The Centre for Theoretical and Computational Chemistry



The dynamics of angular degrees of freedom: new basis set and grid representations of Hamiltonian operators

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Outline

- ✓ Part I: Recipe & ingredients for a vibrational calculation
- ✓ Part II: A new basis set for angular motion & comparison with Legendre functions
- Part III: Overview of localized & delocalized representations, construction of localized mixed basis functions, applications

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✓ Part IV: Conclusions



Recipe & ingredients for a vibrational calculation



Vibrational spectrum of H₂O

- How to compute very accurately the vibrational levels of H₂O in the electronic ground state?
 - ✓ Compute potential energy surface on dense grid in (r_1, r_2, α) space
 - Make decision: definition of potential energy operator V(r₁, r₂, α) directly on grid or via analytical model function
 - Select basis functions for description of vibrational wave functions. If V(r₁, r₂, α) is defined on discrete set of points basis functions are still needed for representation of T operator
- Popular basis functions for radial degrees of freedom:

$$\mu_{n}(x) = \sqrt{\frac{2}{b-a}} \sin\left(\frac{n\pi(x-a)}{(b-a)}\right), n = 1, 2, ..., N$$

$$\nu_{n}(x) = \sigma \cos\left(\frac{n\pi(x-a)}{(b-a)}\right), n = 0, 1, 2, ..., N - 1 \begin{cases} \sigma = \sqrt{\frac{1}{b-a}}, n = 0\\ \sigma = \sqrt{\frac{2}{b-a}}, n \neq 0 \end{cases}$$
See e.g. Colbert & Miller, JCP 96, 1982 (1992)

Vibrational basis sets

- Why are the $\mu_n(x)$ and $\nu_n(x)$ functions popular?
 - ✓ They yield analytic expressions for $<\mu_m|T|\mu_n>$ and $<\nu_m|T|\nu_n>$ for finite and infinite definition intervals [a,b]
 - ✓ They are associated with an equidistant quadrature grid (relation to Chebychev)
 - ✓ The quadrature rule is of Gaussian accuracy (discrete orthogonality)

$$\int_{a}^{b} f(x) dx = w \sum_{k=1}^{N} f(x_{k}) \begin{cases} w = \frac{b-a}{N+1}, \text{ for } \mu_{n} \\ w = \frac{b-a}{N}, \text{ for } \nu_{n} \end{cases}$$

- Which basis sets are appropriate for bending motion?
 - ✓ The bending kinetic energy operator is:

$$\hat{T}_{bend} = -c_{bend} \left(\frac{\partial^2}{\partial x^2} + \cot(x) \frac{\partial}{\partial x} \right) \qquad c_{bend} = \frac{1}{26}$$

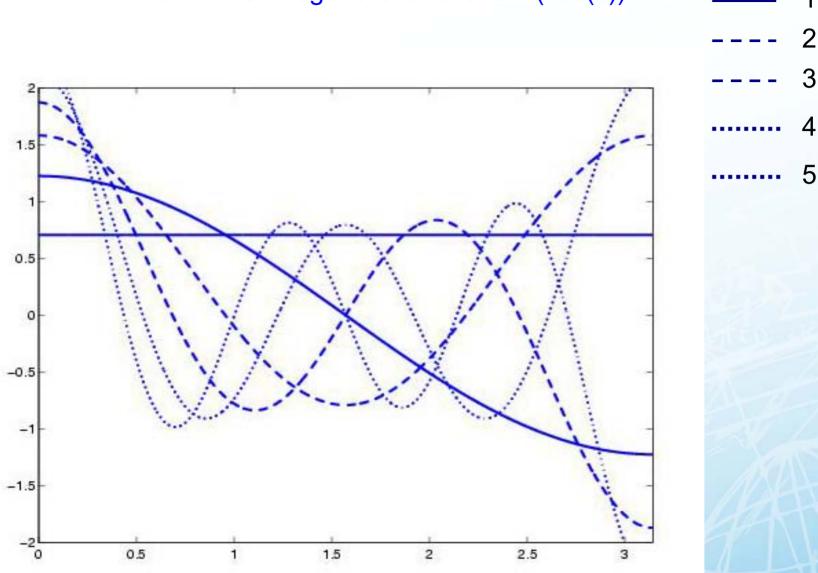
Legendre basis for bending motion

- ✓ In this form, T_{bend} is hermitian on [0, π] with respect to volume element sin(x) dx
- $\checkmark~\mu_n(x)$ and $\nu_n(x)$ functions perform badly as basis functions for T_{bend}
- ✓ The standard basis functions for T_{bend} are derived from Legendre polynomials $P_I(x)$:

$$\sigma_0 P_0(\cos(x)) = \sqrt{\frac{1}{2}} \qquad \qquad \sigma_2 P_2(\cos(x)) = \sqrt{\frac{5}{2^5}} (3\cos(x) + 5\cos(3x))$$

$$\sigma_1 P_1(\cos(x)) = \sqrt{\frac{3}{2}} \cos(x) \qquad \sigma_3 P_3(\cos(x)) = \sqrt{\frac{9}{2^{13}}} (9 + 20\cos(2x) + 35\cos(4x))$$

- ✓ $P_I(cos(x))$ are the eigenfunctions of T_{bend} → diagonal analytic representation
- \checkmark P_I(cos(x)) are associated with quadrature rule of Gaussian accuracy
- ✓ Grid point density increases moderately towards interval limits





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A new basis set for angular motion & comparison with Legendre functions



The η_n (x) angular basis functions

• Can we formulate basis functions for bending motion that are analog to the $\mu_n(x)$ and $\nu_n(x)$ functions?

