

Thermodynamical signature of quantum criticality

Universally diverging Grüneisen ratio

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Phys. Rev. Lett. 91, 066404 (2003).

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R. Küchler et al

Phys. Rev. Lett. 91, 066405 (2003).

Phys. Rev. Lett. 93, 096402 (2004).

Outline

- Brief introduction to Quantum Criticality

- Scaling Analysis of Thermodynamics

Grüneisen ratio = thermal expansion / specific heat $\rightarrow \infty$

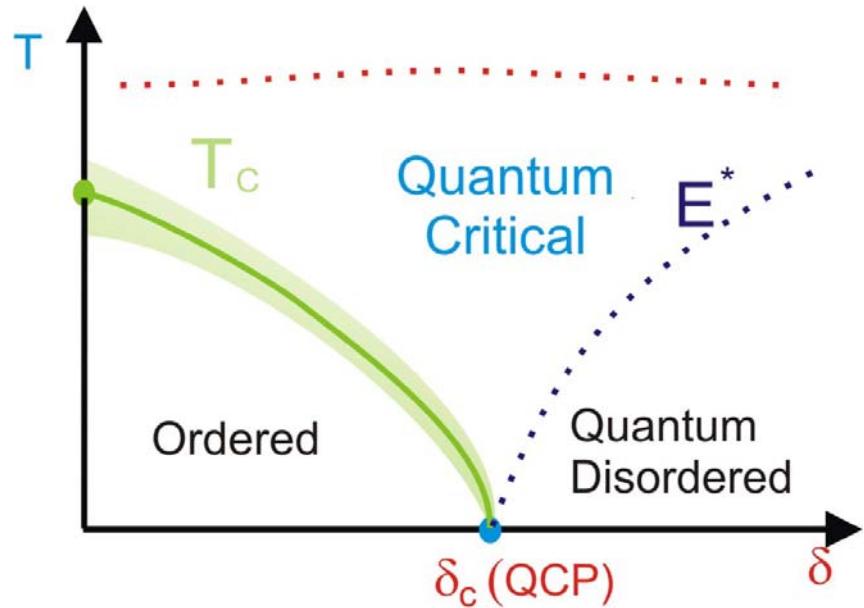
- Thermodynamic probe of quantum critical metals

Experiments

Theory

- Spin Density Wave (SDW)
- Local Quantum Criticality (LQCP)

Quantum phase transitions



S. Sachdev, *Quantum Phase Transitions*
(Cambridge), 1998.
S. L. Sondhi, S. M. Girvin, J. P. Carini, and D.
Shahar, Rev. Mod. Phys. 69, 315 (1997).

$$\begin{aligned}\xi &\sim (\delta - \delta_c)^{-\nu} \\ E^* &\sim (\delta - \delta_c)^{\nu z} \rightarrow 0\end{aligned}$$

Quantum fluctuations

$$\tau_Q \sim \hbar/E^*$$

Thermal fluctuations

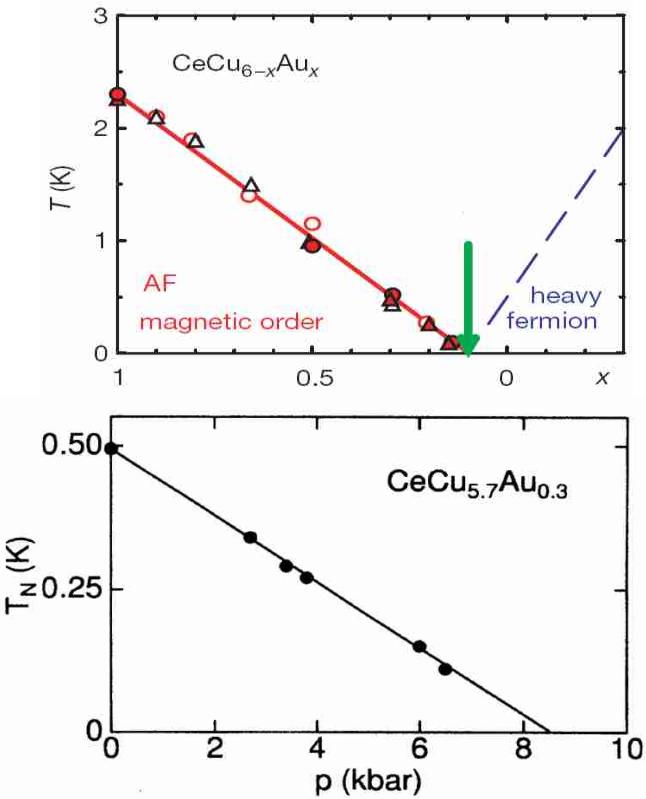
$$\tau_T \sim \hbar/(k_B T)$$

Quantum Critical

- ★ $\tau_Q \gg \tau_T$
- ★ Non-Fermi Liquid, ...

