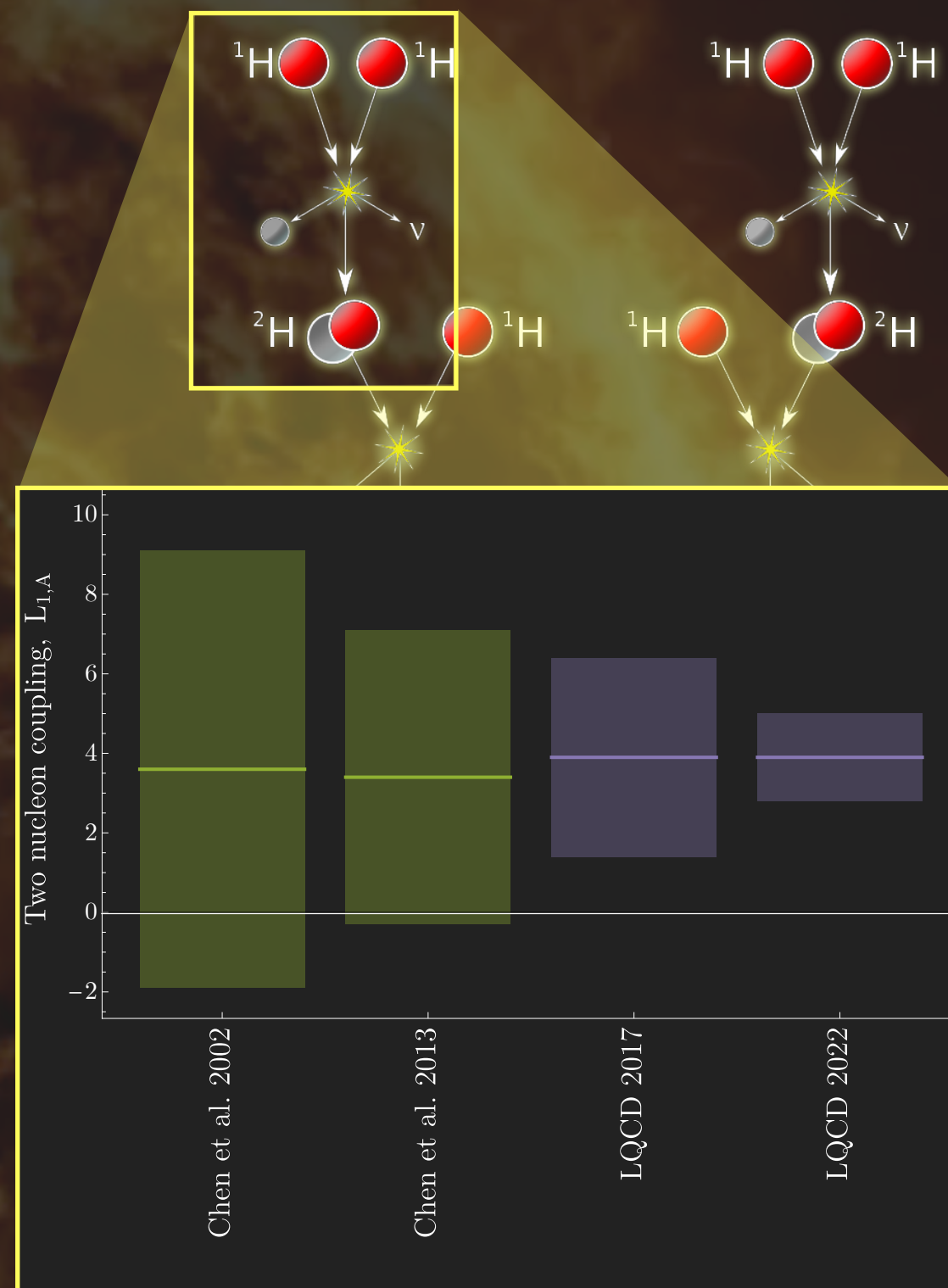
PROTON-PROTON FUSION

- ▶ First step in chain of reactions powering stars like the sun
- ▶ Intricate process involve all three SM forces
- ▶ Difficult to measure (Coulomb barrier)
- ▶ 2017: LQCD calculation of pp fusion rate
 - ▶ Uncertainties competitive with phenomenological extractions
 - ▶ Next generation calculations will improve precision
 - ▶ Improve solar modelling



PROTON-PROTON FUSION

- ▶ First step in chain of reactions powering stars like the sun
- ▶ Intricate process involve all three SM forces
- ▶ Difficult to measure (Coulomb barrier)
- ▶ 2017: LQCD calculation of pp fusion rate
 - ▶ Uncertainties competitive with phenomenological extractions
 - ▶ Next generation calculations will improve precision
 - ▶ Improve solar modelling



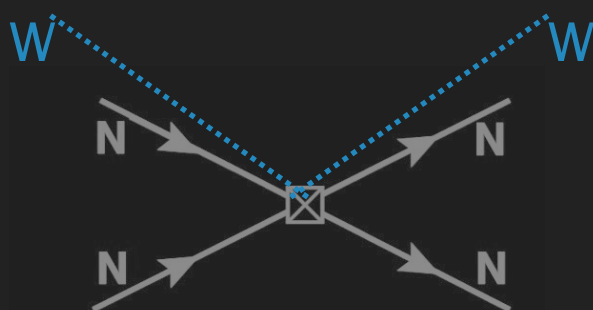
DOUBLE BETA DECAY

- QCD calculation of subprocess

$$nn \rightarrow ppe^-e^-\bar{\nu}\bar{\nu}$$

[NPLQCD, PRL 2017b]

- Revealed significant nuclear effects (even beyond g_A quenching)



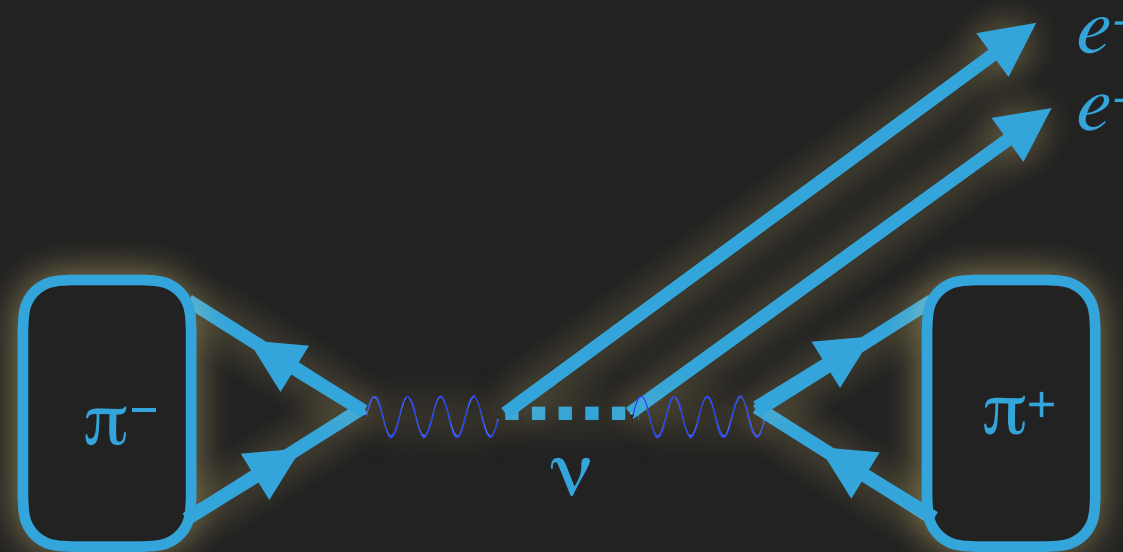
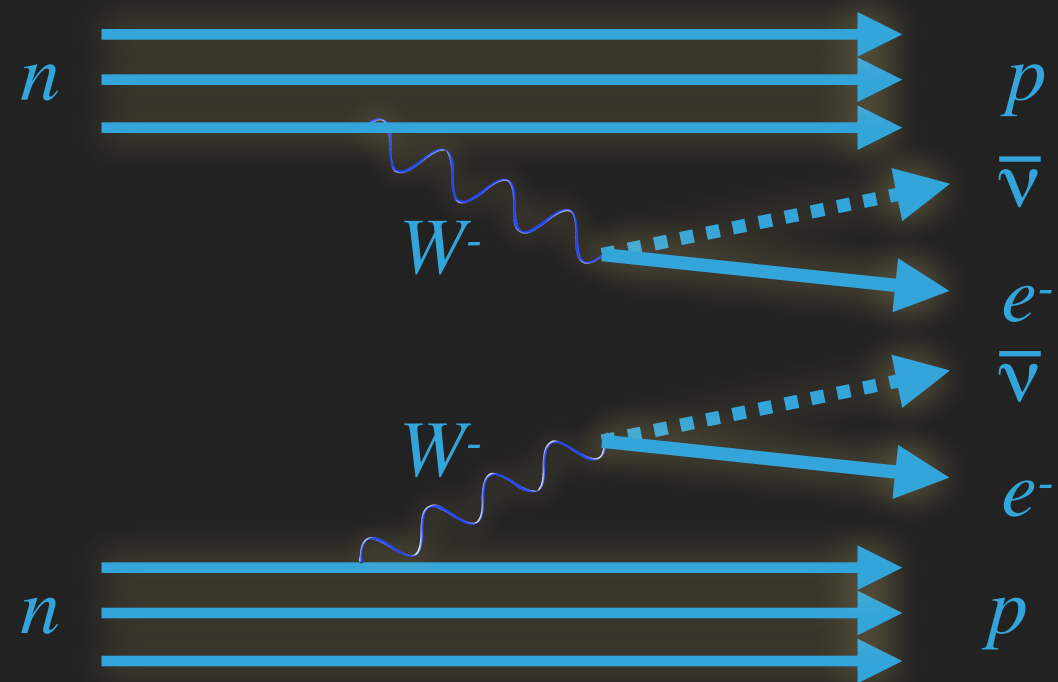
- Beginning calculations of neutrinoless processes

[WD, Murphy 1811.0554]

- Disallowed pion transition as a test

$$\pi^- \rightarrow \pi^+ e^- e^-$$

- Light nuclei are next



DARK MATTER INTERACTIONS

- ▶ DM direct detection experiments search for recoil of nucleus from DM scattering