How about:

$$\eta_n(x) = \sqrt{\frac{2}{\pi}} \frac{\sin(nx)}{\sqrt{\sin(x)}}, n = 1, 2, \dots, N$$

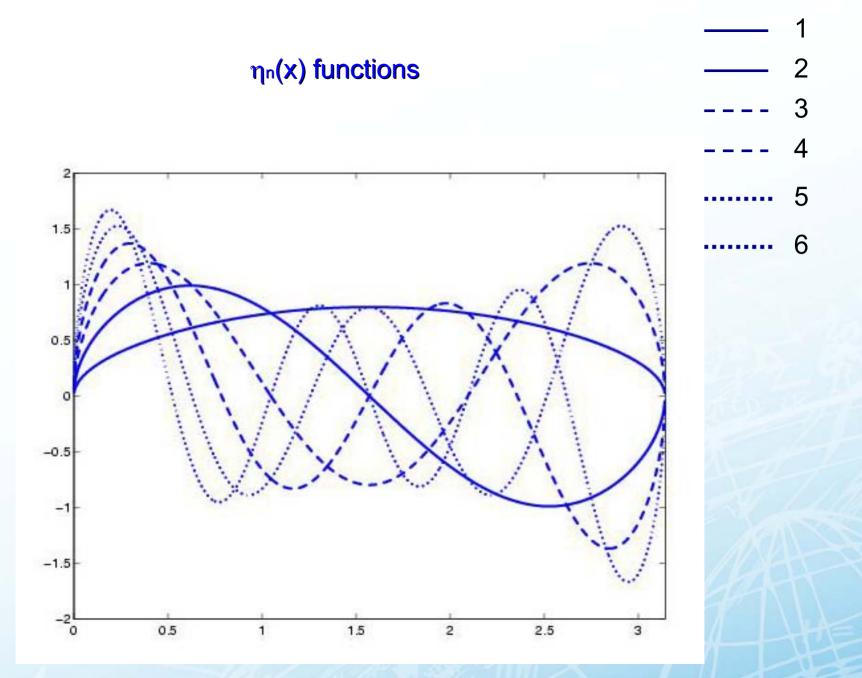
• Properties of $\eta_n(x)$ functions :

- \checkmark they are orthonormal on [0, π] wrt to volume element sin(x) dx
- ✓ the matrix elements $<\eta_m|T_{bend}|\eta_n>$ have simple analytic solutions
- ✓ they are related to an equidistant quadrature grid
- ✓ the quadrature rule

 \checkmark

$$\int_{0}^{\pi} f(x)\sin(x)\,dx = \sum_{k=1}^{N} w_{k}f(x_{k}) \qquad w_{k} = \sqrt{\frac{\pi}{N+1}}\sin\left(\frac{\pi k}{N+1}\right)$$

is of Gaussian accuracy



Definition of model Hamiltonian

- We compare the performance of $\eta_n(x)$ and $P_l(cos(x))$ basis functions
- Model system: pure bending motion of H₂O

 $\hat{H} = \hat{T}_{bend} + c_0 + c_1 (\cos(x)) + c_2 (\cos(x))^2 + c_3 (\cos(x))^3 + c_4 (\cos(x))^4$

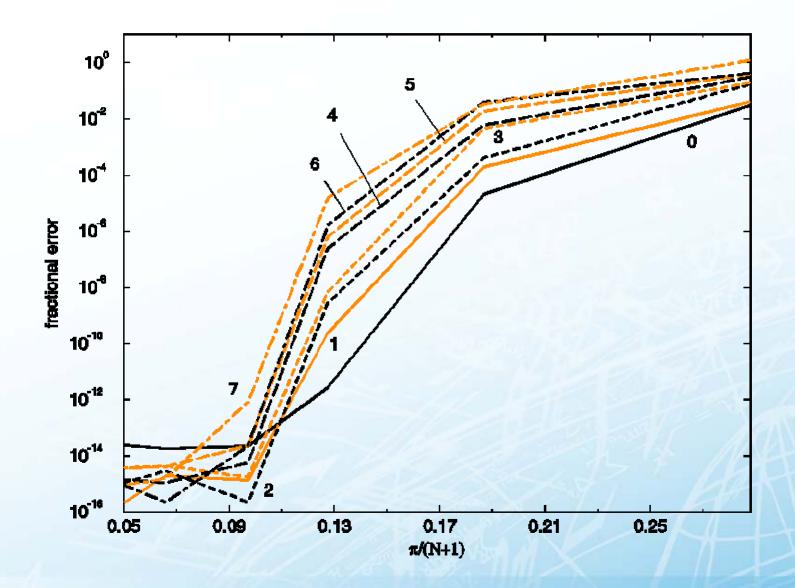
relatively harmonic potential \rightarrow well suited for Legendre basis

Variational Basis Represention (VBR) for H

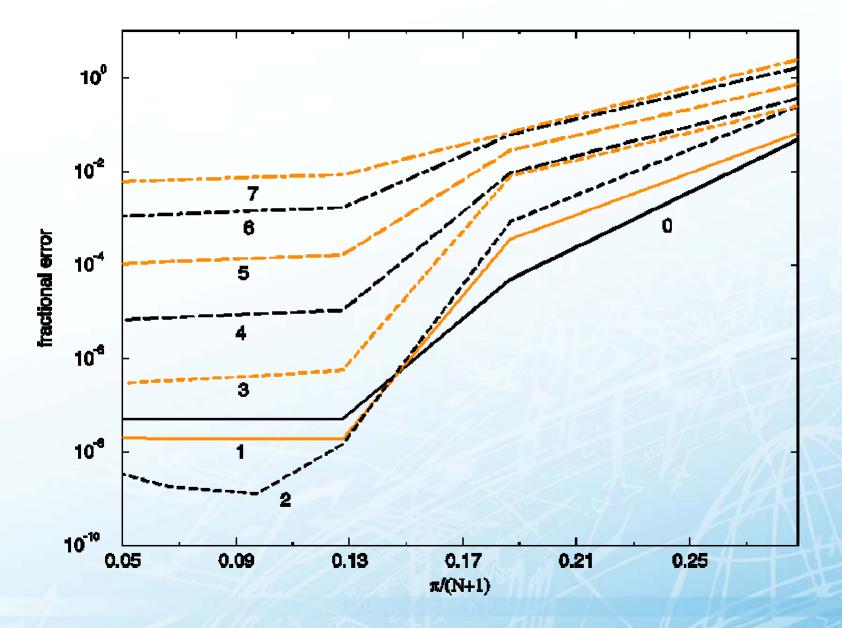
$$\left\langle \varphi_{m}(x) \middle| \hat{O} \middle| \varphi_{n}(x) \right\rangle = \int_{a}^{b} \varphi_{m}(x) \left[\hat{O} \varphi_{n}(x) \right] w(x) dx$$

✓ For a true VBR, all matrix elements must be evaluated exactly

Legendre VBR for H₂O bending



$\eta_n(x)$ VBR for H₂O bending



How to improve performance of $\eta_n(x)$?

