

# Maxwell Model of Inelastic Collisions

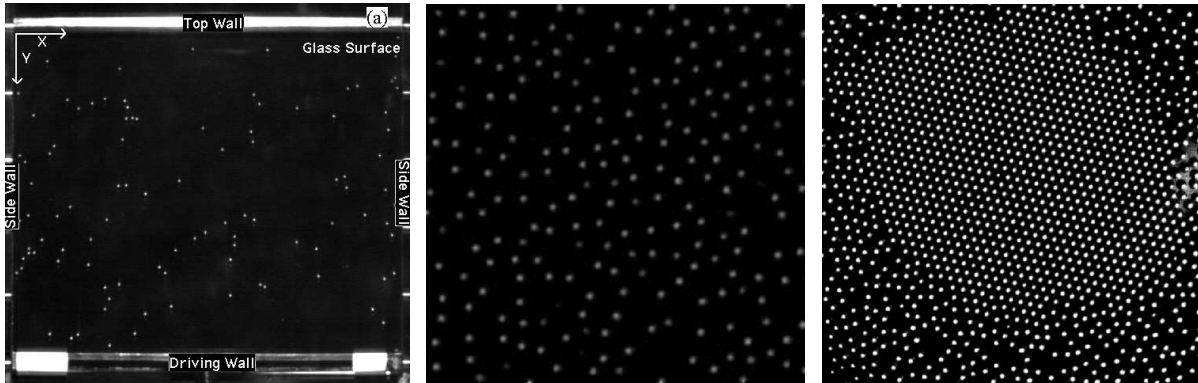
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- I Motivation
- II Freely evolving inelastic gases
- III Impurities
- IV Social dynamics

with: Paul Krapivsky (Boston)

# Experiments: Granular Gases



- Vibration: vertical, horizontal, electrostatic  
Gollub, Swinney, Menon, Aronson, Kudrolli, Urbach
- non-Maxwellian velocity statistics

$$P(v) \sim \exp(-|v|^\alpha) \quad 1 < \alpha < 1.5$$

- Clustering, density inhomogeneities
- Collective phenomena: phase transitions, pattern formation, shocks

# A nonequilibrium gas

## Characteristics

- Hard sphere (contact) interactions
- Dissipative (inelastic) collisions

## Consequences of energy dissipation

- No energy equipartition
- No ergodicity
- Strong velocity correlations

## Challenges

- Hydrodynamics: flow equations
- Kinetic Theory: velocity distributions
- Sharp validity criteria are missing

# The Elastic Maxwell Model

J.C. Maxwell, Phil. Tran. Roy. Soc **157**, 49 (1867)

- Infinite particle system
- Binary collisions
- Random collision partners
- Random impact directions  $\mathbf{n}$
- Elastic collisions ( $\mathbf{g} = \mathbf{v}_1 - \mathbf{v}_2$ )

$$\mathbf{v}_1 \rightarrow \mathbf{v}_1 - \mathbf{g} \cdot \mathbf{n} \mathbf{n}$$

- Mean-field collision process
- Purely Maxwellian velocity distributions

$$P(\mathbf{v}) = \frac{1}{(2\pi T)^{d/2}} \exp\left(-\frac{v^2}{2T}\right)$$

**What about inelastic, dissipative collisions?**

# The Inelastic Maxwell Model

- Inelastic collisions  $r = 1 - 2\epsilon$

$$\mathbf{v}_{1,2} = \mathbf{u}_{1,2} \mp (1 - \epsilon) (\mathbf{g} \cdot \mathbf{n}) \mathbf{n}$$

- Boltzmann equation  $\mathbf{g} \cdot \mathbf{n} \rightarrow \langle g \rangle$

$$\frac{\partial P(\mathbf{v}, t)}{\partial t} = \int d\mathbf{n} \int d\mathbf{u}_1 \int d\mathbf{u}_2 \langle g \rangle P(\mathbf{u}_1, t) P(\mathbf{u}_2, t) \times \left\{ \delta(\mathbf{v} - \mathbf{v}_1) - \delta(\mathbf{v} - \mathbf{u}_1) \right\}$$