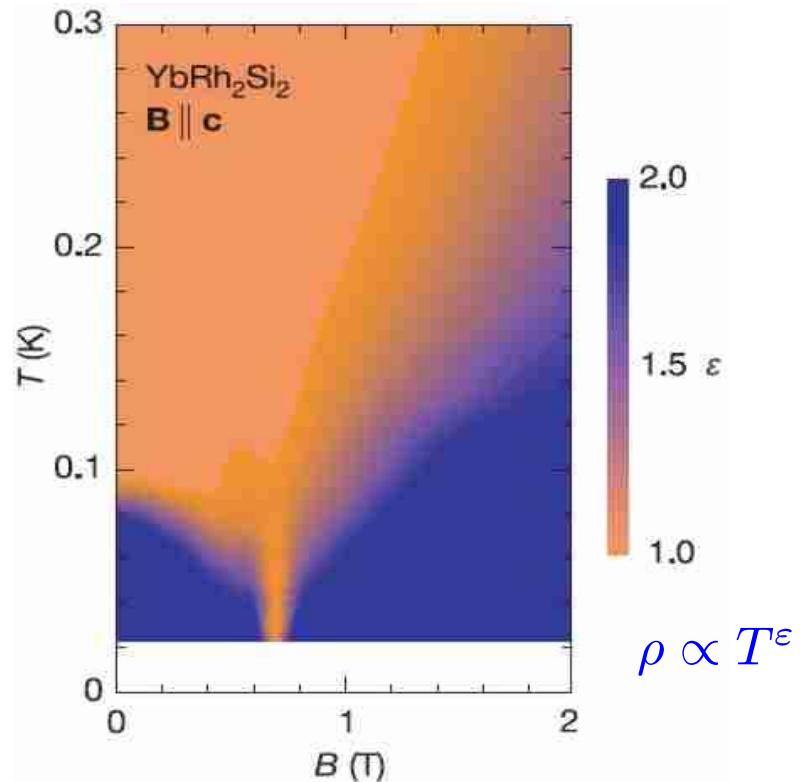
QPT in heavy fermion metals

$\text{CeCu}_{6-x}\text{Au}_x$



H. v. Löhneysen et al, Phys. Rev. Lett. **72**, 3262 (1994). B. Bogenberger and H. v. Löhneysen, *ibid*, **74**, 1016 (1995).
A. Schröder et al, Nature **407**, 351 (2000).

YbRh_2Si_2



O. Trovarelli et al., Phys. Rev. Lett. **85**, 626 (2000).
J. Clusters et al, Nature **424**, 524 (2003) .

How to understand QCPs?

Quantum Criticality

New Physics: Non-Fermi liquid,
Unconventional superconductivity? ...

Nature of QCP

$d+z$ Ginzburg-Landau Theory ?

Experimental probes

thermodynamics

Grüneisen ratio

$$\Gamma = \frac{\alpha}{c_p}$$

Thermal expansion
Specific heat

Why Grüneisen ratio particularly revealing to QCPs?

Capture two relevant directions

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_{p,N} = -\frac{1}{V} \left(\frac{\partial S}{\partial p} \right)_{T,N}$$

$$c_p = \frac{T}{N} \left(\frac{\partial S}{\partial T} \right)_p$$

Classical phase transition

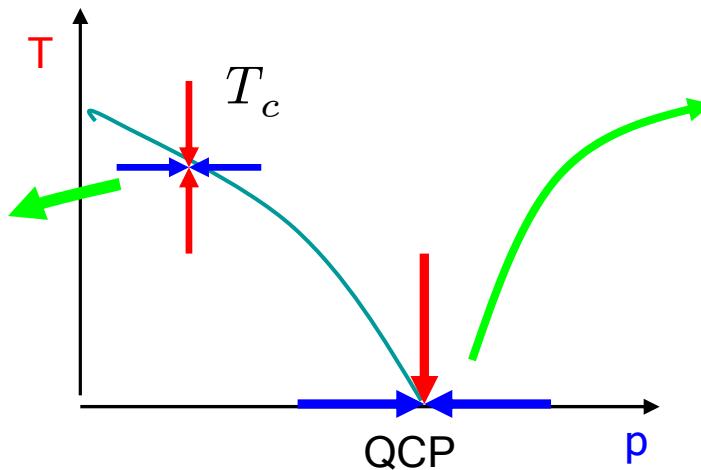
$$t = \frac{T-T_c}{T_c}$$

$$\xi \sim |t|^{-\nu_c}$$

$$c_p \sim |t|^{-(2-d\nu_c)}$$

$$\sim \alpha$$

$$\rightarrow \Gamma \sim const$$



Quantum phase transition

- T, p not equivalent !

- Capture both relevant directions

- $c_p \rightarrow 0$

- α is more singular than c_p

- divergent Γ

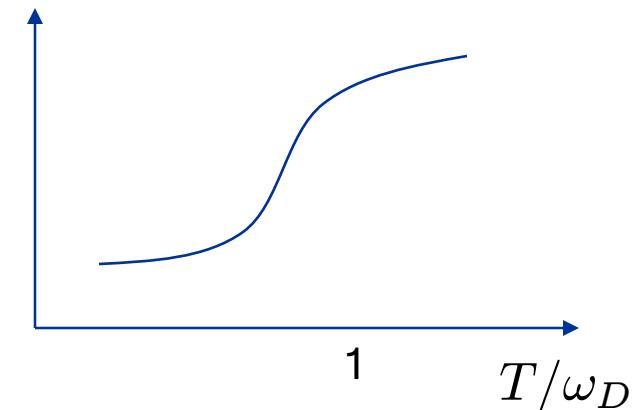
Why Grüneisen ratio particularly revealing to QCPs?

Grüneisen parameter in solids

- Phonons: Debye Frequency ω_D
- Electrons: Fermi Energy E_F

Characteristic Energy E^*

$$\begin{aligned} S/N &= f(T/E^*) \\ \Gamma &= \frac{1}{V_m E^*} \left(\frac{\partial E^*}{\partial p} \right). \end{aligned}$$



Finite!