- Obviously, a basis formed exclusively by $\eta_n(x)$ functions is incomplete
- Can we use the complementary functions?

$$\sigma_n \frac{\cos(nx)}{\sqrt{\sin(x)}}, n = 0, 1, 2..., N-1$$

• Can we supplement the $\eta_n(x)$ functions? For example:

$$s_{r}(x) = \exp(-r\sin(x)) \quad s_{t}(x) = \exp(-t\sin(x)) \\ \{\eta_{m}(x)(1-s_{r}(x)), P_{n}(\cos(x))s_{t}(x)\}$$

- ✓ Switching functions $(1-s_r(x))$, $s_t(x)$ keep basis orthogonal
- Evaluation of matrix elements tedious

Supplementation of $\eta_n(x)$ functions

Is direct basis extension an option? For example:

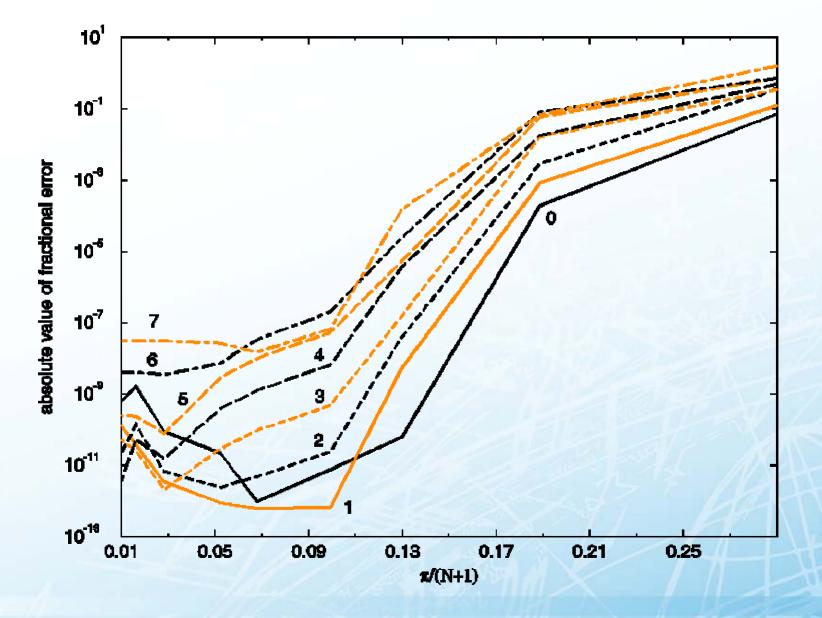
$$\begin{cases} \eta_1(x), \eta_2(x), \dots, \eta_N(x), \\ \sigma_0 P_0(\cos(x)), \sigma_1 P_1(\cos(x)), \dots, \sigma_M P_M(\cos(x)) \end{cases} \end{cases}$$

- \checkmark We construct representation of H in this mixed basis \rightarrow symmetric matrix **A**
- ✓ Loewdin (symmetric) orthogonalization of mixed basis:
 - Diagonalization of overlap matrix yields eigenvector matrix U and the matrix X=diag($1/\sqrt{\epsilon_1}$, $1/\sqrt{\epsilon_2}$, ..., $1/\sqrt{\epsilon_N}$)
 - Orthogonalization of mixed basis according to:

$$(U X U^T) A (U X U^T) = H$$

 \checkmark **H** is the desired representation of H in orthonormal mixed basis

Mixed basis (+ P₀,P₁) VBR H₂O bending





Overview of localized & delocalized representations, construction of localized mixed basis functions, applications



Overview: representations

- We differentiate between infinitely localized, nearly localized and delocalized basis functions
- Discrete representations are only possible for local operators
- Discrete Variable Representation (DVR) of local operator O is a matrix diagonal over the grid points x_k.
- If the grid is related to orthogonal basis functions φ_m(x) through a quadrature rule of the form:

$$\int_{a}^{b} f(x) w(x) dx = \sum_{k=1}^{N} w_{k} f(x_{k})$$

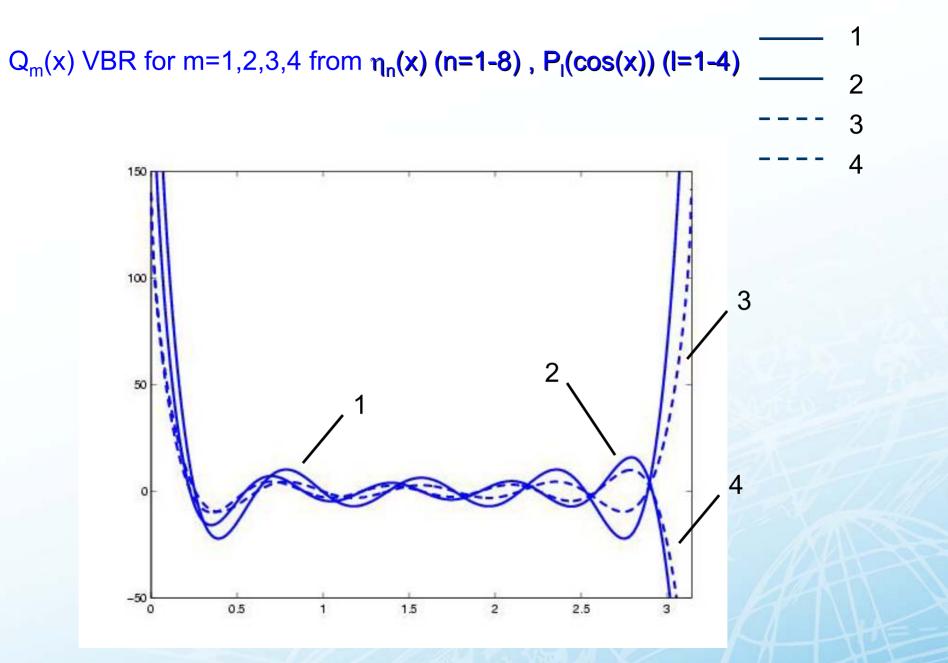
then we can define a Finite Basis Representation (FBR):

$$\int_{a}^{b} \varphi_{m}(x) \left[\hat{O} \varphi_{n}(x) \right] w(x) dx \approx \sum_{k=1}^{N} \varphi_{m}(x_{k}) \left[\hat{O} \varphi_{n}(x_{k}) \right] w_{k}$$