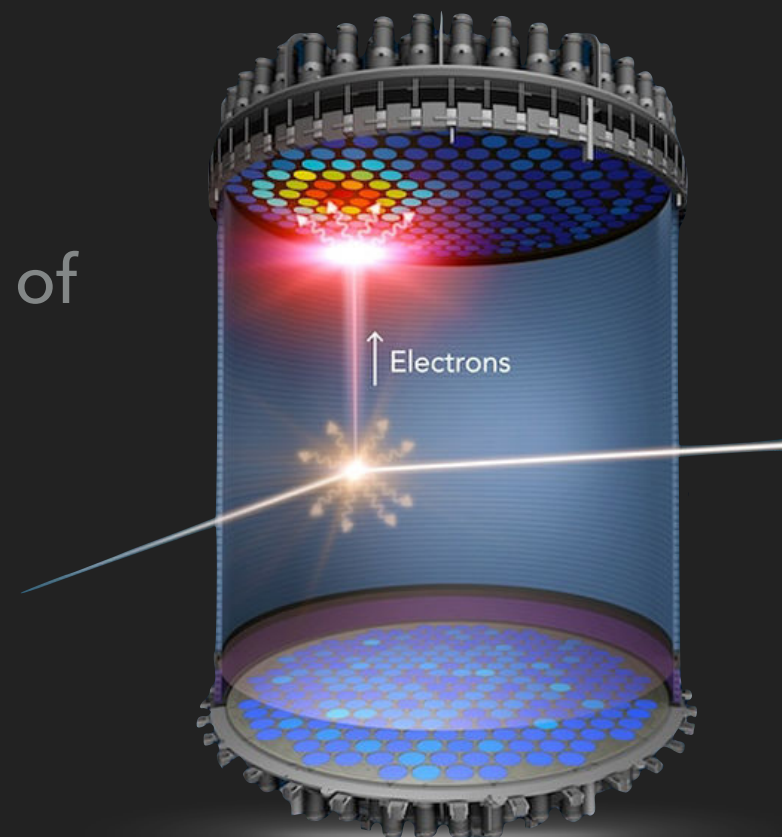
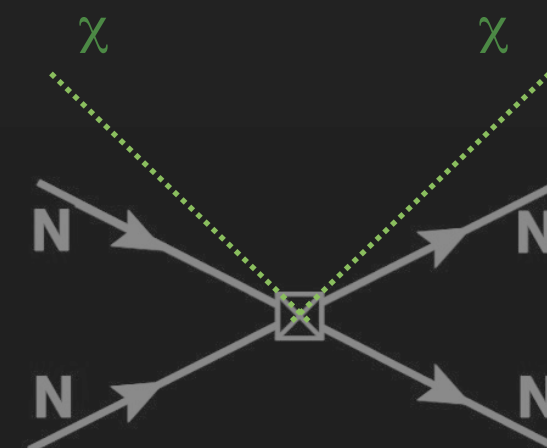
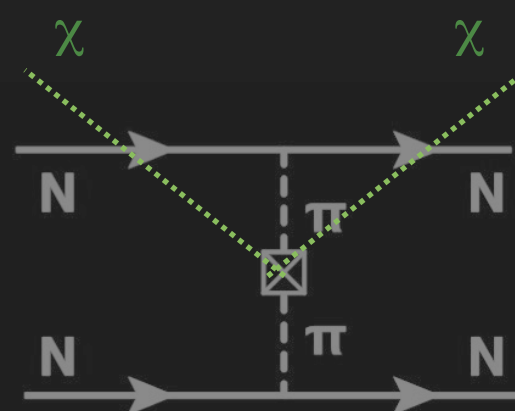
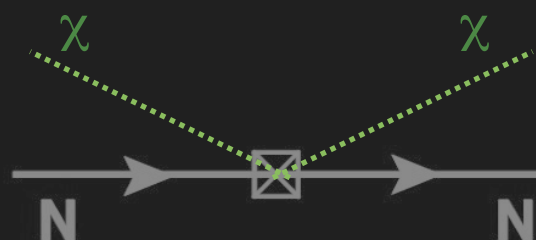
- ▶ One popular class of DM interactions is through scalar exchange

$$\mathcal{L} = \frac{G_F}{2} \sum_q \kappa_q (\bar{\chi}\chi)(\bar{q}q)$$

- ▶ Direct detection depends on nuclear matrix element

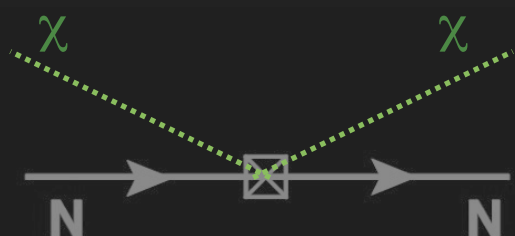
$$\overline{m} \langle Z, N | \bar{u}u + \bar{d}d | Z, N \rangle$$