Classical phase transition T_c

Quantum critical point

$$E^* \rightarrow 0 \Rightarrow \Gamma \rightarrow \infty$$

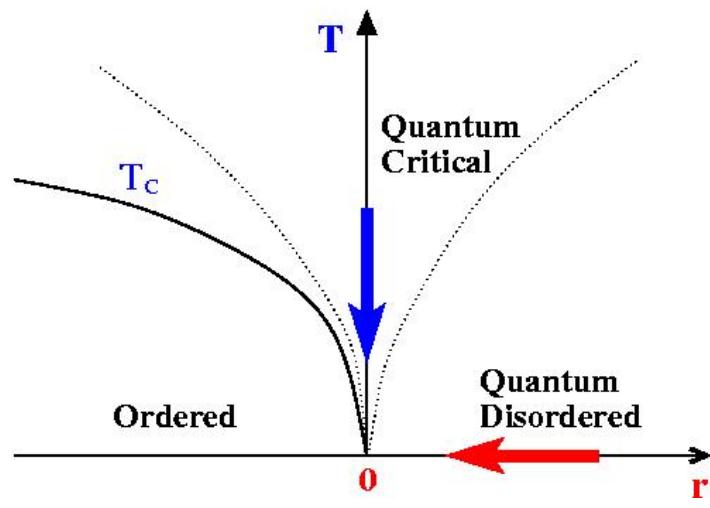
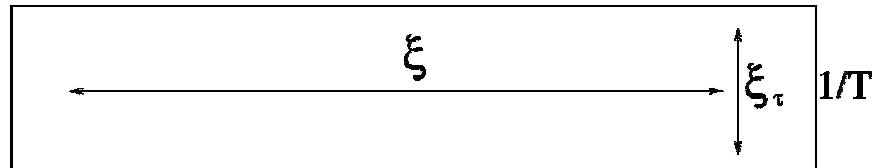
Divergent!

Quantum scaling (Hyperscaling)

Finite size scaling

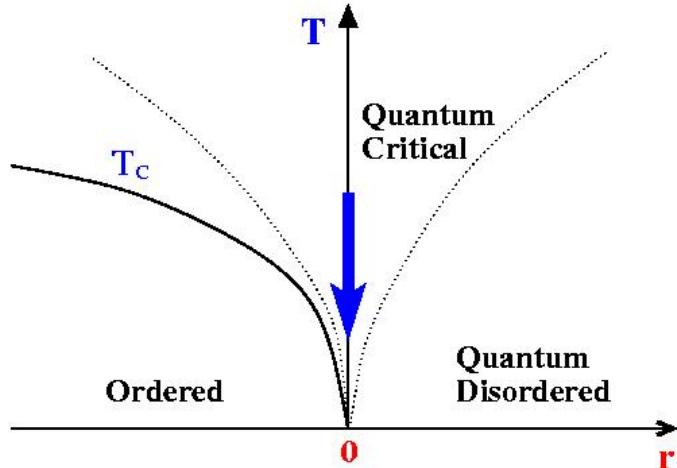
$$\xi \sim |r|^{-\nu}, \quad \xi_\tau \sim \xi^z \sim 1/T$$

d -dim Quantum System $\Leftrightarrow d+z$ Classical System



$$\begin{aligned} \frac{F_{cr}}{V} &= -\frac{k_B T}{V} \ln Z \\ &\sim \boxed{\xi^{-(d+z)}} f_1\left(\frac{1/T}{\xi^z}\right) \\ &= -\rho_0 r^{\nu(d+z)} \tilde{f}\left(\frac{T}{r^{\nu z}}\right) \\ &= -\rho_0 T^{\frac{d+z}{z}} f\left(\frac{r}{T^{1/\nu z}}\right) \end{aligned}$$

Quantum critical regime



$$r = (p - p_c)$$

$$\alpha \sim \frac{\partial^2 F}{\partial r \partial T}$$

$$F_{cr} = -\rho_0 T^{\frac{d+z}{z}} f\left(\frac{r}{T^{\frac{1}{\nu z}}}\right)$$

$$f(x \rightarrow 0) \approx f(0) + x f'(0) + \dots$$

$$c_{cr}(T, r = 0) \sim T^{d/z}$$

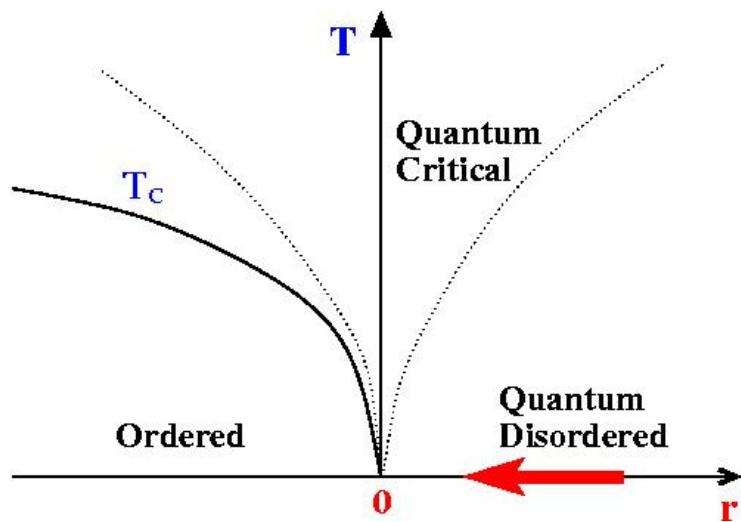
$$\alpha_{cr}(T, r = 0) \sim T^{(d - \frac{1}{\nu})/z}$$

$$\Gamma_{cr}(T, r = 0) = \frac{\alpha_{cr}}{c_{cr}} = -G_T \frac{1}{T^{1/(\nu z)}}$$

$$\Gamma \sim 1/T^x, \quad x = 1/(\nu z)$$

Measures the scaling dimension of the tuning parameter

Quantum disordered regime



$$F_{cr} = -\rho_0 r^{\nu(d+z)} \tilde{f}\left(\frac{T}{r^{\nu z}}\right)$$

$$\tilde{f}(x \rightarrow 0) = \tilde{f}(0) + c x^{y_0+1}$$

$$c_{cr}(T \rightarrow 0, r) \sim T^{y_0} r^{\nu(d-y_0z)}$$

$$\alpha_{cr}(T \rightarrow 0, r) \sim T^{y_0} \frac{r^{\nu(d-y_0z)}}{p_c r}$$

$$\Gamma_{cr}(T \rightarrow 0, r) = -G_r \frac{1}{V_m(p - p_c)}.$$

$$G_r = \boxed{\frac{\nu(d - y_0 z)}{y_0}}.$$

Results for hyperscaling

- Any generic QCP, $\Gamma \rightarrow \infty \Rightarrow$ identify a QCP
- Characterize a QCP
 - QC, $\Gamma \sim 1/T^x$, x relates the most relevant parameter
 - QD, $\Gamma = G_r/[V_m(p - p_c)]$, G_r universal