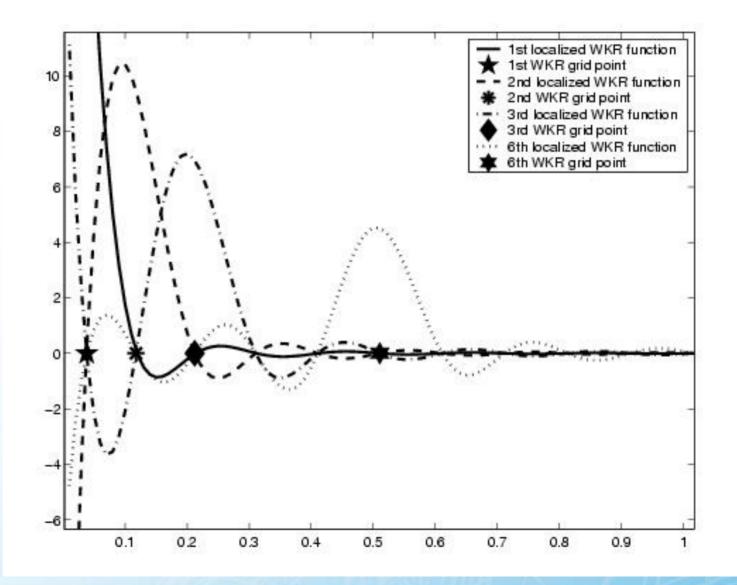
VBR, FBR, DVR, NDVR

- FBR matrix is an approximation to the VBR matrix
- FBR and DVR are equivalent: $FBR O = \Lambda O \Lambda^{t}$
- In analogy we can define:

- $^{VBR}O = \Lambda ^{NDVR}O \Lambda^{t}$
- NDVR matrix is an approximation to the DVR matrix:
 - \checkmark NDVR basis functions are approximately localized at grid points x_k.
- The term "DVR calculation" is not exact: DVR H = NDVR T + DVR V
 - ✓ DVR results are not variational!
- How to perform "DVR calculation" for mixed basis functions Q_n(x)?
 - ✓ Definition of grid points through zeroes of $Q_{N+M+1}(x)$ (from $\eta_{N+1}(x)$ and $P_M(cos(x))$) → explicit derivation of orthogonal mixed basis
 - Stablishment of quadrature rule for mixed basis \rightarrow derivation of Λ matrix

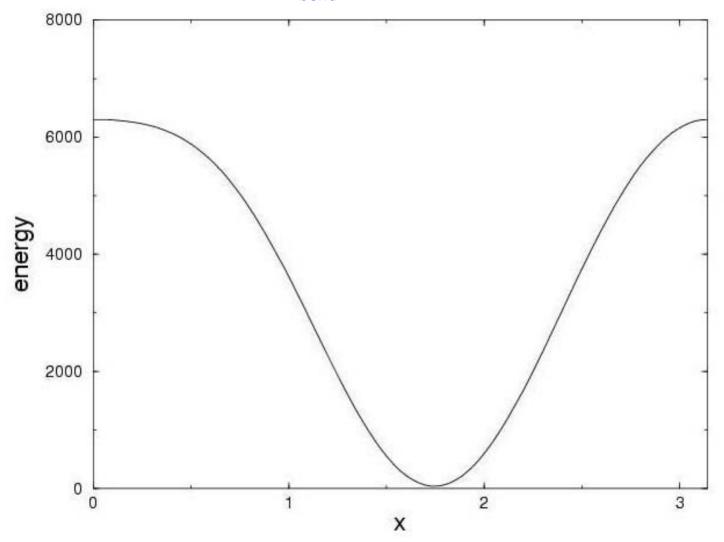


$Q_m(x)$ NDVR localized at x_k , k=1,2,3,6 from $\eta_n(x)$ (n=1-30), P_I(cos(x)) (l=1-2)

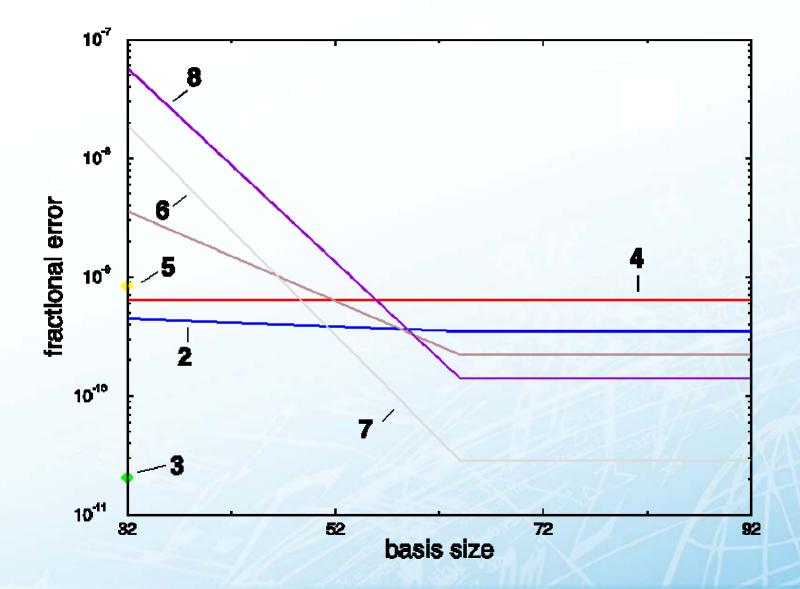


 $V_1 = 300 + 3000 \cos(x) + 8000 (\cos(x))^2 - 3000 (\cos(x))^3 - 2000 (\cos(x))^4$