- ▶ At hadronic/nuclear level



NUCLEON SCALAR COUPLING

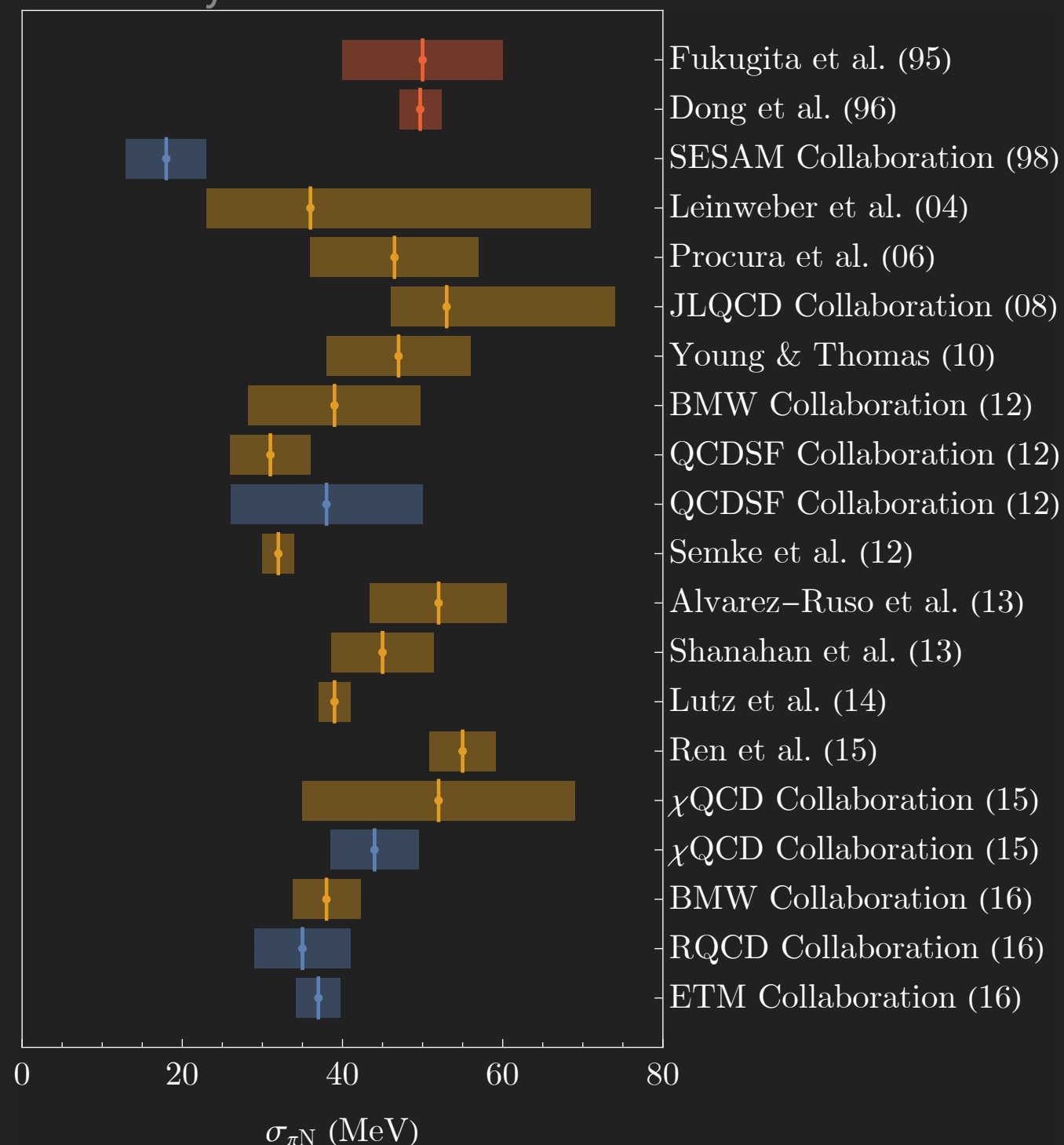
► Single nucleon contribution



► Calculated in LQCD

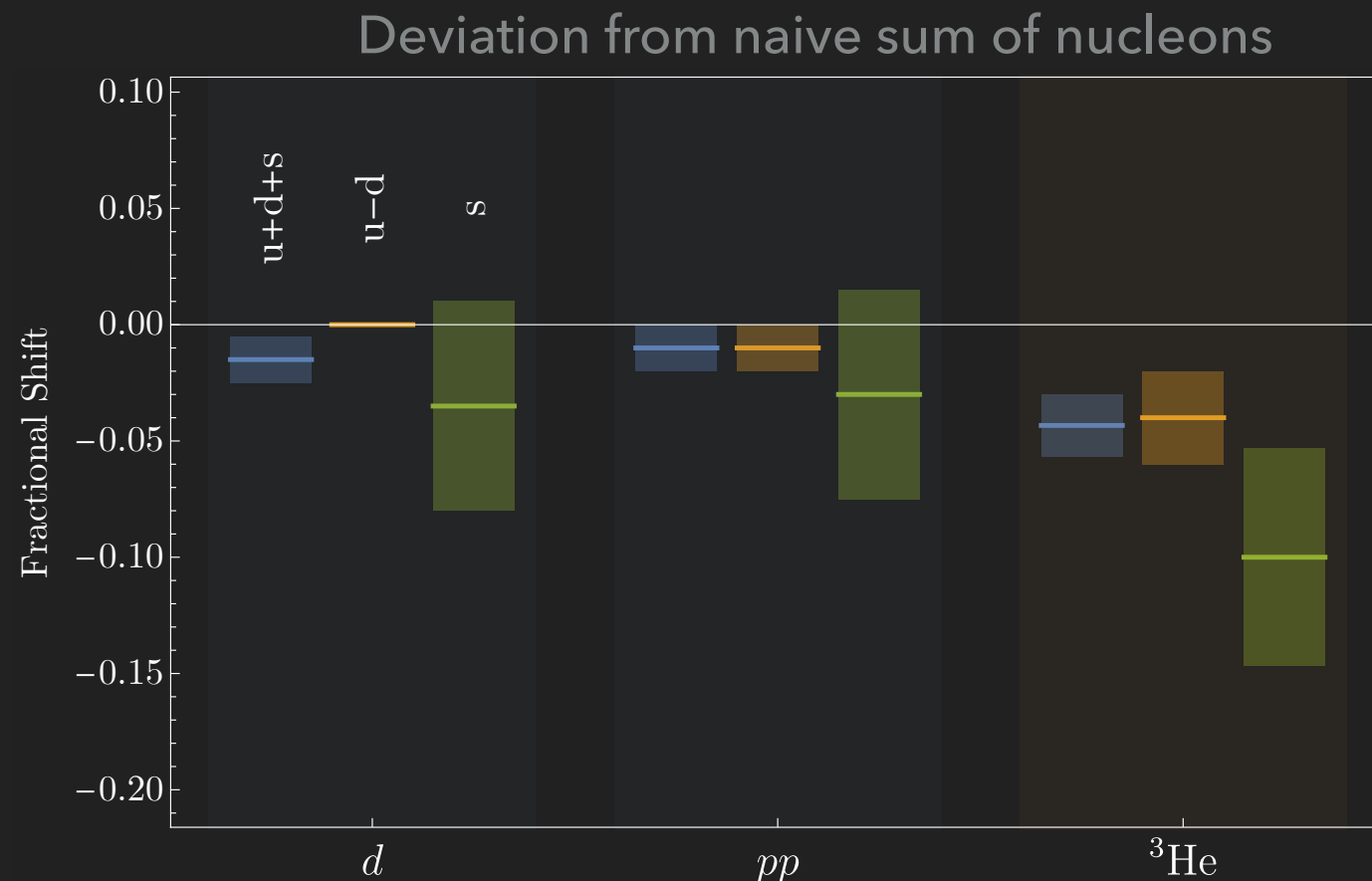
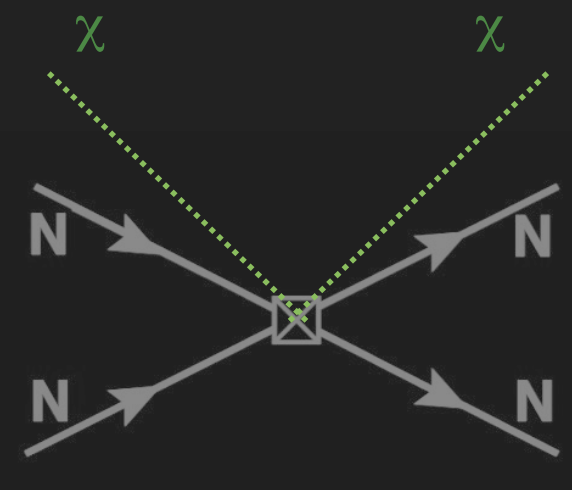
► Results from many groups

Summary from Shanahan 2016



NUCLEAR EFFECTS CAN BE LARGE!

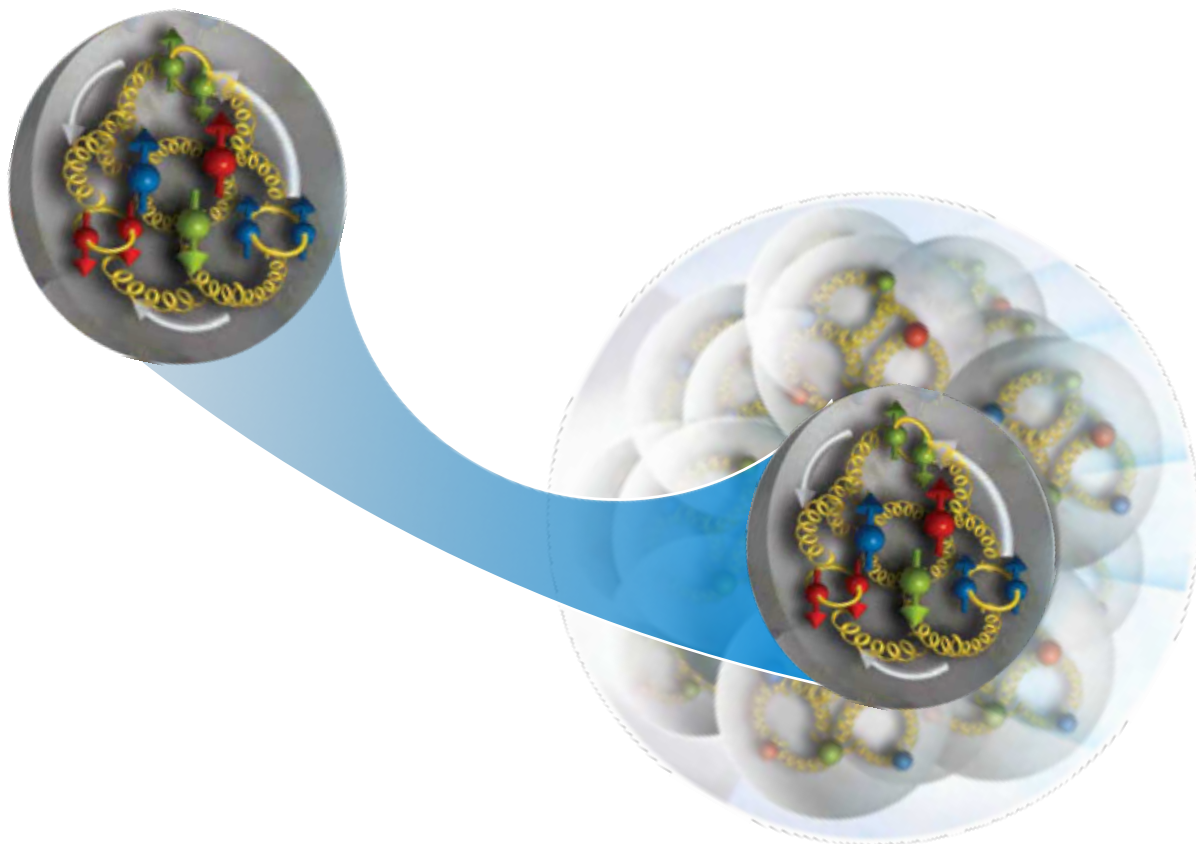
- ▶ LQCD study of scalar couplings for $A=1,2,3$
- ▶ Unexpectedly large ($\sim 10\%$) deviation from sum of nucleon matrix elements for $A=3$
- ▶ Naive extrapolation to ^{136}Xe implies significant consequences for dark matter detection sensitivity



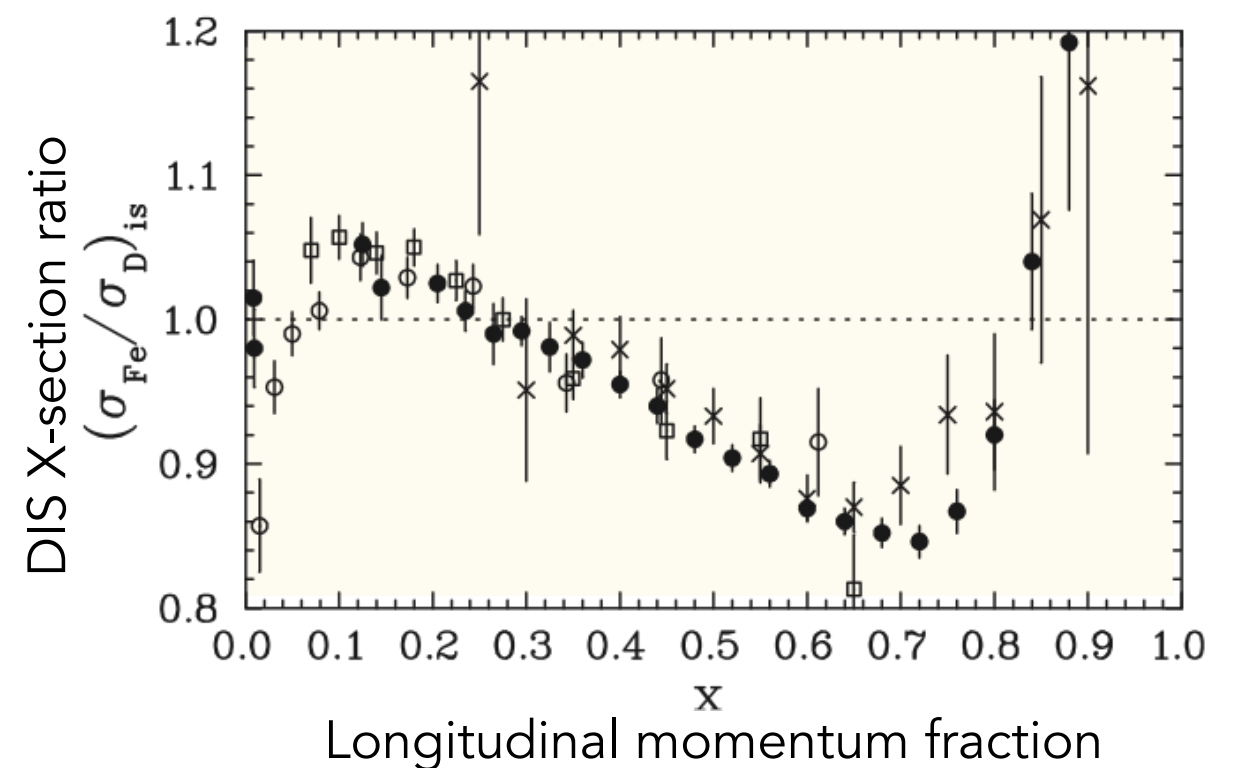
EMC-type effects from Lattice QCD

Understanding the quark and gluon structure of matter

How is the partonic structure of nucleons modified in nuclei?