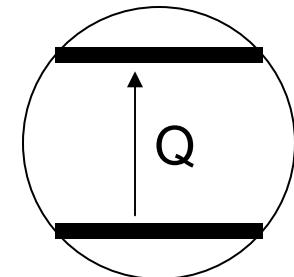
Questions:

Whether hyperscaling holds

Spin-density-wave Picture

Hertz, Millis, Moriya,

An effective d+z Ginzburg-Landau Theory



Long-wavelength fluctuations of the order parameter $\vec{\phi}$
 $d_{\text{eff}} = d + z \geq 4$

$$S[\vec{\phi}] = \frac{1}{2} \sum_{\mathbf{q}, i\omega_l} \chi_0^{-1}(\mathbf{q}, i\omega_l) |\vec{\phi}(\mathbf{q}, i\omega_l)|^2 + u \int_0^\beta d\tau \int d^d \mathbf{r} |\vec{\phi}(\mathbf{r}, \tau)|^4,$$

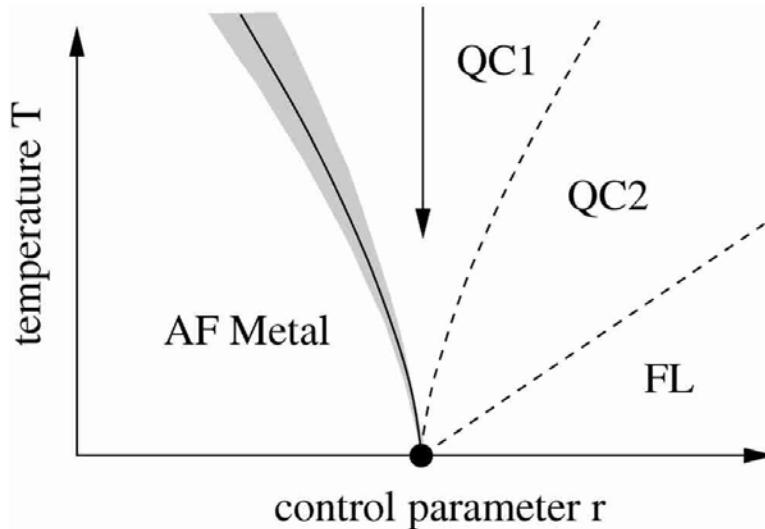
$$\chi_0^{-1}(\mathbf{q}, i\omega_l) = \delta + (q\xi_0)^2 + \frac{|\omega_l|}{\Gamma_q}$$

$$\omega \sim q^z$$

$$\text{AFM fluctuation } Q = (\pi, \dots, \pi) \quad \Gamma_q = \Gamma_0 Q \quad z = 2$$

$$\delta \sim q^{1/\nu} \quad \nu = 1/2$$

SDW - 3D Antiferromagnetic fluctuation



$$F_{cr} \sim T^{(d+z)/z} f \left(\frac{r}{T^{1/(\nu z)}}, u T^{(d+z+\frac{1}{\nu}-4)/z} \right)$$

FL regime

$$F_{cr} \sim r^{(d+z)\nu} \left(\frac{T}{r^{\nu z}} \right)^2$$

QC regime

$$F_{cr} \sim T^{(d+z)/z} \left(1 + \frac{r}{T^{1/\nu z}} + u T^{(d+z-4)/z} \right)$$