c_{bend} = 10.0

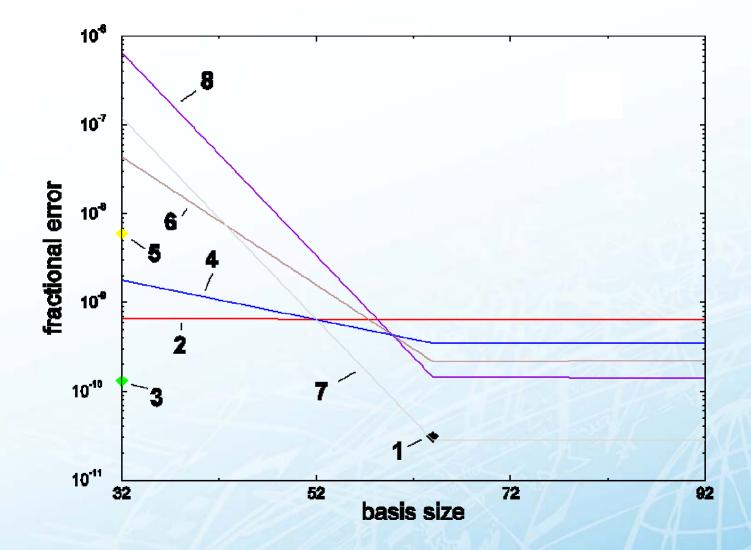


Legendre VBR for V₁

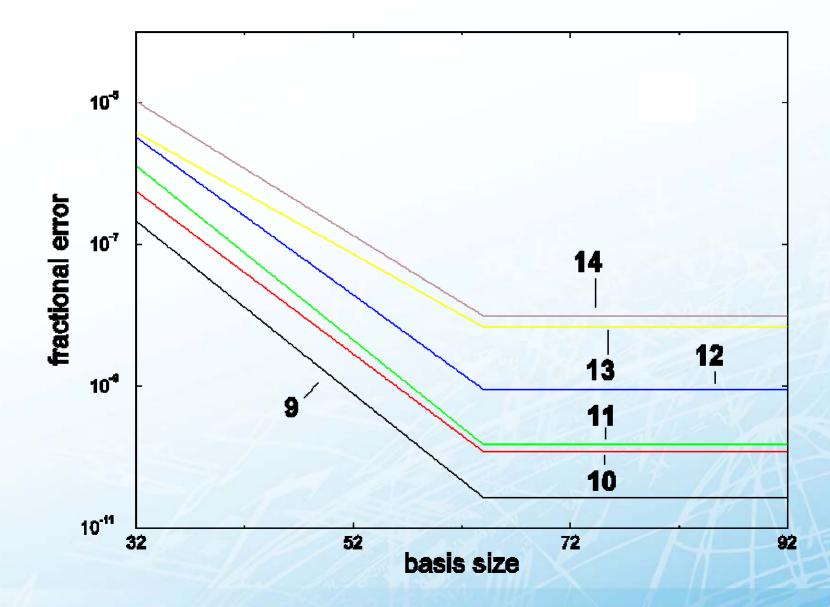


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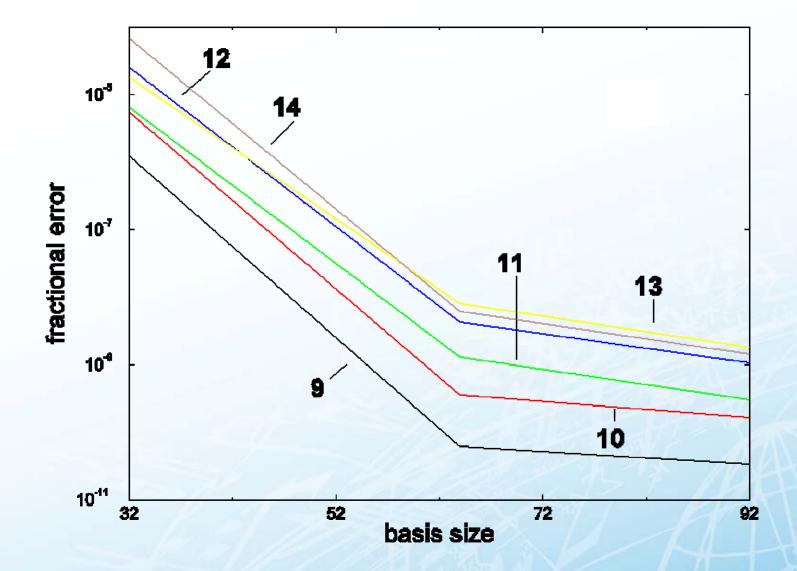
 $Q_m(x)$ VBR for V₁ from $\eta_n(x)$ (n=1-N-2), P_I(cos(x)) (l=1-2)



Legendre VBR for V₁

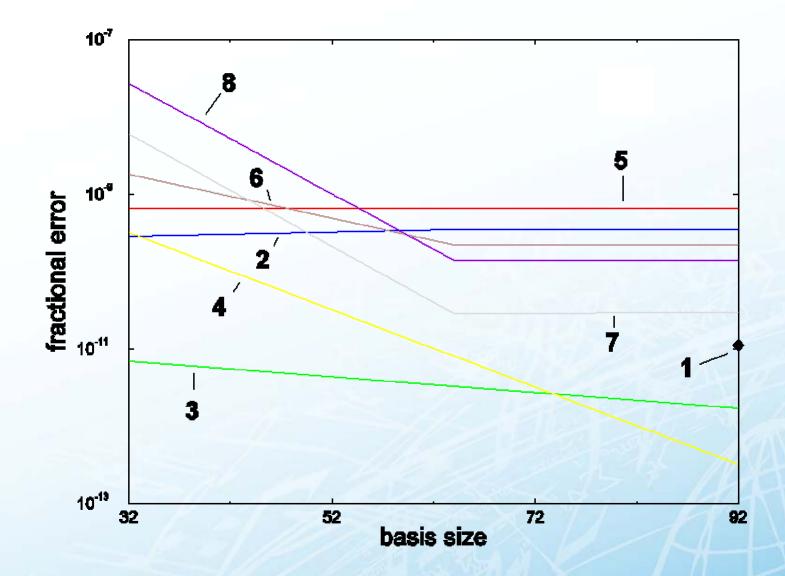


 $Q_m(x)$ VBR for V₁ from $\eta_n(x)$ (n=1-N-2), P_I(cos(x)) (l=1-2)