Encoded in EMC-type effects

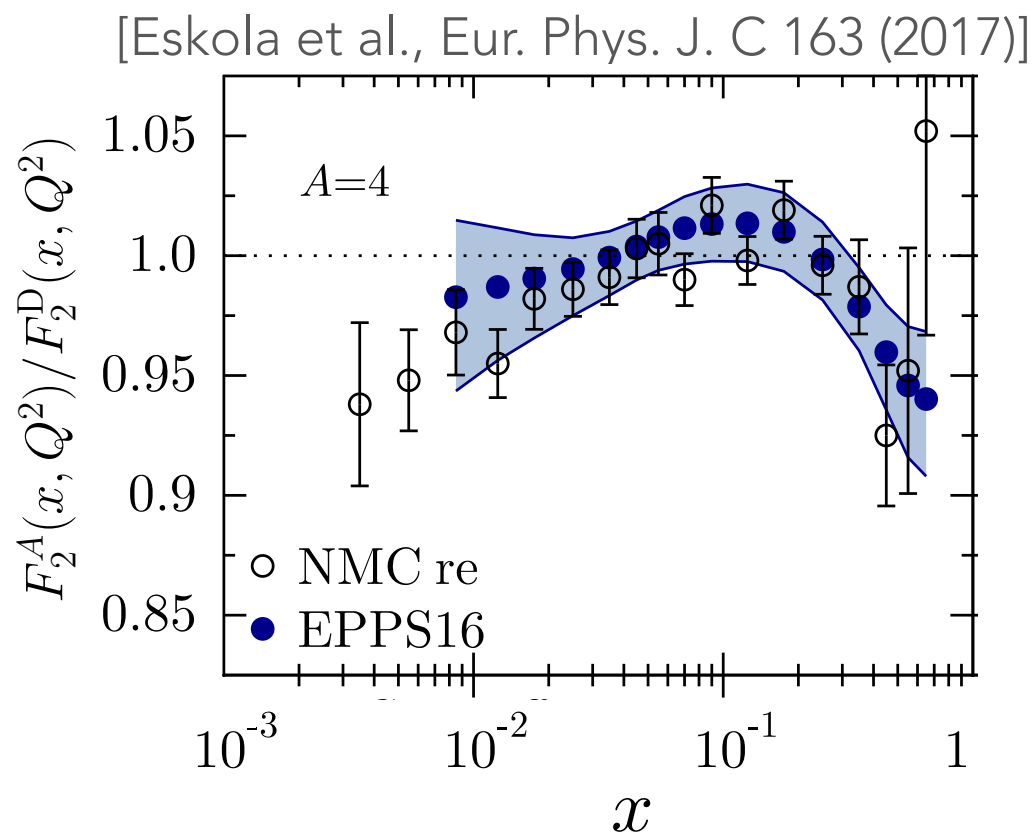


(EMC: Aubert et al., 1983)

EMC effects in Mellin moments

First investigation of EMC-type effects from LQCD:
Nuclear effects in Mellin moments of PDFs

- Calculable from local operators
- **BUT** EMC effects in moments are very small



Classic EMC effect is defined in F_2 :

$$F_2(x, Q^2) = \sum_{q=u,d,s,\dots} x e_q^2 [q(x, Q^2) + \bar{q}(x, Q^2)]$$

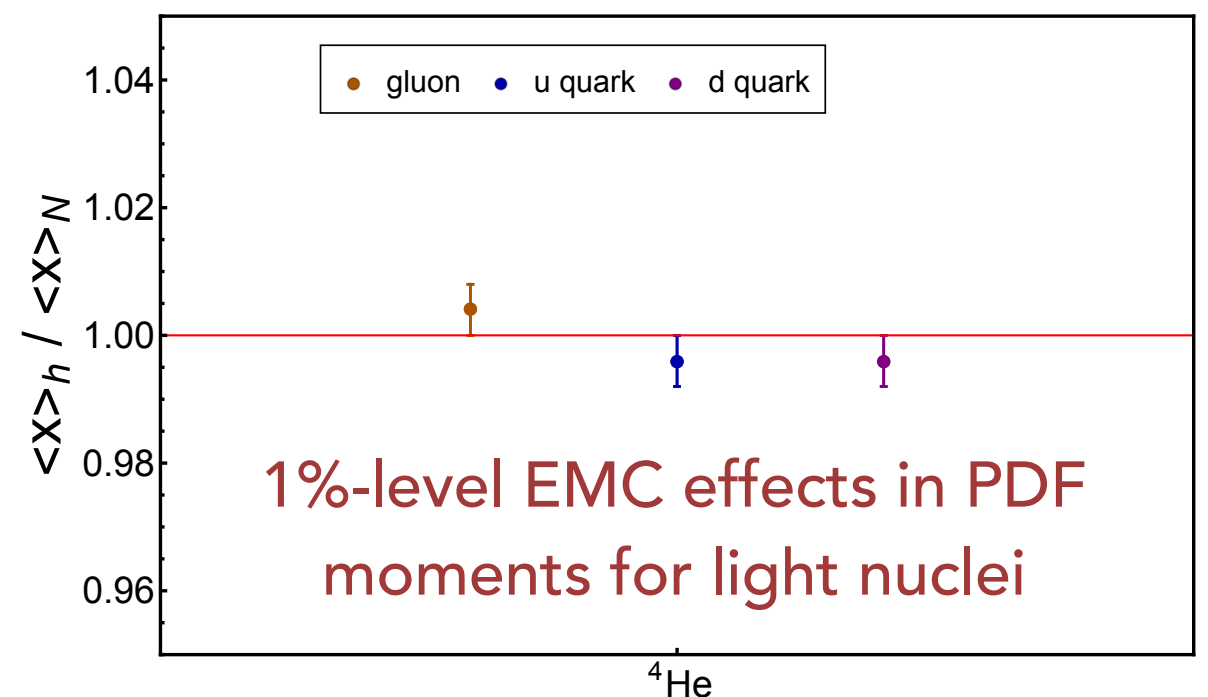
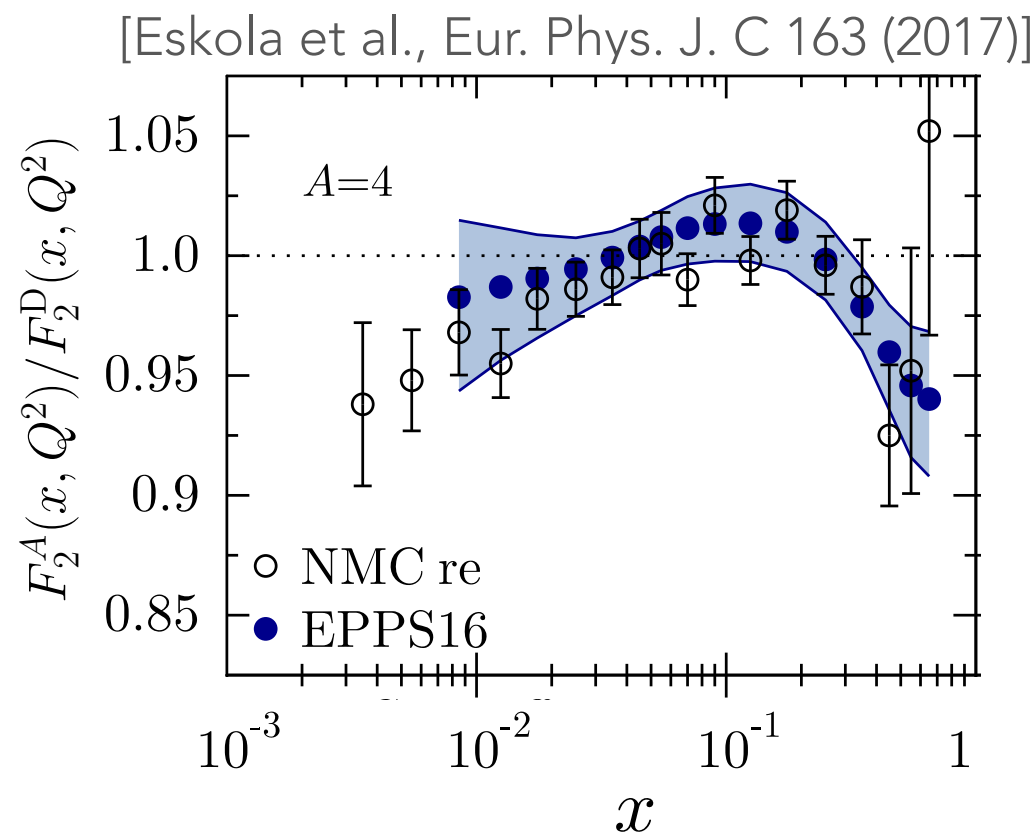
Number density of partons of flavour q

→ x-integrals of numerator and denominator $\int_0^1 dx x^n q(x, Q^2)$

EMC effects in Mellin moments

First investigation of EMC-type effects from LQCD:
Nuclear effects in Mellin moments of PDFs

- Calculable from local operators
- **BUT** EMC effects in moments are very small



Momentum fraction of nuclei

First investigation of EMC-type effects from LQCD: Nuclear effects in Mellin moments of PDFs

- Lowest Mellin moment of spin-independent PDF defines fraction of momentum of nucleus A carried by parton of type f

$$\langle x \rangle_A^f = \int_0^1 dx x f^A(x) \qquad \sum_{f=q,g} \langle x \rangle_h^f = 1$$

- Momentum sum rule implies **nucleus-independent ratio of quark and gluon EMC effects** in the first moment

$$\left(\frac{\langle x \rangle_A^f}{\langle x \rangle_p^f} - 1 \right) = E_A^f$$

$$\frac{E_A^g}{E_A^q} = - \frac{\langle x \rangle_p^q}{\langle x \rangle_p^g} \approx -1.4$$