- Fourier transform

Krupp 1967

$$F(\mathbf{k}, t) = \int d\mathbf{v} e^{i\mathbf{k} \cdot \mathbf{v}} P(\mathbf{v}, t)$$

- Closed equations  $\mathbf{q} = (1 - \epsilon)\mathbf{k} \cdot \mathbf{n} \mathbf{n}$

$$\frac{\partial}{\partial t} F(\mathbf{k}, t) + F(\mathbf{k}, t) = \int d\mathbf{n} F[\mathbf{k} - \mathbf{q}, t] F[\mathbf{q}, t],$$

**Theory is analytically tractable**

# One Dimension

- **Scaling of isotropic velocity distribution**

$$P(\mathbf{v}, t) \rightarrow \frac{1}{T^{d/2}} \Phi \left( \frac{|\mathbf{v}|}{T^{1/2}} \right) \quad \text{or} \quad F(k, t) \rightarrow f(k^2 T)$$

- **Nonlinear and nonlocal equation**

$$-2\epsilon(1 - \epsilon)f'(x) + f(x) = f(\epsilon^2 x)f((1 - \epsilon)^2 x)$$

- **Exact solution**

$$f(x) = (1 + \sqrt{x}) e^{-\sqrt{x}} \cong 1 - \frac{1}{2}x + \frac{1}{3}x^{3/2} + \dots$$

- **Lorentzian<sup>2</sup> velocity distribution**

$$\Phi(v) = \frac{2}{\pi} \frac{1}{(1 + v^2)^2}$$

- **Algebraic tail**

Baldassari 2001

$$\Phi(v) \sim v^{-4} \quad w \gg 1$$

**Universal scaling function, exponent**

# Scaling, Nontrivial Exponents

- Freely cooling case

$$T = \langle v^2 \rangle \sim t^{-2}$$

- Governing equation for scaling function

$$-\lambda x \Phi'(x) + \Phi(x) = \int d\mathbf{n} \Phi(x\xi) \Phi(x\eta)$$

$$\lambda = 2\epsilon(1 - \epsilon)/d, \quad \xi = 1 - (1 - \epsilon^2) \cos^2 \theta, \quad \eta = (1 - \epsilon)^2 \cos^2 \theta$$

- Power-law tails

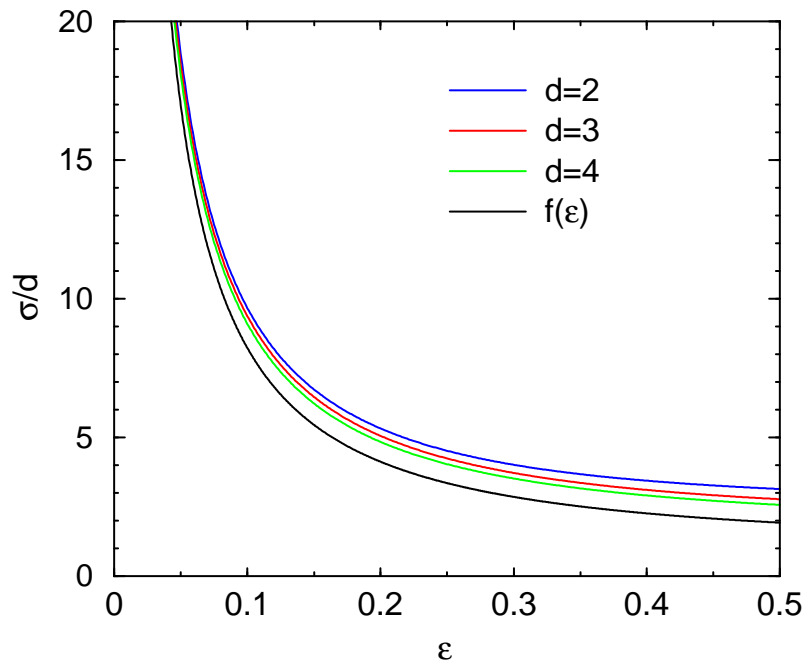
$$\Phi(v) \sim v^{-\sigma}, \quad v \rightarrow \infty.$$