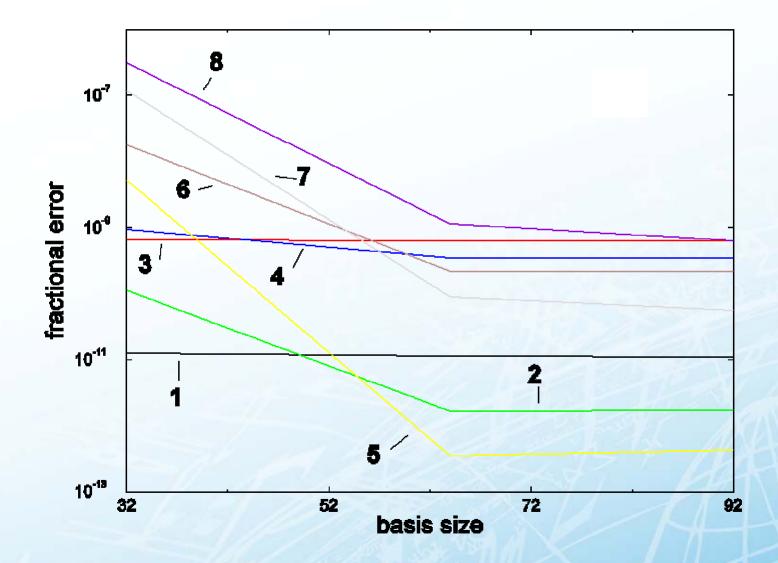


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Legendre DVR for V_1

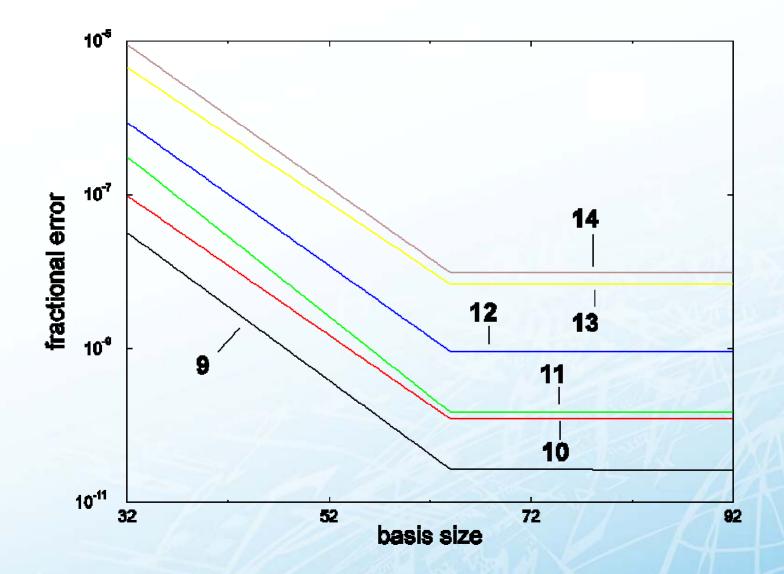


 $Q_m(x)$ DVR for V₁ from $\eta_n(x)$ (n=1-N-2), P_I(cos(x)) (l=1-2)

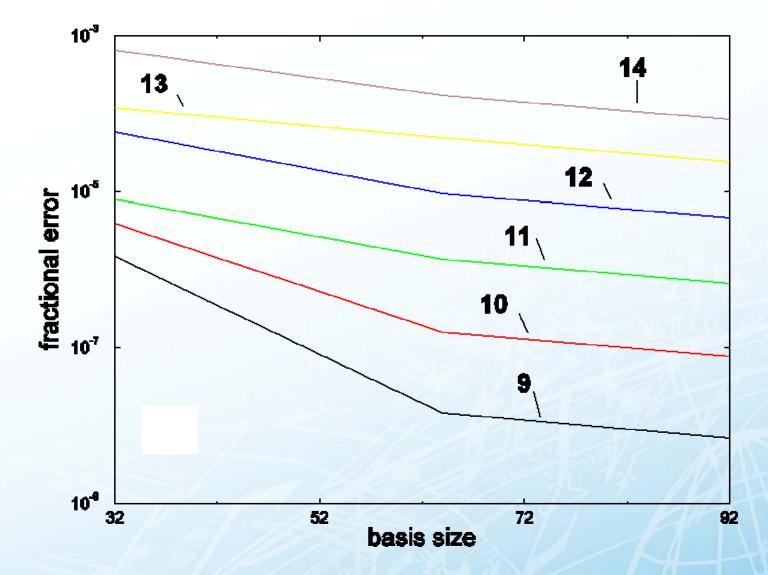


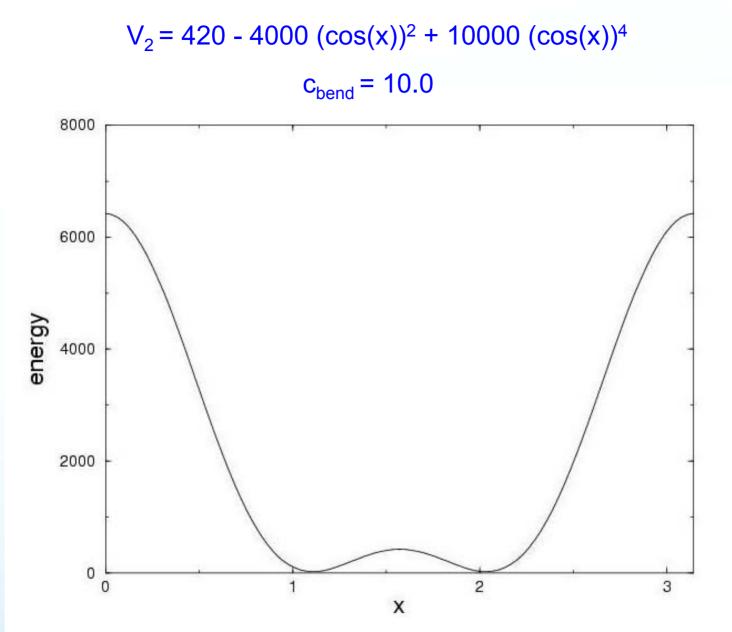
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Legendre DVR for V₁



 $Q_{m}(x)$ DVR for V₁ from $\eta_{n}(x)$ (n=1-N-2), P_I(cos(x)) (l=1-2)