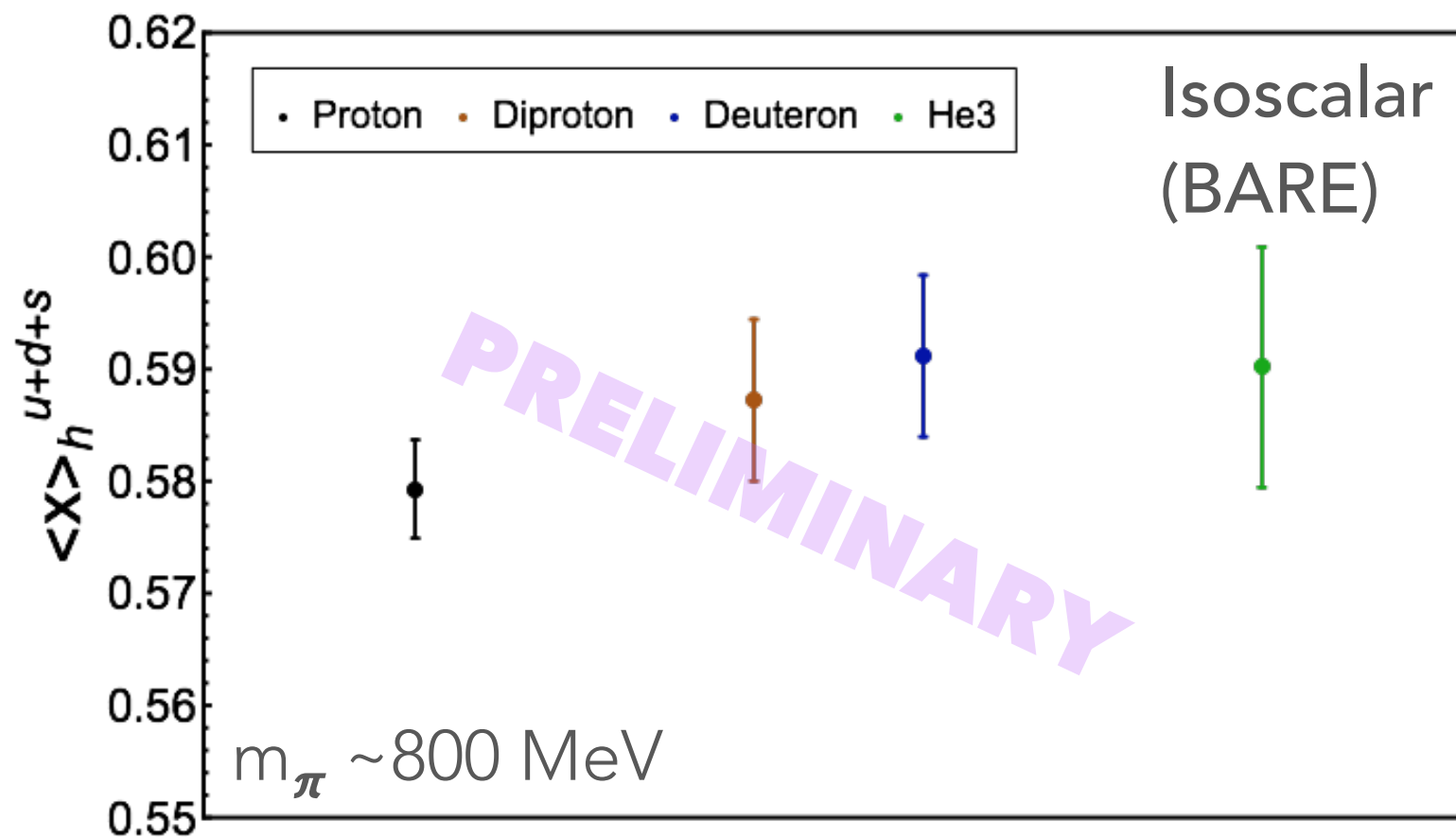
$\overline{\text{MS}} (\mu = 2\text{GeV})$

Momentum fraction of nuclei

Matrix elements of the Energy-Momentum Tensor in light nuclei

→ first QCD determination of momentum fraction of nuclei

- Few-percent determination of quark momentum fraction
~10% determination of strange quark contributions

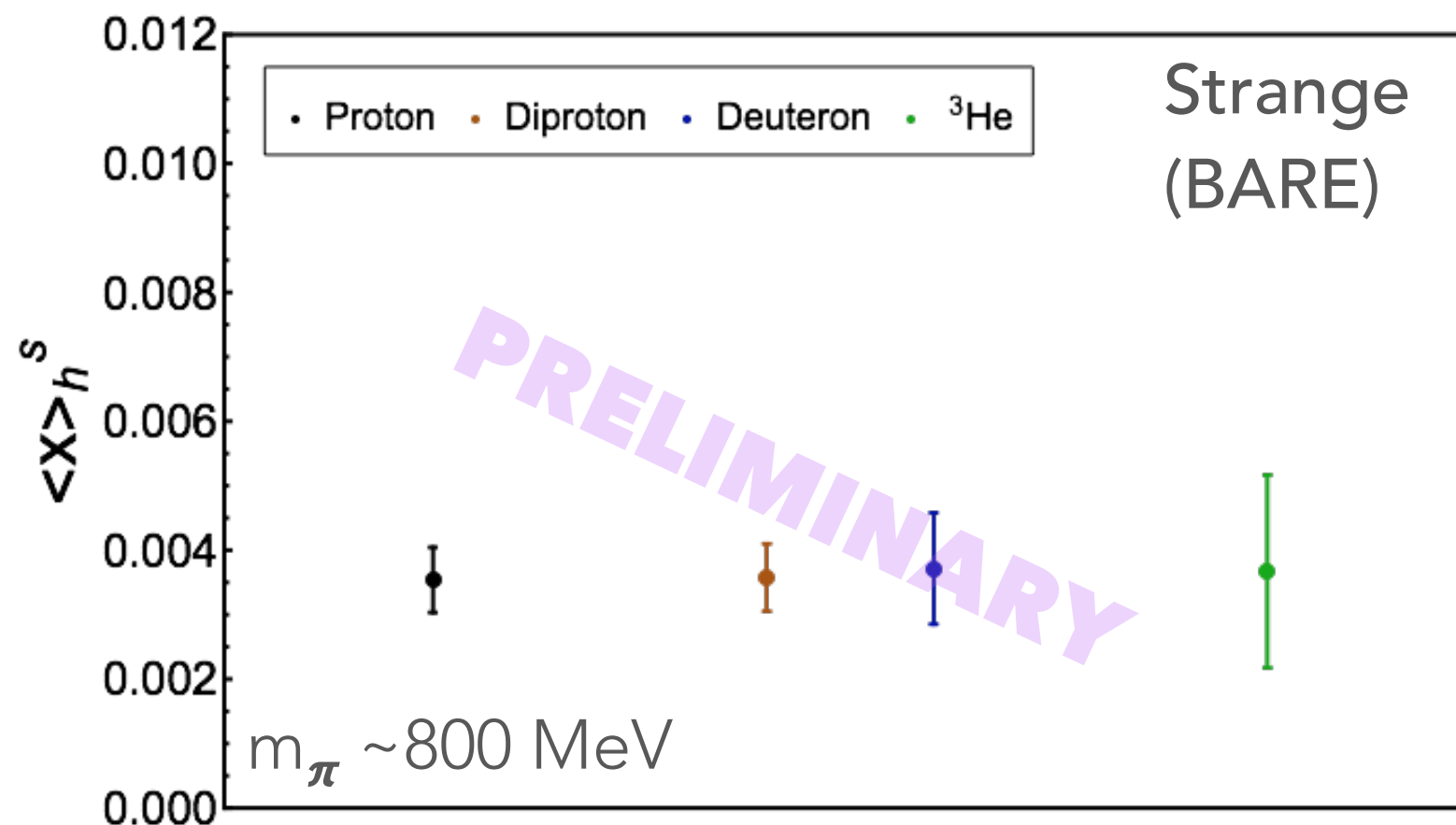


Momentum fraction of nuclei

Matrix elements of the Energy-Momentum Tensor in light nuclei

→ first QCD determination of momentum fraction of nuclei

- Few-percent determination of quark momentum fraction
~10% determination of strange quark contributions



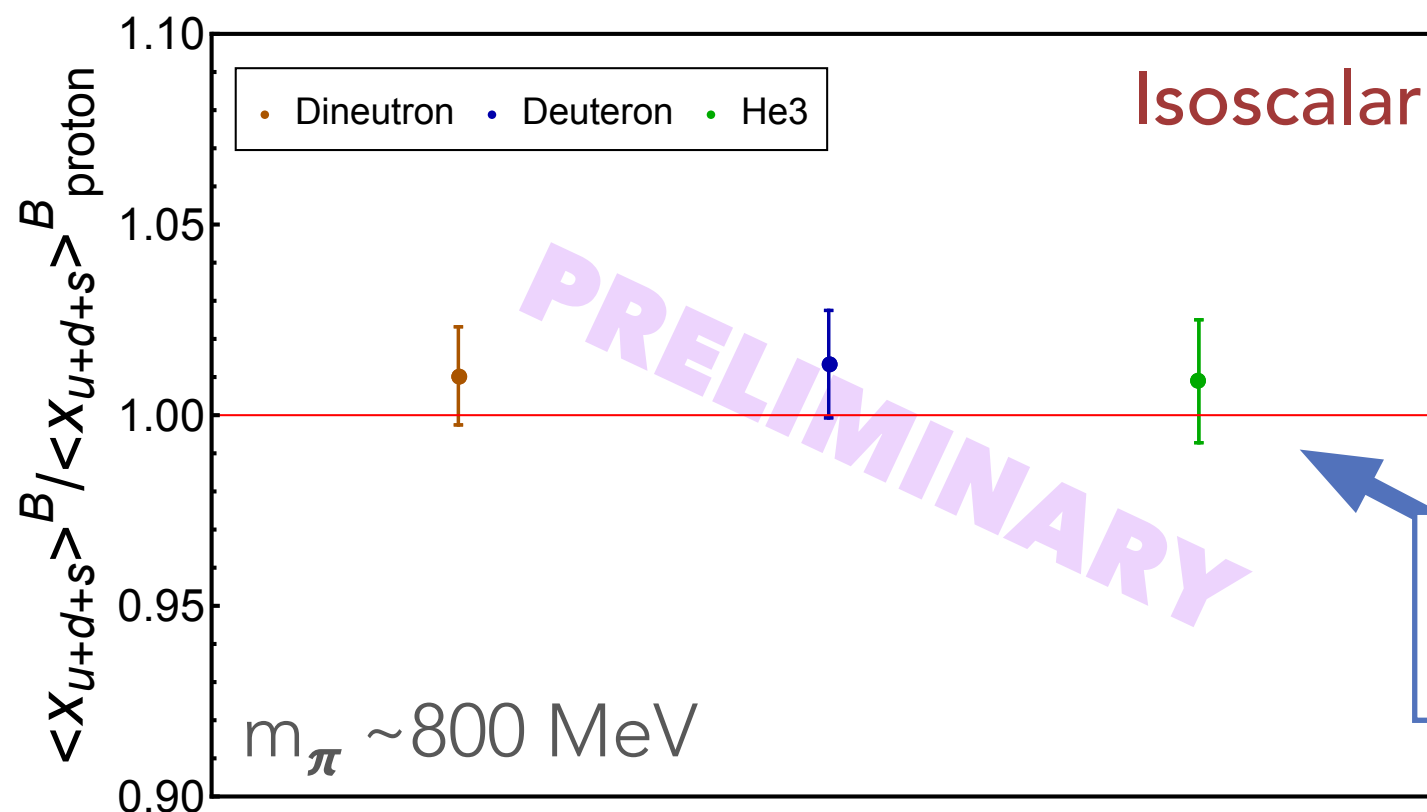
Momentum fraction of nuclei

Matrix elements of the Energy-Momentum Tensor in light nuclei

→ first QCD determination of momentum fraction of nuclei

- Bounds on EMC effect in moments at ~few percent level, consistent with phenomenology