SDW - Results

FL Regime

$$d = 3, z = 2 \quad d = 2, z = 2$$

$$\alpha_{cr} \quad Tr^{-1/2} \quad Tr^{-1}$$

$$c_{cr} \sim \quad -Tr^{1/2} \quad T \log \frac{1}{r}$$

$$\Gamma_{r,cr} = \quad -(2r)^{-1} \quad \left(r \log \frac{1}{r}\right)^{-1}$$

QC Regime

$$d = 3, z = 2 \quad d = 2, z = 2$$

Hyperscaling

$$\alpha_{cr} \sim$$

$$T^{1/2}$$

$$\log \log \frac{1}{T}$$

$$T^{\left(d - \frac{1}{\nu}\right)/z}$$

$$c_{cr} \sim$$

$$-T^{3/2}$$

$$T \log \frac{1}{T}$$

$$T^{d/z}$$

$$\Gamma_{r,cr} \sim$$

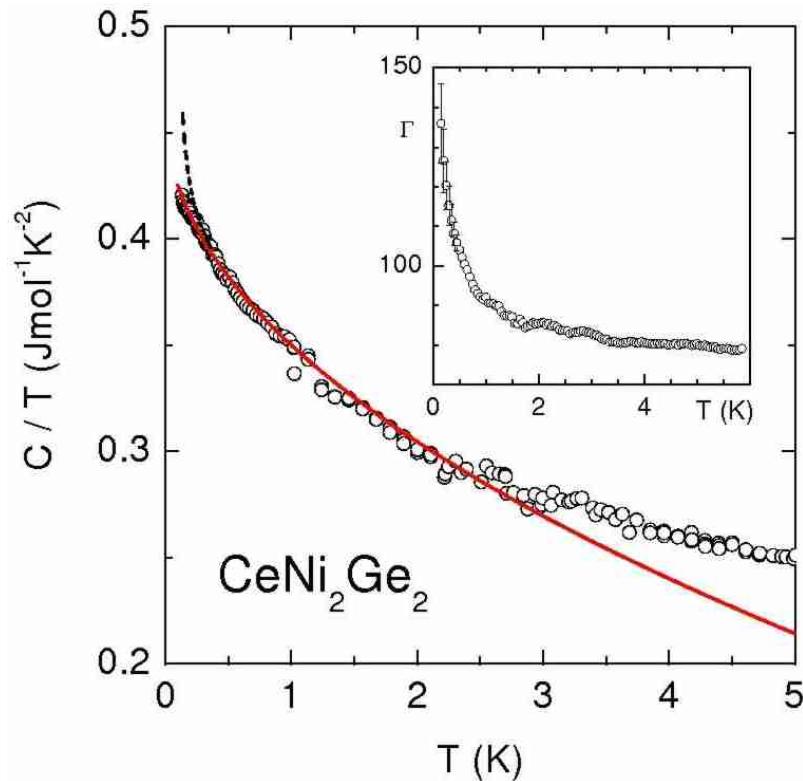
$$-T^{-1}$$

$$\frac{\log \log \frac{1}{T}}{T \log \frac{1}{T}}$$

$$T^{-1/(\nu z)}$$

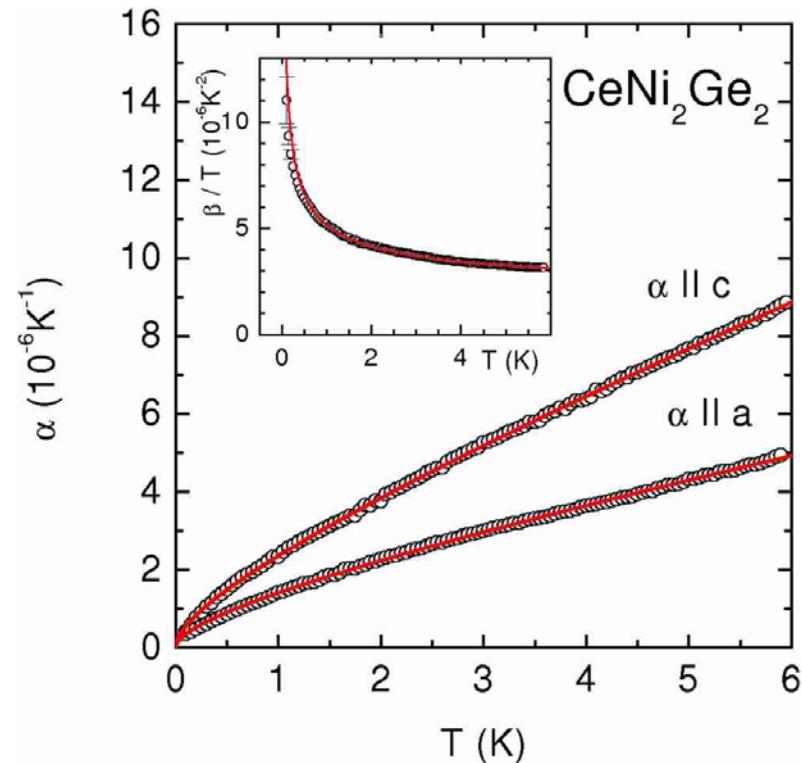
CeNi₂Ge₂

Specific heat



$$C/T = \text{const} - T^{1/2}$$

Thermal expansion

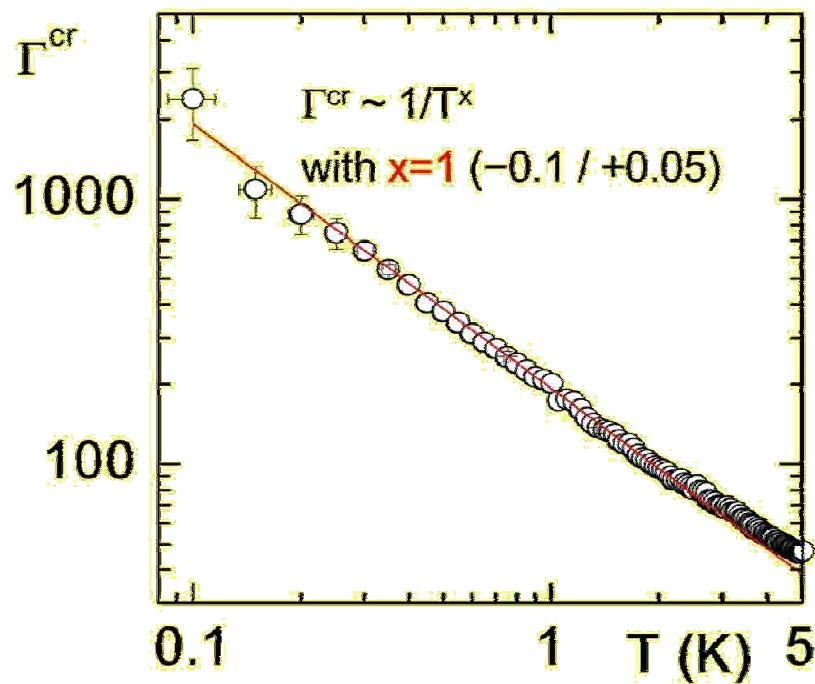


$$\alpha = aT^{1/2} + bT$$

R. Küchler, N. Oeschler, P. Gegenwart, T. Cichorek, K. Neumaier, O. Tegus, C. Geibel, J. A. Mydosh, F. Steglich, L. Zhu, and Q. Si, Phys. Rev. Lett. 91, 066405 (2003).