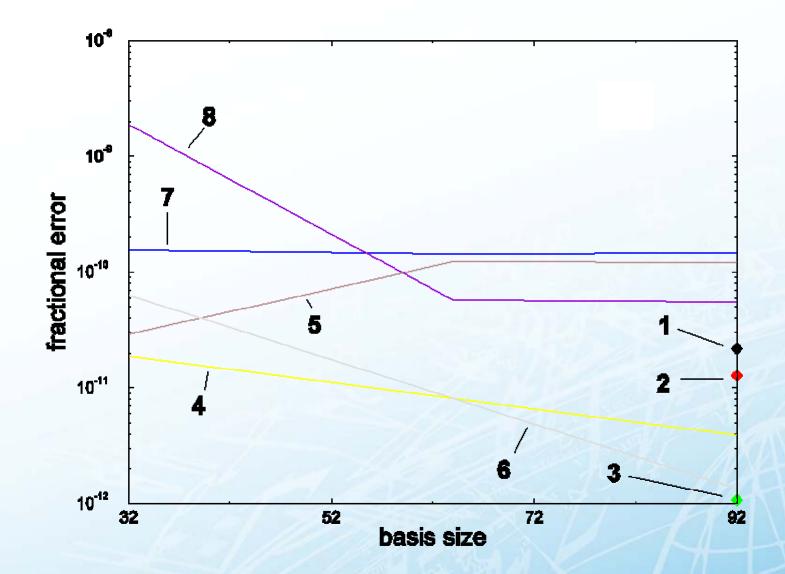




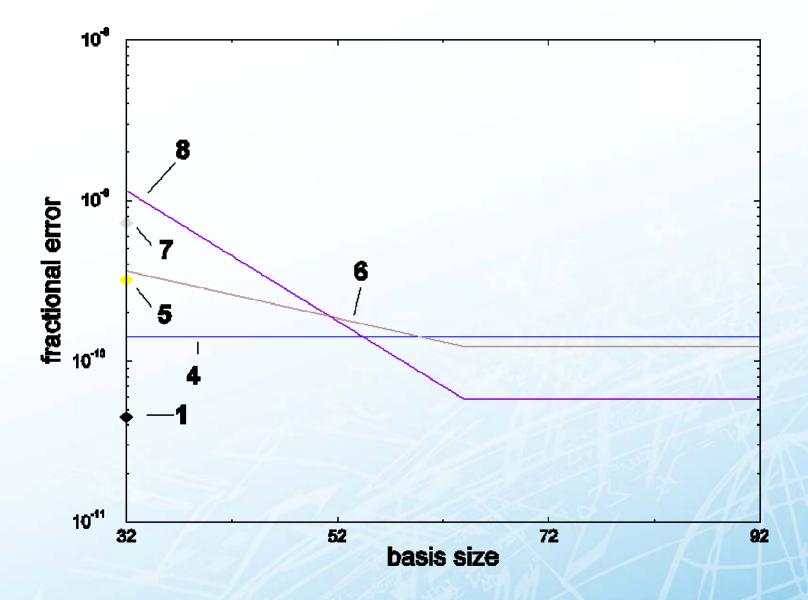
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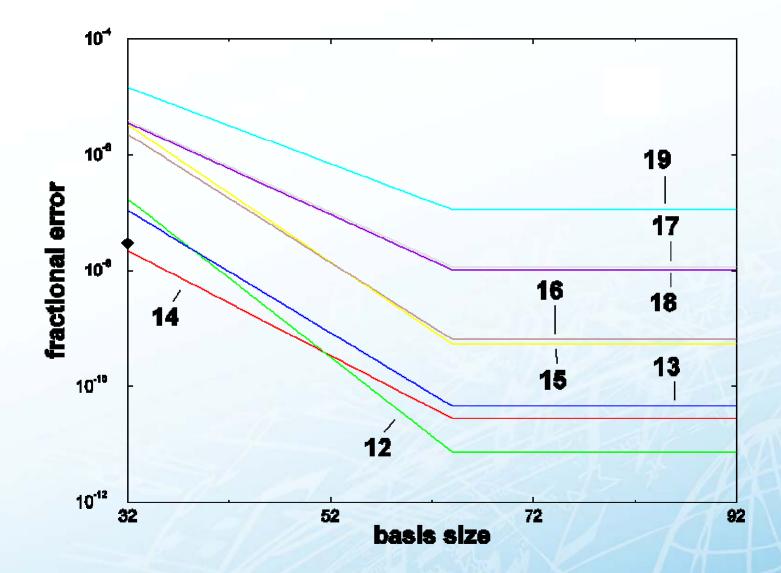
Legendre VBR for V₂



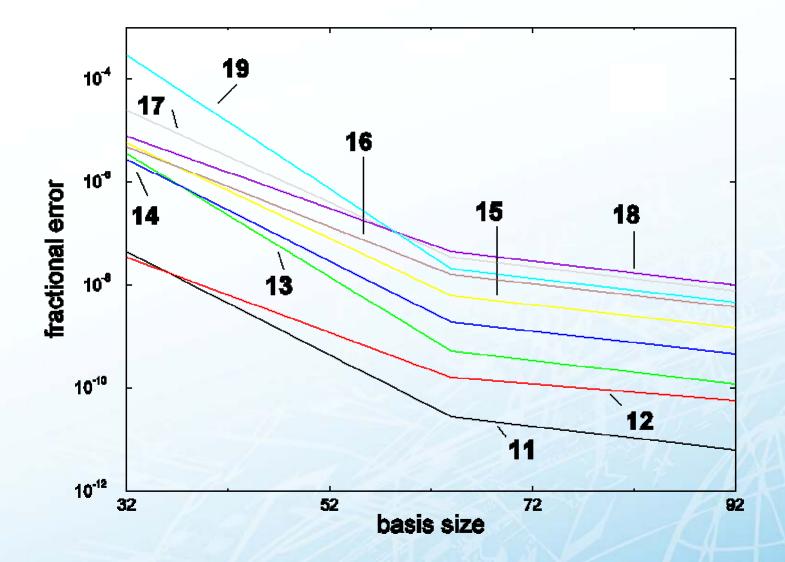
 $Q_m(x)$ VBR for V₂ from $\eta_n(x)$ (n=1-N-2) , P_I(cos(x)) (l=1-2)