Ratio of quark momentum fraction in nucleus to nucleon



- Small mixing with gluon EMT operators (neglected)
- **Sum rule constraint**



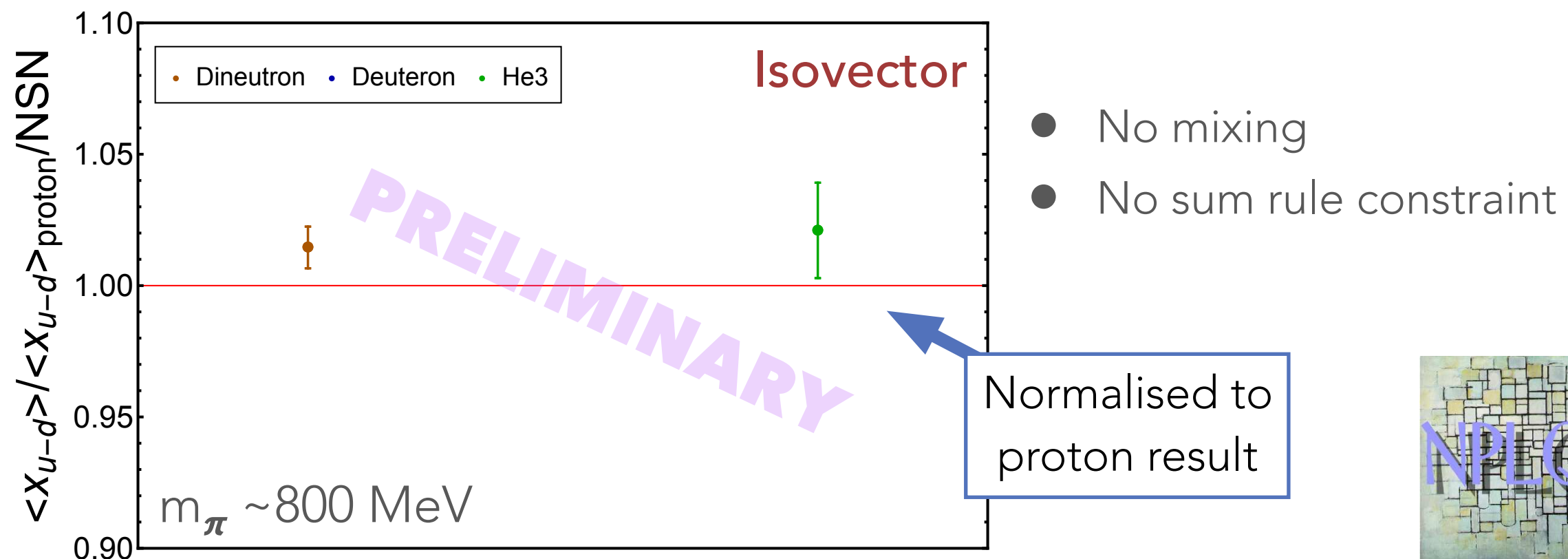
Momentum fraction of nuclei

Matrix elements of the Energy-Momentum Tensor in light nuclei

→ first QCD determination of momentum fraction of nuclei

- Bounds on EMC effect in moments at ~few percent level, consistent with phenomenology

Ratio of quark momentum fraction in nucleus to nucleon



Gluon momentum fraction of nuclei

Matrix elements of the spin-independent gluon operator in nucleon + light nuclei

→ first determination of gluon momentum fraction of nuclei

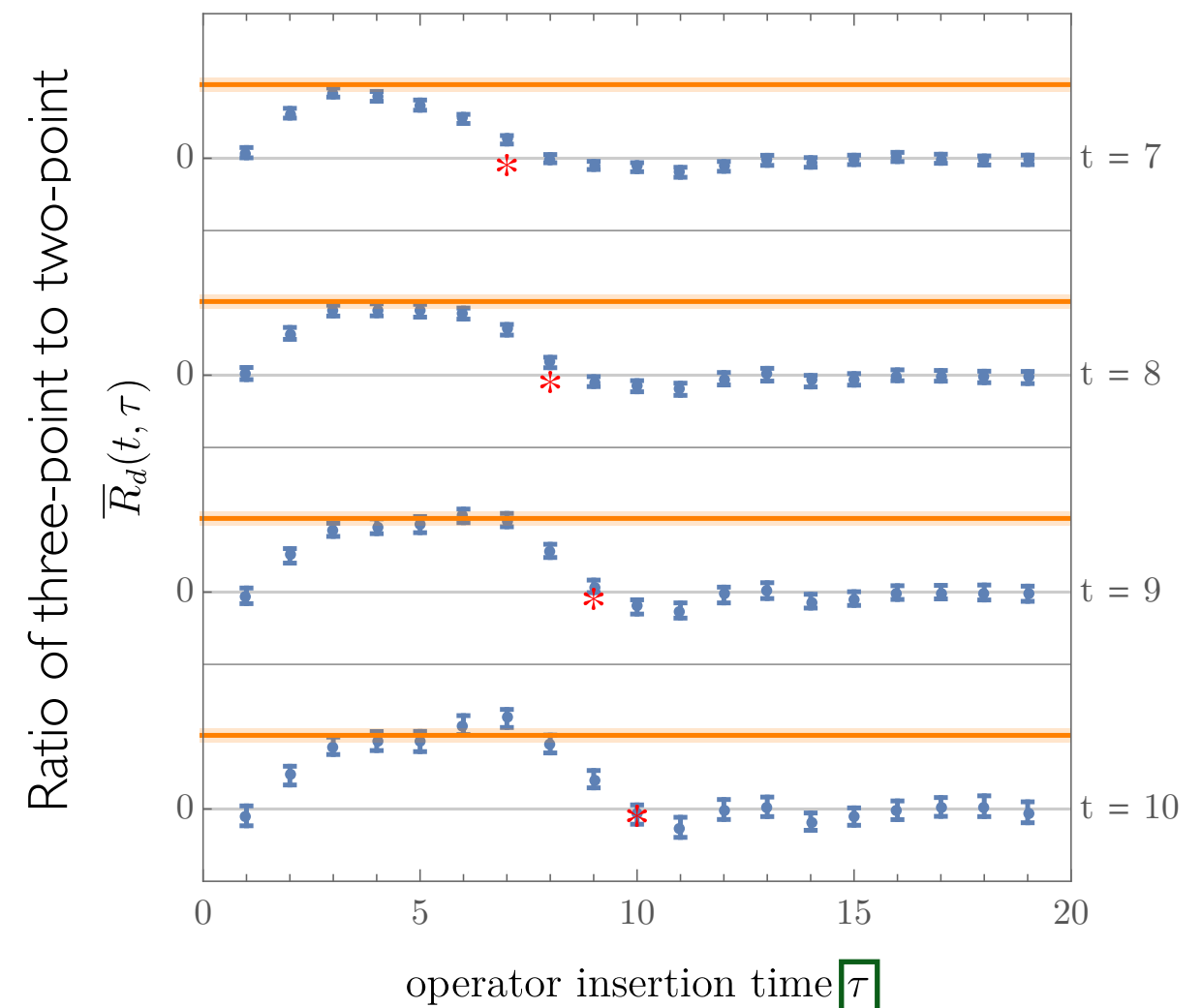
Doubly challenging:

- Nuclear matrix element
- Gluon observable (suffer from poor signal-to-noise)
- BUT: clean signals at $\sim 5\%$ precision



[NPLQCD PRD96 094512 (2017)]

Deuteron gluon momentum fraction

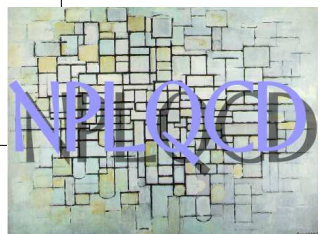
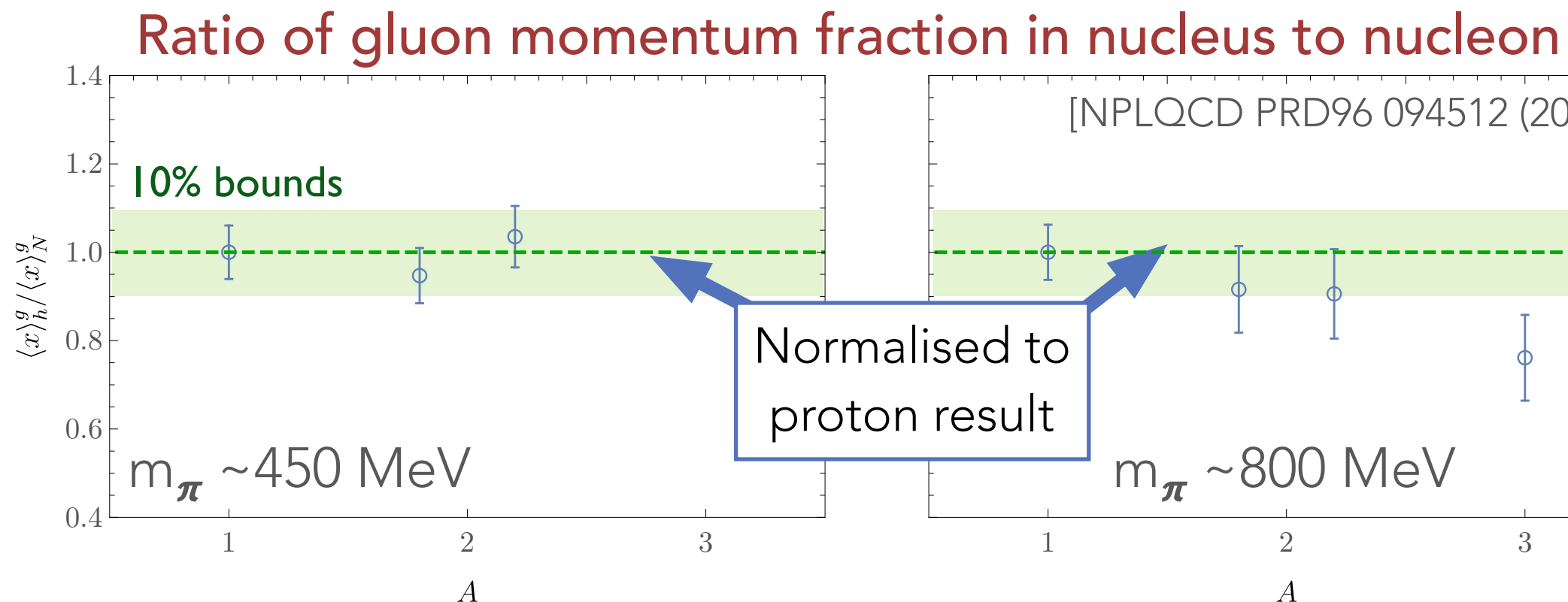