- Exact solution for the exponent  $\sigma$

$$1 - \epsilon(1 - \epsilon) \frac{\sigma - d}{d} = {}_2F_1 \left[ \frac{d - \sigma}{2}, \frac{1}{2}; \frac{d}{2}; 1 - \epsilon^2 \right] + (1 - \epsilon)^{\sigma - d} \frac{\Gamma\left(\frac{\sigma - d + 1}{2}\right) \Gamma\left(\frac{d}{2}\right)}{\Gamma\left(\frac{\sigma}{2}\right) \Gamma\left(\frac{1}{2}\right)}$$

**Nonuniversal tails, exponents depend on  $\epsilon$ ,  $d$**

# The exponent $\sigma$



- Maxwellian distributions:  $d = \infty, \epsilon = 0$

- Diverges in high dimensions

$$\sigma \cong d f(\epsilon)$$

- Diverges for low dissipation

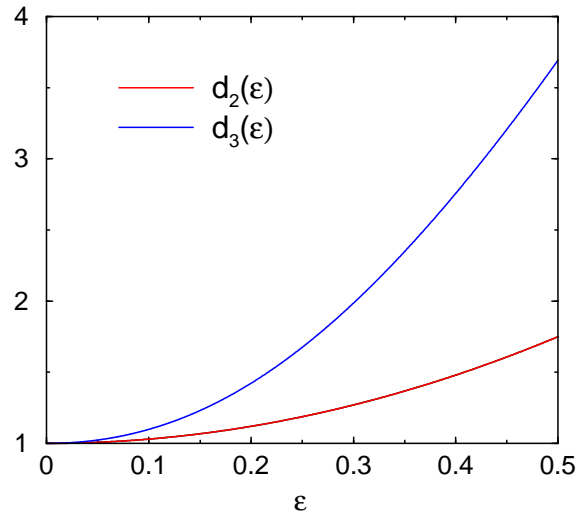
$$\sigma \cong d \epsilon^{-1}$$

- In practice, huge

$$\sigma(d = 3, r = 0.8) \cong 30!$$



# Dynamics



- Moments of the velocity distribution

$$M_{2n}(t) = \int d\mathbf{v} |\mathbf{v}|^{2n} P(\mathbf{v}, t)$$

- Multiscaling asymptotic behavior

$$M_{2n} \sim M_2^{\xi_n/2} \quad \xi_n = \begin{cases} n & d < d_n(\epsilon), \\ \alpha_n(\epsilon) & d > d_n(\epsilon). \end{cases}$$

- Nonlinear multiscaling spectrum (1D):

$$\alpha_n(\epsilon) = \frac{1 - \epsilon^{2n} - (1 - \epsilon)^{2n}}{1 - \epsilon^2 - (1 - \epsilon)^2}$$

**Sufficiently large moments exhibit multiscaling**

# Velocity Correlations

- Definition (correlation between  $v_x^2$  and  $v_y^2$ )