CeNi₂Ge₂ - *cont'd*

Grüneisen ratio



Consistent with

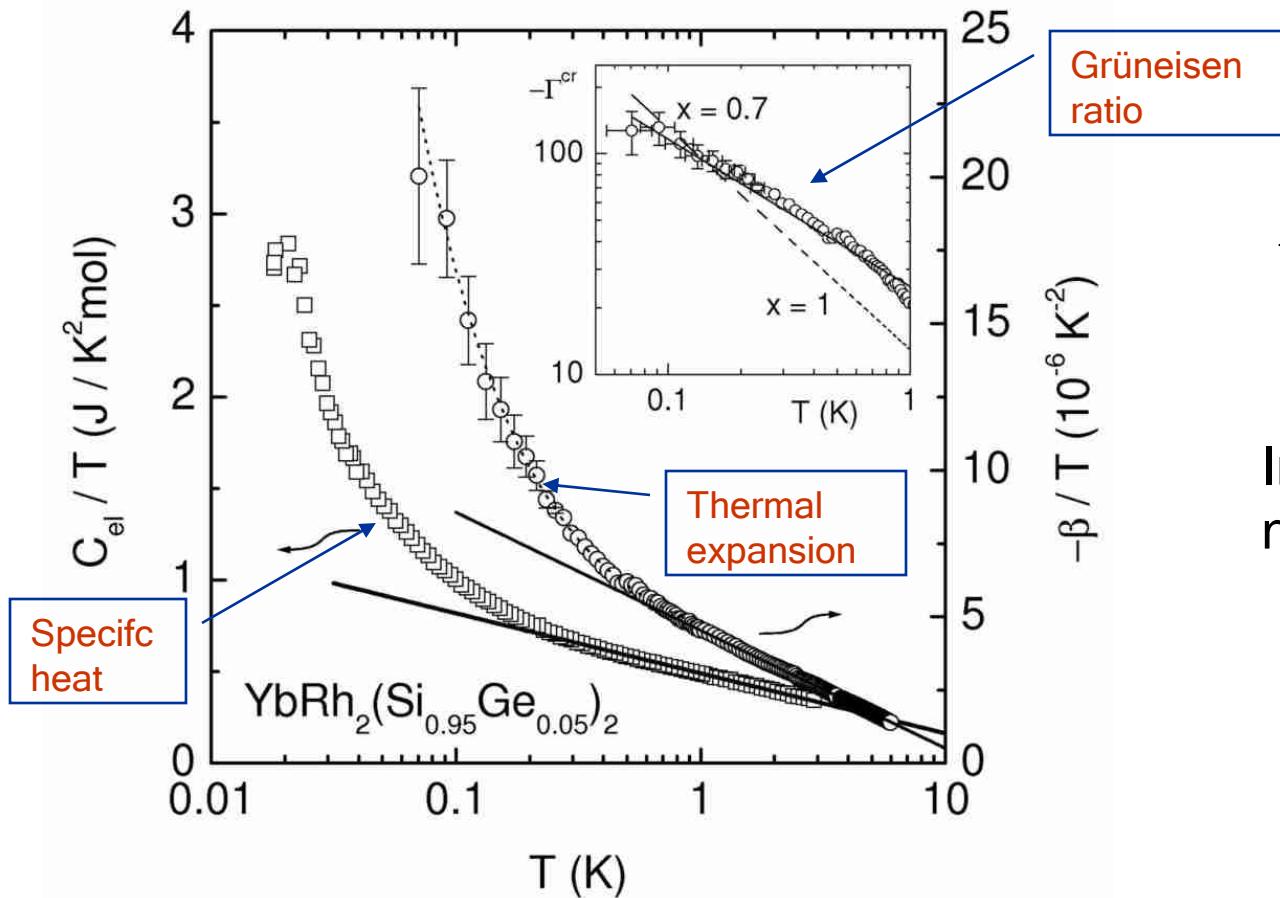
3d SDW picture

$$z = 2, \nu = 1/2$$

$$\Gamma \sim 1/T^{1/(\nu z)}$$

R. Küchler et al., Phys. Rev. Lett. 91, 066405 (2003).

$\text{YbRh}_2(\text{Si}_{0.95}\text{Ge}_{0.05})_2$



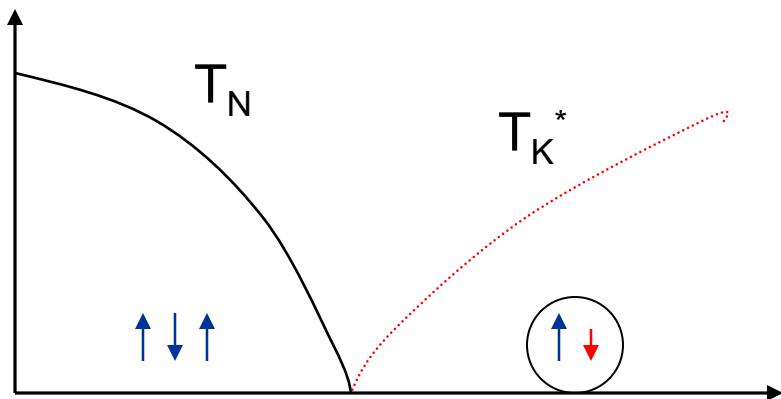
$$-\Gamma \sim 1/T^{0.7}$$

Inconsistent with
neither 3D or 2D SDW

R. Küchler et al., Phys. Rev. Lett. 91, 066405 (2003).

Theories

■ Locally quantum criticality



- critical local dynamics
- ω/T scaling, fractional exponent
- T_K^* finite for SDW transition

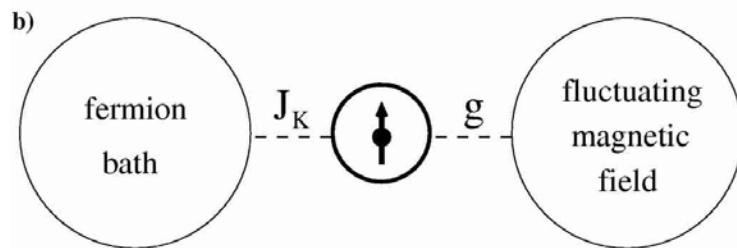
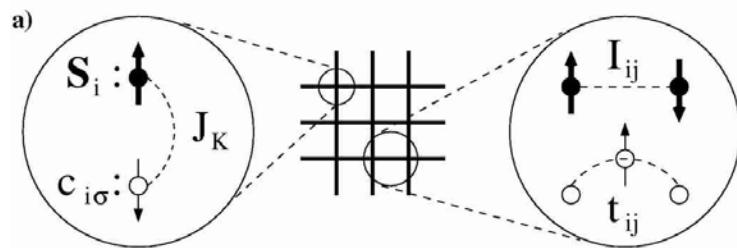
Q. Si, S. Rabello, K. Ingersent, and J.L. Smith, Nature (London) **413**, 804 (2001).
P. Coleman, C. Pépin, Q. Si, and R. Ramazashvili, J. Phys.: Condens. Matter **13**, R723 (2001).