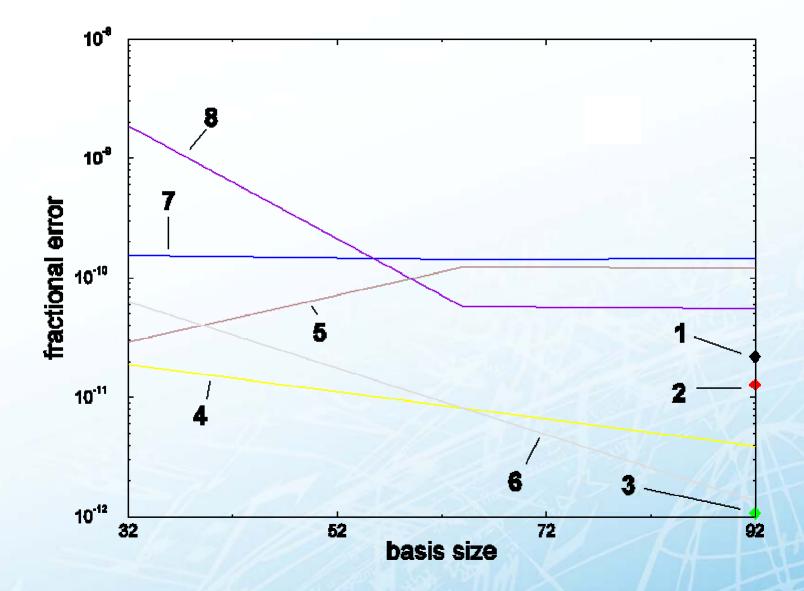
Legendre VBR for V₂



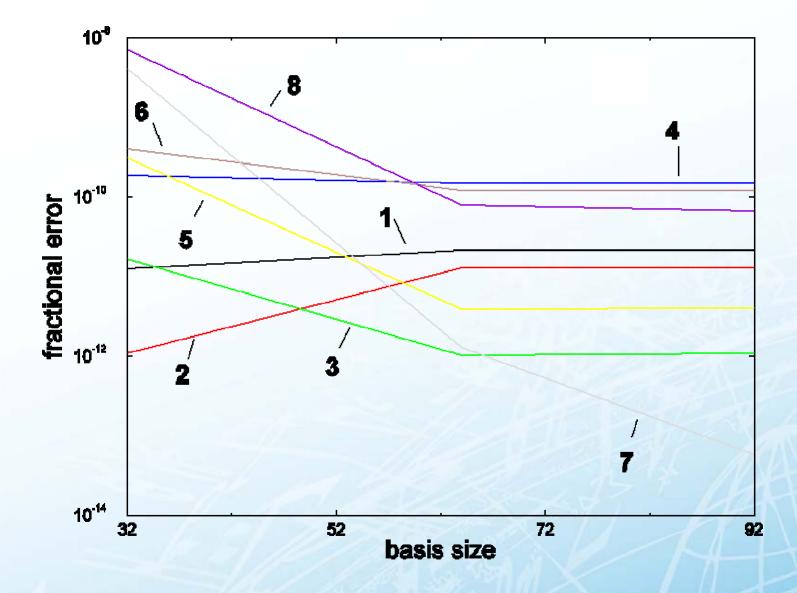
 $Q_m(x)$ VBR for V₂ from $\eta_n(x)$ (n=1-N-2), P_I(cos(x)) (l=1-2)



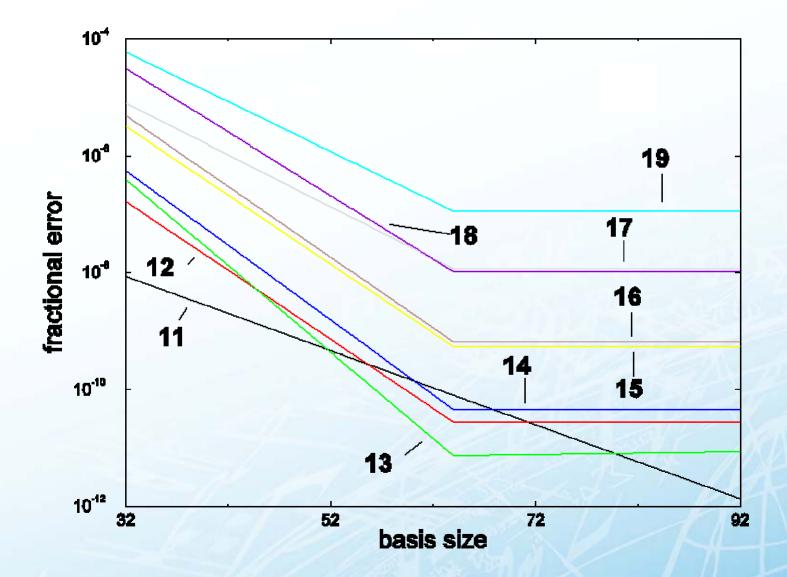
Legendre DVR for V_2



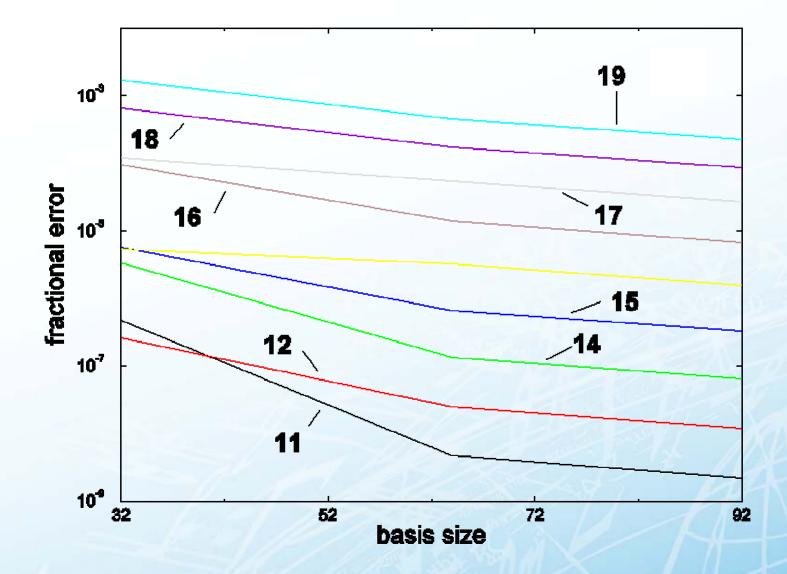
 $Q_m(x)$ DVR for V₂ from $\eta_n(x)$ (n=1-N-2), P_I(cos(x)) (l=1-2)



Legendre DVR for V₂



 $Q_m(x)$ DVR for V₂ from $\eta_n(x)$ (n=1-N-2), P_I(cos(x)) (l=1-2)





Conclusions



- Legendre functions form for many applications a good basis set for bending degrees of freedom
- However, they offer a limited flexibility in particular for the description of states with larger density close to the center of the interval
- The combination of η_n(x) and P_l(cos(x)) functions appears to be an interesting alternative to the pure Legendre basis

For mixed basis VBR calculations:

- ✓ all matrix elements can be evaluated analytically
- ✓ computationally efficient because of Loewdin orthogonalization
- ✓ more homogeneous accuracy distribution for different eigenstates

For mixed basis DVR calculations:

- ✓ explicit orthogonalization complicated → but needs to be performed only once since kinetic energy operator is always the same
- ✓ quadrature rule for mixed basis set has been derived
- ✓ accuracy similar to Legendre DVR can be reached, but further improvement necessary

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