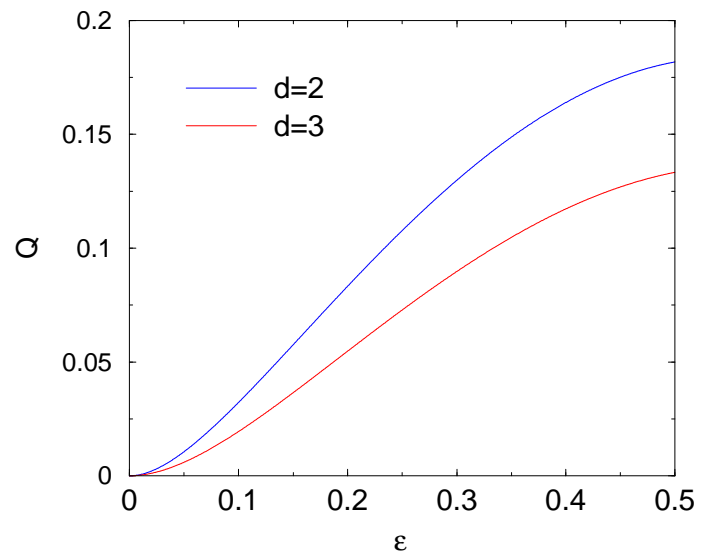
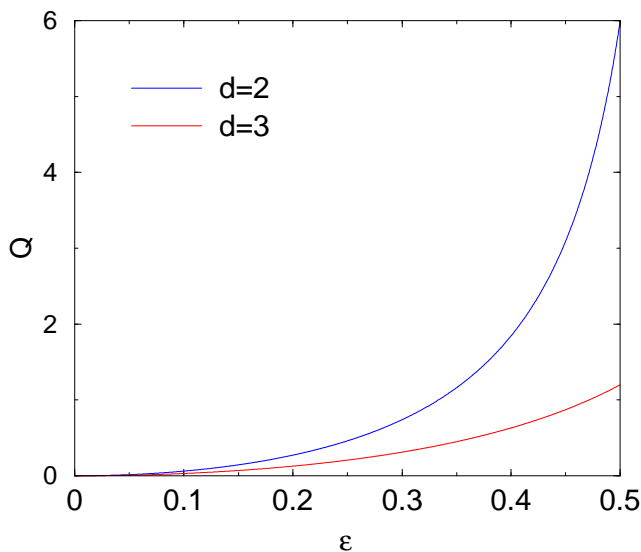
$$Q = \frac{\langle v_x^2 v_y^2 \rangle - \langle v_x^2 \rangle \langle v_y^2 \rangle}{\langle v_x^2 \rangle \langle v_y^2 \rangle}$$

- Unforced case (freely evolving)  $P(v) \sim v^{-\sigma}$

$$Q = \frac{6\epsilon^2}{d - (1 + 3\epsilon^2)}$$

- Forced case (white noise)  $P(v) \sim e^{-|v|}$

$$Q = \frac{6\epsilon^2(1 - \epsilon)}{(d + 2)(1 + \epsilon) - 3(1 - \epsilon)(1 + \epsilon^2)}$$



**Correlations diminish with energy input**

# The “brazil nut” problem

- Fluid background: mass 1
- Impurity: mass  $m$
- Theory: Lorentz-Boltzmann equation
- Series of transition masses

$$1 < m_1 < m_2 < \dots < m_\infty$$

- Ratio of moments diverges asymptotically

$$\frac{\langle v_I^{2n} \rangle}{\langle v_F^{2n} \rangle} \sim \begin{cases} c_n & m < m_n; \\ \infty & m > m_n. \end{cases}$$

- Light impurity: moderate violation of equipartition, impurity mimics the fluid
- Heavy impurity: extreme violation of equipartition, impurity sees a static fluid

**series of phase transitions**

# Conclusions

- non-Maxwellian velocity distributions
- Power-law high energy tails
- non-universal exponents
- Multiscaling of the moments, Temperature insufficient to characterize large moments
- Correlations between velocity components

Ben-Naim and Krapivsky, PRE 61, 5 (00); JPA 35, L147 (02); PRE 66, 011309 (02); EPJE 8, 507 (02).

Ernst and Brito, EL 58, 182 (02); PRE 65, 040301 (02); JSP 109, 407 (02).

Bobylev, Carrido and Gamba, JSP 98, 743 (00).

Baldassarri, Marconi and Puglisi, EL 58, 14 (02); PRE 65, 051301 (02); 66, 011301 (02).

Bobylev and Cercignani, JSP 106, 1039 (02).