Other theories

T.Senthil, M. Vojta, and S. Sachdev, Phys. Rev. B **69**, 035111 (2004)
C. Pépin, Phys. Rev. Lett. **94**, 066402 (2005)
V.R. Shaginyan JETP LETTERS **79**, 286 (2004)

Locally critical scenario

Self-consistent Bose-Fermi Kondo Model



$$J_K \mathbf{S} \cdot \mathbf{s}_c + \sum_{p,\sigma} E_p c_{p\sigma}^\dagger c_{p\sigma}$$

$$+ g \sum_p \mathbf{S} \cdot (\vec{\phi}_p + \vec{\phi}_{-p}^\dagger) + \sum_p w_p \vec{\phi}_p^\dagger \cdot \vec{\phi}_p.$$

$$\sum_p \delta(\omega - w_p) = (K_0^2/\pi) |\omega|^{1-\epsilon}$$

Free energy $f_{imp} \sim T\Phi(r/T^{\lambda_r})$
 Grüneisen ratio $\Gamma \sim 1/T^{\lambda_r} \leftarrow 1/(\nu z)$

anisotropic case, with $g_z, g_\perp, J_z, J_\perp$

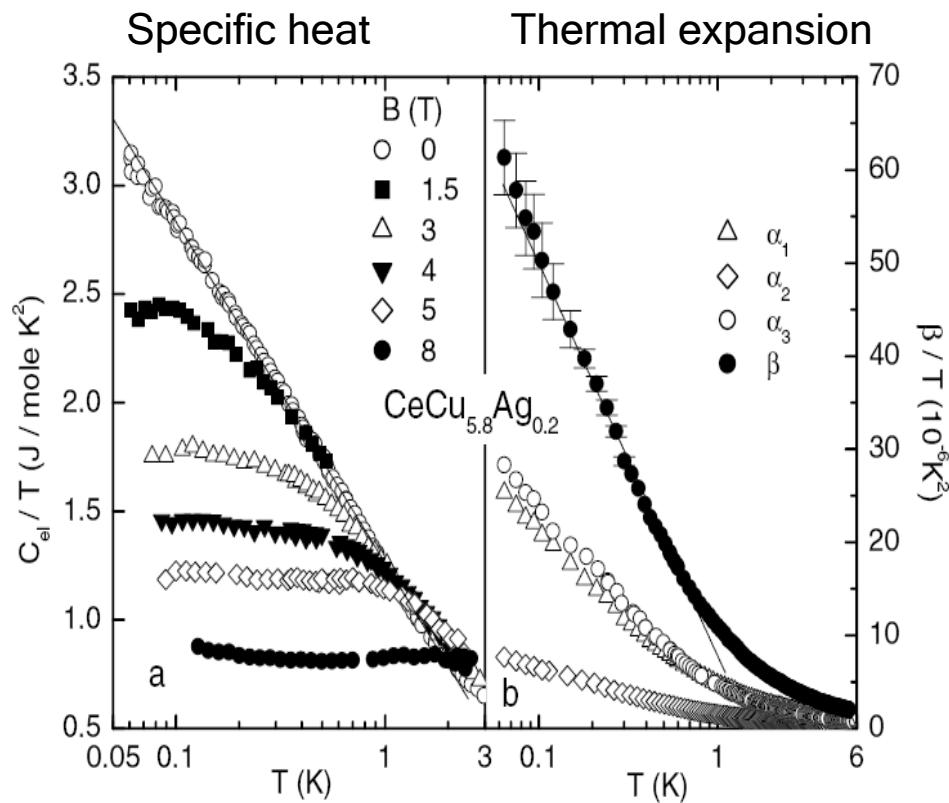
$$\text{XY} \quad g_z = 0$$

$$\lambda_r \approx 0.618034\epsilon + 0.0376644\epsilon^2 \approx 0.66$$

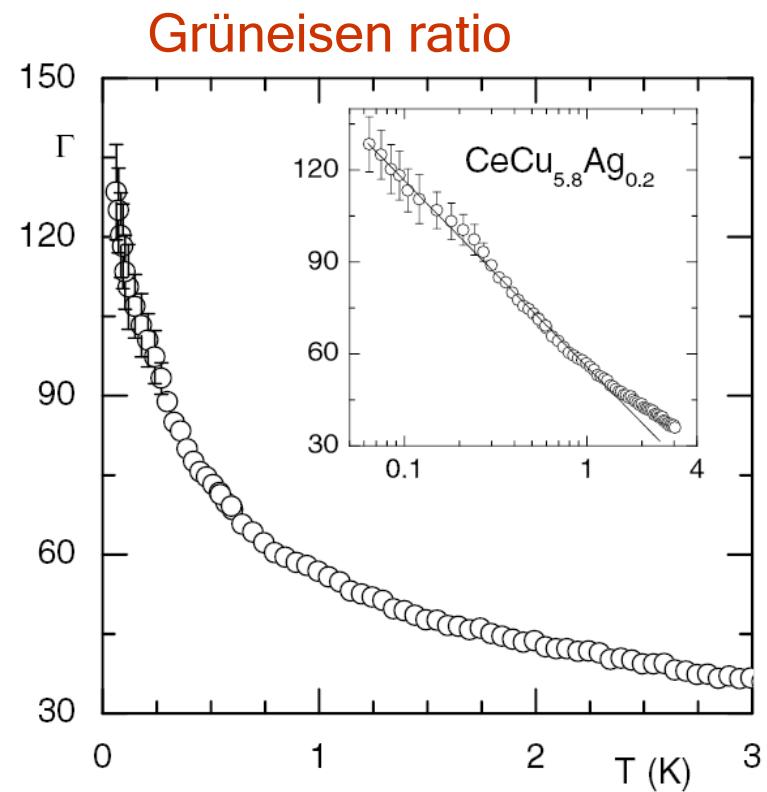
Self-consistency
when $\epsilon = 1^-$

Consistent with the observation in $\text{YbRh}_2(\text{Si}_{0.95}\text{Ge}_{0.05})_2$

