# The dynamics of angular degrees of freedom: new basis set and grid representations of Hamiltonian operators 

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## Outline

$\checkmark$ Part I: Recipe \& ingredients for a vibrational calculation
$\checkmark$ Part II: A new basis set for angular motion \& comparison with Legendre functions
$\checkmark$ Part III: Overview of localized \& delocalized representations, construction of localized mixed basis functions, applications
$\checkmark$ Part IV: Conclusions

## Part I

Recipe \& ingredients for a vibrational calculation

## Vibrational spectrum of $\mathrm{H}_{2} \mathrm{O}$

- How to compute very accurately the vibrational levels of $\mathrm{H}_{2} \mathrm{O}$ in the electronic ground state?
$\checkmark$ Compute potential energy surface on dense grid in $\left(r_{1}, r_{2}, \alpha\right)$ - space
$\checkmark$ Make decision: definition of potential energy operator $V\left(r_{1}, r_{2}, \alpha\right)$ directly on grid or via analytical model function
$\checkmark$ Select basis functions for description of vibrational wave functions. If $V\left(r_{1}, r_{2}, \alpha\right)$ is defined on discrete set of points basis functions are still needed for representation of T operator
- Popular basis functions for radial degrees of freedom:

$$
\begin{aligned}
& \mu_{n}(x)=\sqrt{\frac{2}{b-a}} \sin \left(\frac{n \pi(x-a)}{(b-a)}\right), n=1,2, \ldots, N \\
& v_{n}(x)=\sigma \cos \left(\frac{n \pi(x-a)}{(b-a)}\right), n=0,1,2, \ldots, N-1\left\{\begin{array}{l}
\sigma=\sqrt{\frac{1}{b-a}}, n=0 \\
\sigma=\sqrt{\frac{2}{b-a}}, n \neq 0
\end{array}\right.
\end{aligned}
$$

See e.g. Colbert \& Miller, JCP 96, 1982 (1992)

## Vibrational basis sets

- Why are the $\mu_{\mathrm{n}}(\mathrm{x})$ and $v_{\mathrm{n}}(\mathrm{x})$ functions popular?
$\checkmark$ They yield analytic expressions for $\left\langle\mu_{m}\right| T \mid \mu_{n}>$ and $\left\langle v_{m}\right| T \mid v_{n}>$ for finite and infinite definition intervals [a,b]
$\checkmark$ They are associated with an equidistant quadrature grid (relation to Chebychev)
$\checkmark$ The quadrature rule is of Gaussian accuracy (discrete orthogonality)

$$
\int_{a}^{b} f(x) d x=w \sum_{k=1}^{N} f\left(x_{k}\right)\left\{\begin{array}{l}
w=\frac{b-a}{N+1}, \text { for } \mu_{n} \\
w=\frac{b-a}{N}, \text { for } v_{n}
\end{array}\right.
$$

- Which basis sets are appropriate for bending motion?
$\checkmark$ The bending kinetic energy operator is:

$$
\hat{T}_{\text {bend }}=-c_{\text {bend }}\left(\frac{\partial^{2}}{\partial x^{2}}+\cot (x) \frac{\partial}{\partial x}\right) \quad c_{\text {bend }}=\frac{1}{2 \Theta}
$$

## Legendre basis for bending motion

$\checkmark$ In this form, $T_{\text {bend }}$ is hermitian on $[0, \pi]$ with respect to volume element $\sin (x) d x$
$\checkmark \mu_{n}(x)$ and $v_{n}(x)$ functions perform badly as basis functions for $T_{\text {bend }}$
$\checkmark$ The standard basis functions for $T_{\text {bend }}$ are derived from Legendre polynomials $P_{\text {I }}(x)$ :

$$
\begin{array}{ll}
\sigma_{0} P_{0}(\cos (x))=\sqrt{\frac{1}{2}} & \sigma_{2} P_{2}(\cos (x))=\sqrt{\frac{5}{2^{5}}}(3 \cos (x)+5 \cos (3 x)) \\
\sigma_{1} P_{1}(\cos (x))=\sqrt{\frac{3}{2}} \cos (x) & \sigma_{3} P_{3}(\cos (x))=\sqrt{\frac{9}{2^{13}}}(9+20 \cos (2 x)+35 \cos (4 x))
\end{array}
$$

$\checkmark P_{1}(\cos (x))$ are the eigenfunctions of $T_{\text {bend }} \rightarrow$ diagonal analytic representation
$\checkmark \quad P_{1}(\cos (x))$ are associated with quadrature rule of Gaussian accuracy
$\checkmark$ Grid point density increases moderately towards interval limits

- $\quad P_{\text {}}(\cos (x))$ are suitable bending basis functions for harmonic type potential functions $\rightarrow$ performance good because the density of excited state wave functions accumulates at interval borders

Normalized Legendre functions $P_{l}(\cos (x))$


## Part II

A new basis set for angular motion \& comparison with Legendre functions

## The $\eta_{\mathrm{n}}(\mathrm{x})$ angular basis functions

- Can we formulate basis functions for bending motion that are analog to the $\mu_{n}(x)$ and $v_{n}(x)$ functions?
$\checkmark$ How about:

$$
\eta_{n}(x)=\sqrt{\frac{2}{\pi}} \frac{\sin (n x)}{\sqrt{\sin (x)}}, n=1,2, \ldots, N
$$