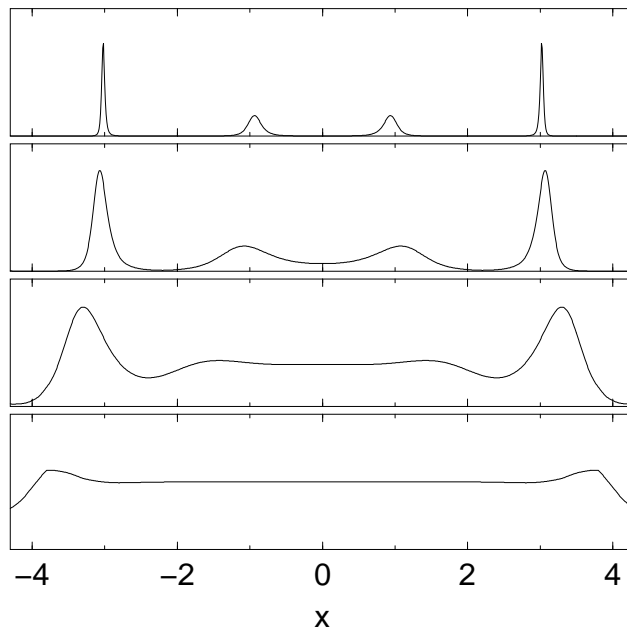
Santos and Dufty, PRL 86, 4823 (01); PRE 64, 051305 (01)

# The Compromise Model

- Opinion  $-\Delta < x < \Delta$
- Reach compromise in pairs Weisbuch 2001

$$(x_1, x_2) \rightarrow \left( \frac{x_1 + x_2}{2}, \frac{x_1 + x_2}{2} \right)$$

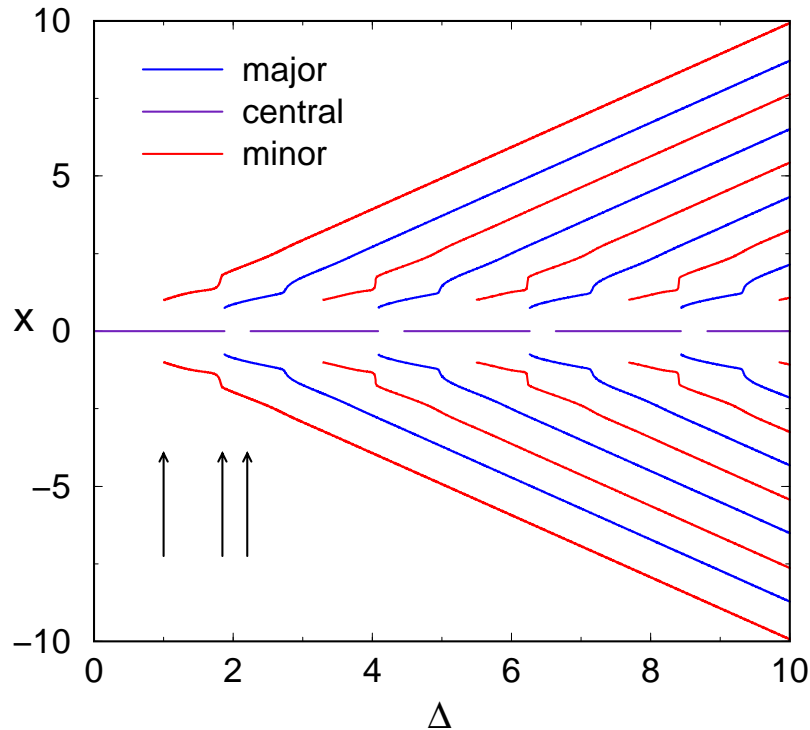
- As long as we are close  $|x_1 - x_2| < 1$



$$P_{\infty}(x) = \sum_i m_i \delta(x - x_i)$$

**Final State: localized clusters**

# Bifurcations and Patterns



- Periodic bifurcations

$$x(\Delta) = x(\Delta + L)$$

- Alternating major-minor pattern
- Critical behavior

$$m \sim (\Delta - \Delta_c)^\alpha \quad \alpha = 3 \text{ or } 4.$$

**Self-similar structure**