Gluon momentum fraction of nuclei

Matrix elements of the spin-independent gluon operator in nucleon + light nuclei [NPLQCD PRD96 094512 (2017)]

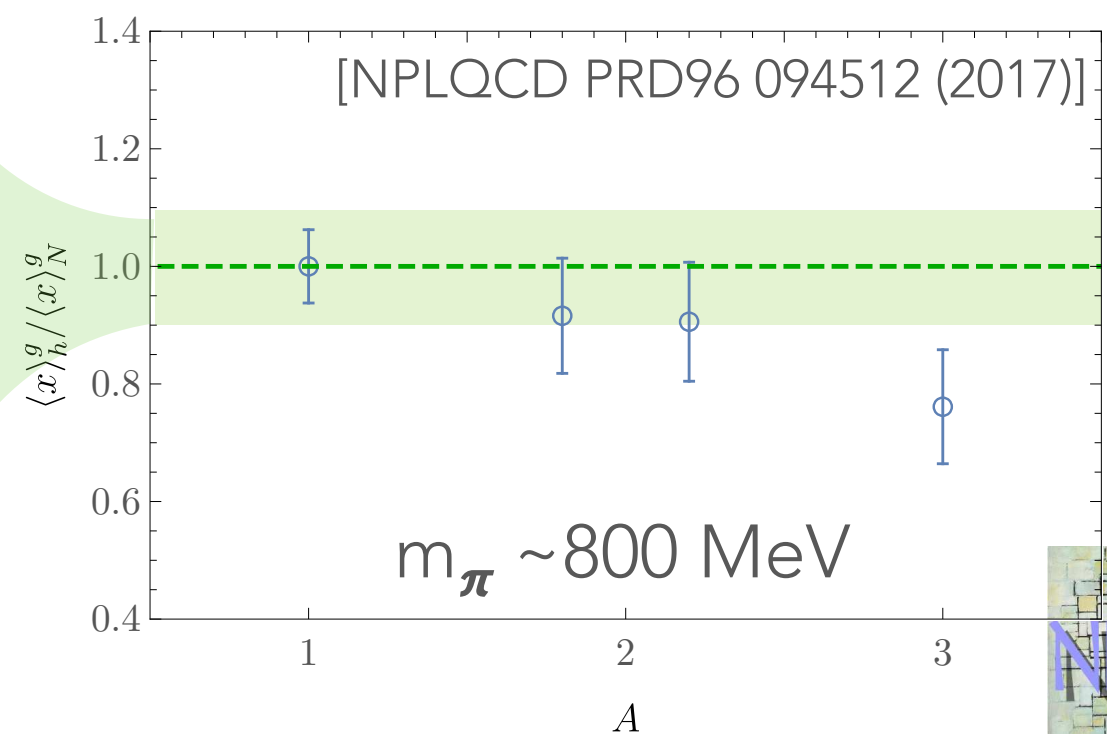
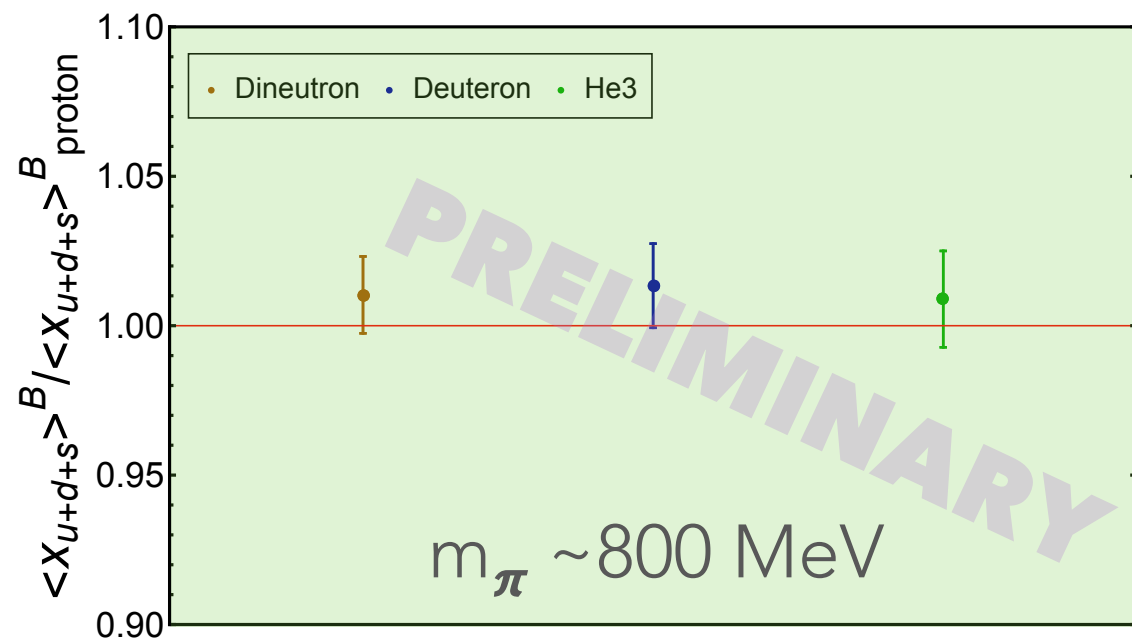
➡ first determination of gluon momentum fraction of nuclei

- Constraints at $\sim 10\%$ level on EMC-effect in gluon momentum fraction
- Small mixing with quark EMT operators (neglected)
- **Sum rule constraint**



Momentum fractions of nuclei

- First determination of all components of momentum decomposition of light nuclei
- Small mixing between quark and gluon EMT operators neglected
- Constraint on either quark or gluon EMC in this quantity implies constraint on the other from sum rules:



Exotic glue in the deuteron

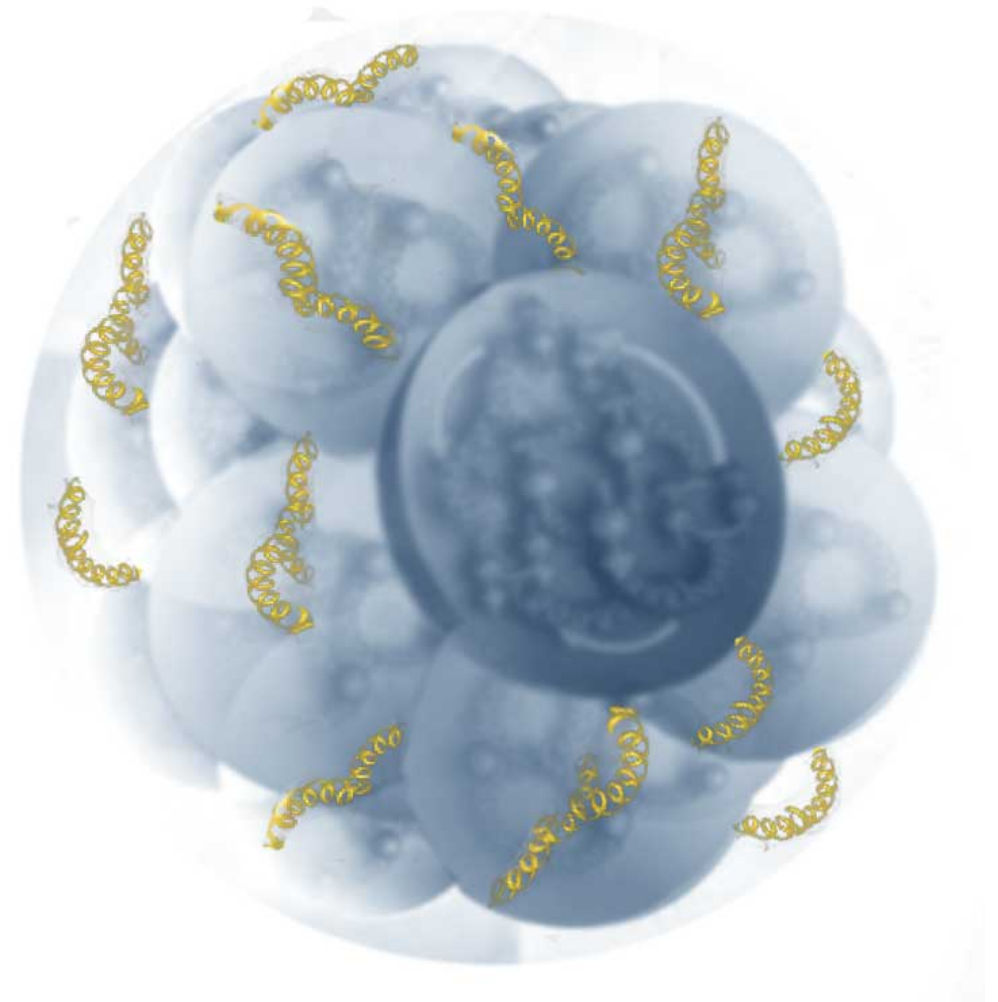
a “pure” EMC-type effect

Contributions to nuclear structure from gluons not associated with individual nucleons in nucleus