CeCu_{5.8}Ag_{0.2}

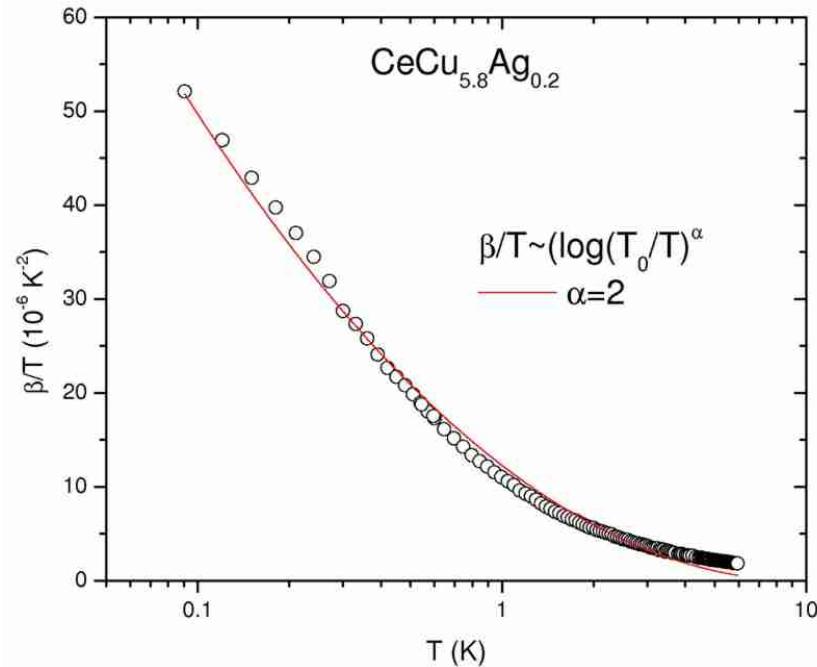


R. Küchler, P. Gegenwart, K. Heuser, E.-W. Scheidt, G. R. Stewart, and F. Steglich, Phys. Rev. Lett. 93, 096402 (2004).



Inconsistent with SDW
Different from YbRh₂Si₂

CeCu_{5.8}Ag_{0.2} - *cont'd*



P. Gegenwart, private communication

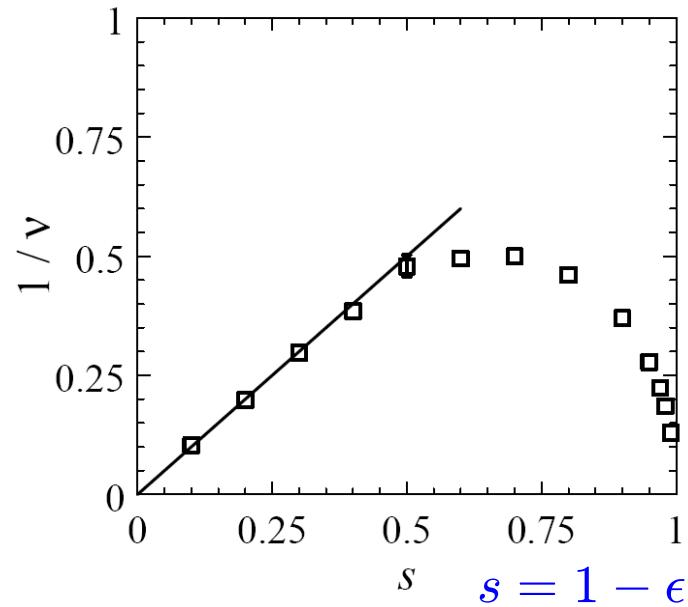
$$c/T \sim \log(T_0/T)$$

$$\beta/T \sim \log^2(T_0/T)$$

$$\Gamma \sim \log(T_0/T)$$

CeCu_{5.8}Ag_{0.2} - *cont'd*

Local criticality in Ising anisotropic case



R. Bulla, N.-H. Tong, and M. Vojta, Phys. Rev. Lett. 91, 170601 (2003); Phys. Rev. Lett. 94, 070604 (2005).

$$1/\nu = 1 - \epsilon$$

When $\epsilon \rightarrow 1^-$

$$\Gamma \sim 1/T^{1/(\nu z)} = 1/T^{0^+} \\ \rightarrow \log T$$

Consistent with the
observation in
CeCu_{5.8}Ag_{0.2}

Conclusion

- Quantum critical point \Leftrightarrow universally **diverging** Grüneisen ratio
 - Identify, characterize QCPs
 - Theory
 - Hypercaling
 - Spin Density Wave
 - Local Quantum Criticality
 - Experiments in quantum critical heavy fermion metals
 - **Observe divergence as predicted**
 - CeNi_2Ge_2 compatible with SDW
 - $\text{YbRh}_2(\text{Si}_{0.95}\text{Ge}_{0.05})_2$ CeCu_{5.8}Ag_{0.2} Inconsistent with SDW
Consistent with LQCP
 - **LQCP XY Ising anisotropic cases**
 - Other strongly correlated systems
-