Exotic glue operator:

nucleon $\langle p | \mathcal{O} | p \rangle = 0$

nucleus $\langle N, Z | \mathcal{O} | N, Z \rangle \neq 0$

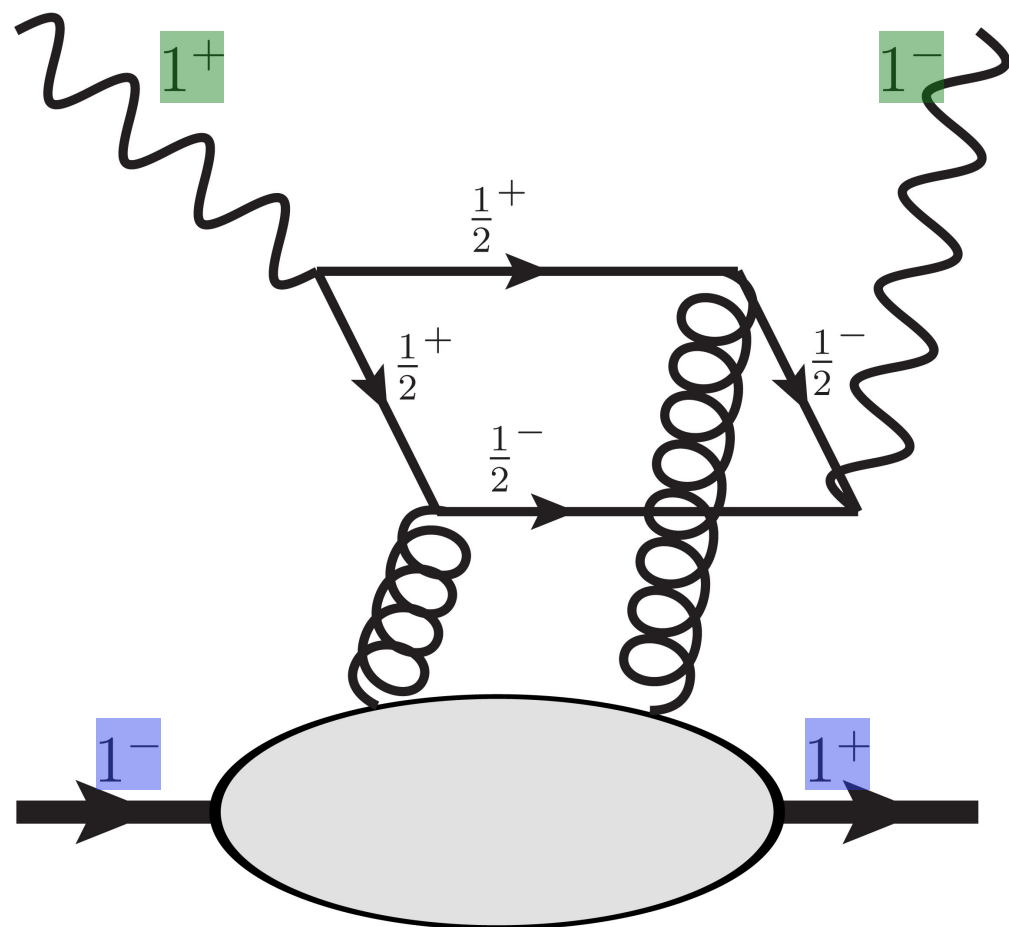


Jaffe and Manohar, “Nuclear Gluonometry”
Phys. Lett. B223 (1989) 218

Exotic glue in the deuteron

a "pure" EMC-type effect

Double helicity flip structure function $\Delta(x, Q^2)$:
changes both photon and target helicity by 2 units

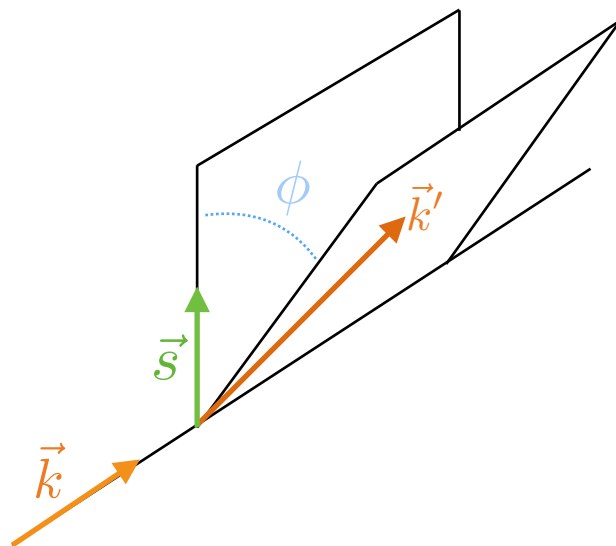


- **Unambiguously gluonic:** no analogous quark PDF at twist-2
- Non-vanishing in forward limit for targets with $\text{spin} \geq 1$
- **Experimentally measurable** in unpolarised electron DIS on polarised target
 - Nitrogen target: JLab Lol 2015
 - Polarised nuclei at EIC
- Moments calculable in LQCD

Exotic glue in the deuteron

a “pure” EMC-type effect

Double helicity flip structure function $\Delta(x, Q^2)$:
changes both photon and target helicity by 2 units



Measure azimuthal variation

$$\lim_{Q^2 \rightarrow \infty} \frac{d\sigma}{dx dy d\phi} = \frac{e^4 ME}{4\pi^2 Q^4} \left[xy^2 F_1(x, Q^2) + (1-y) F_2(x, Q^2) - \frac{x(1-y)}{2} \Delta(x, Q^2) \cos 2\phi \right]$$

- **Unambiguously gluonic:** no analogous quark PDF at twist-2
- Non-vanishing in forward limit for targets with $\text{spin} \geq 1$
- **Experimentally measurable** in unpolarised electron DIS on polarised target
 - Nitrogen target: JLab Lol 2015
 - Polarised nuclei at EIC
- **Moments calculable in LQCD**

Exotic glue in the deuteron

a “pure” EMC-type effect

Double helicity flip structure function $\Delta(x, Q^2)$:
changes both photon and target helicity by 2 units

Parton model interpretation: gluonic transversity

$$\Delta(x, Q^2) = -\frac{\alpha_s(Q^2)}{2\pi} \text{Tr} Q^2 x^2 \int_x^1 \frac{dy}{y^3} [g_{\hat{x}}(y, Q^2) - g_{\hat{y}}(x, Q^2)]$$

$g_{\hat{x}, \hat{y}}(y, Q^2)$: probability of finding a gluon with momentum fraction y
linearly polarised in \hat{x} , \hat{y} direction

Non-nucleonic glue in deuteron

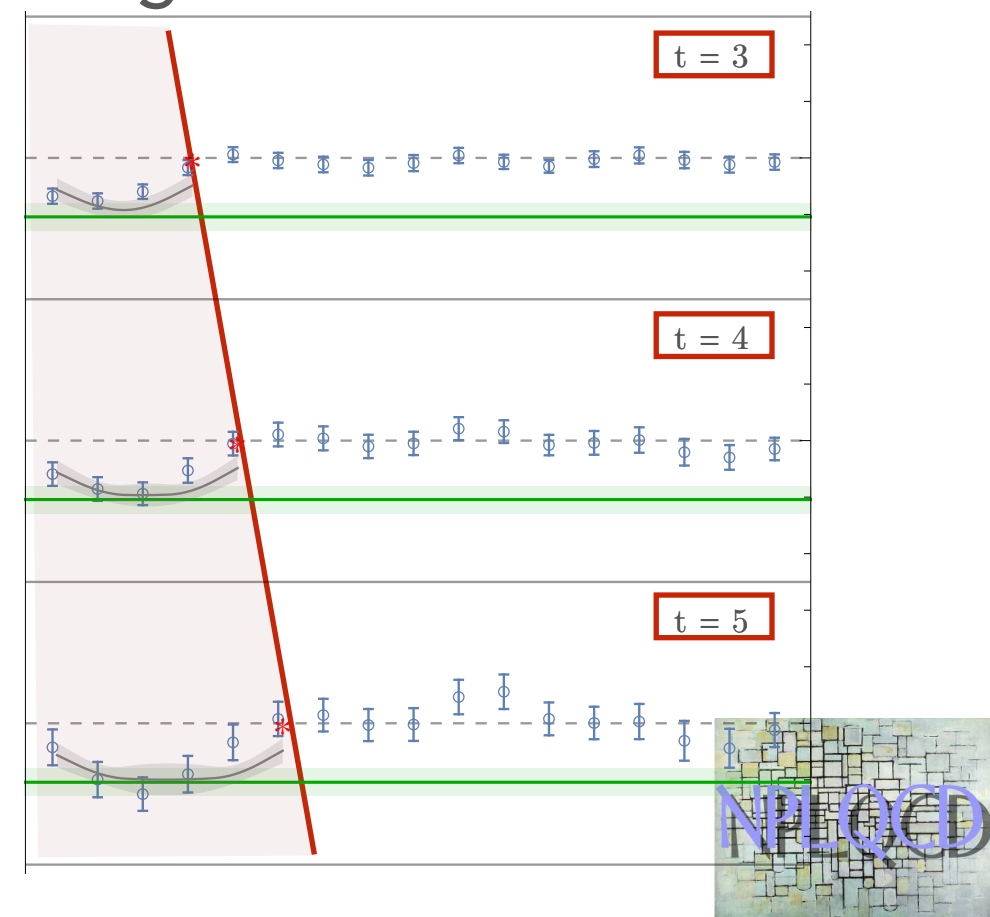
Contributions to nuclear structure from gluons not associated with individual nucleons in nucleus

$$\text{nucleon: } \langle p | \mathcal{O} | p \rangle = 0$$

$$\text{nucleus: } \langle N, Z | \mathcal{O} | N, Z \rangle \neq 0$$

- First moment of gluon transversity distribution in the deuteron
[Jaffe, Manohar PLB223 (1989) 218]
- First evidence for non-nucleonic gluon contributions to nuclear structure: LQCD with $m_\pi \sim 800$ MeV [NPLQCD PRD96 (2017)]
- Magnitude relative to momentum fraction as expected from large- N_c

Signal in LQCD data



OUTLOOK

- ▶ Nuclei are under study directly from QCD
 - ▶ Spectroscopy of light nuclei and exotic nuclei
 - ▶ Structure: magnetic moments, parton structure
 - ▶ Interactions: $np \rightarrow d\gamma$, $pp \rightarrow de + \nu$, $nn \rightarrow pp$, DM
- ▶ Prospect of a quantitative connection to QCD makes this an exciting time for nuclear physics
 - ▶ Critical role in current and upcoming intensity frontier experimental program
- ▶ Exponential improvements needed for larger nuclei: machine learning & quantum computing