- Properties of $\eta_{\mathrm{n}}(x)$ functions:
$\checkmark$ they are orthonormal on $[0, \pi]$ wrt to volume element $\sin (x) \mathrm{dx}$
$\checkmark$ the matrix elements $<\eta_{m}\left|T_{\text {bend }}\right| \eta_{n}>$ have simple analytic solutions
$\checkmark$ they are related to an equidistant quadrature grid
$\checkmark$ the quadrature rule

$$
\int_{0}^{\pi} f(x) \sin (x) d x=\sum_{k=1}^{N} w_{k} f\left(x_{k}\right) \quad w_{k}=\sqrt{\frac{\pi}{N+1}} \sin \left(\frac{\pi k}{N+1}\right)
$$

is of Gaussian accuracy
$-1$
$\eta_{n}(x)$ functions

- 2
---- 3
---- 4



## Definition of model Hamiltonian

- We compare the performance of $\eta_{n}(x)$ and $P_{1}(\cos (x))$ basis functions
- Model system: pure bending motion of $\mathrm{H}_{2} \mathrm{O}$
$\hat{H}=\hat{T}_{\text {bend }}+c_{0}+c_{1}(\cos (x))+c_{2}(\cos (x))^{2}+c_{3}(\cos (x))^{3}+c_{4}(\cos (x))^{4}$
relatively harmonic potential $\rightarrow$ well suited for Legendre basis
- Variational Basis Represention (VBR) for H
$\left\langle\varphi_{m}(x)\right| \hat{O}\left|\varphi_{n}(x)\right\rangle=\int_{a}^{b} \varphi_{m}(x)\left[\hat{O} \varphi_{n}(x)\right] w(x) d x$
$\checkmark$ For a true VBR, all matrix elements must be evaluated exactly

Legendre VBR for $\mathrm{H}_{2} \mathrm{O}$ bending

$\eta_{n}(x)$ VBR for $\mathrm{H}_{2} \mathrm{O}$ bending


## How to improve performance of $\eta_{\mathrm{n}}(\mathrm{x})$ ?

- Obviously, a basis formed exclusively by $\eta_{\mathrm{n}}(\mathrm{x})$ functions is incomplete
- Can we use the complementary functions?

$$
\sigma_{n} \frac{\cos (n x)}{\sqrt{\sin (x)}}, n=0,1,2 \ldots, N-1
$$

- Can we supplement the $\eta_{n}(x)$ functions? For example:

$$
\begin{gathered}
s_{r}(x)=\exp (-r \sin (x)) \quad s_{t}(x)=\exp (-t \sin (x)) \\
\left\{\eta_{m}(x)\left(1-s_{r}(x)\right), P_{n}(\cos (x)) s_{t}(x)\right\}
\end{gathered}
$$

$\checkmark$ Switching functions $\left(1-\mathrm{s}_{\mathrm{r}}(\mathrm{x})\right), \mathrm{s}_{\mathrm{t}}(\mathrm{x})$ keep basis orthogonal
$\checkmark$ Evaluation of matrix elements tedious

## Supplementation of $\eta_{\mathrm{n}}(\mathrm{x})$ functions

- Is direct basis extension an option? For example:
$\left\{\begin{array}{l}\eta_{1}(x), \eta_{2}(x), \ldots, \eta_{N}(x), \\ \sigma_{0} P_{0}(\cos (x)), \sigma_{1} P_{1}(\cos (x)), \ldots, \sigma_{M} P_{M}(\cos (x))\end{array}\right\}$
$\checkmark$ We construct representation of H in this mixed basis $\rightarrow$ symmetric matrix $\boldsymbol{A}$
$\checkmark$ Loewdin (symmetric) orthogonalization of mixed basis:
- Diagonalization of overlap matrix yields eigenvector matrix $\boldsymbol{U}$ and the matrix $X=\operatorname{diag}\left(1 / \sqrt{ } \varepsilon_{1}, 1 / \sqrt{ } \varepsilon_{2}, \ldots, 1 / \sqrt{ } \varepsilon_{N}\right)$
- Orthogonalization of mixed basis according to:

$$
\left(U X U^{T}\right) A\left(U X U^{T}\right)=H
$$

$\checkmark \boldsymbol{H}$ is the desired representation of H in orthonormal mixed basis

Mixed basis (+ $\left.\mathrm{P}_{0}, \mathrm{P}_{1}\right)$ VBR $\mathrm{H}_{2} \mathrm{O}$ bending


## Part III

## Overview of localized \& delocalized representations, construction of localized mixed basis functions, applications

## Overview: representations

- We differentiate between infinitely localized, nearly localized and delocalized basis functions
- Discrete representations are only possible for local operators
- Discrete Variable Representation (DVR) of local operator O is a matrix diagonal over the grid points $x_{k}$.
- If the grid is related to orthogonal basis functions $\varphi_{m}(x)$ through a quadrature rule of the form:

$$
\int_{a}^{b} f(x) w(x) d x=\sum_{k=1}^{N} w_{k} f\left(x_{k}\right)
$$

then we can define a Finite Basis Representation (FBR):

$$
\int_{a}^{b} \varphi_{m}(x)\left[\hat{O} \varphi_{n}(x)\right] w(x) d x \approx \sum_{k=1}^{N} \varphi_{m}\left(x_{k}\right)\left[\hat{O} \varphi_{n}\left(x_{k}\right)\right] w_{k}
$$

## VBR, FBR, DVR, NDVR

- FBR matrix is an approximation to the VBR matrix
- FBR and DVR are equivalent:
- In analogy we can define:

$$
\begin{aligned}
& { }^{F B R} O=\Lambda^{D V R} O \Lambda^{t} \\
& { }^{V B R} O=\Lambda^{N D V R} O \Lambda^{t}
\end{aligned}
$$

- NDVR matrix is an approximation to the DVR matrix:
$\checkmark$ NDVR basis functions are approximately localized at grid points $\mathrm{X}_{\mathrm{k}}$.
- The term "DVR calculation" is not exact: $\quad{ }^{D V R} H={ }^{N D V R} T+{ }^{D V R} V$
$\checkmark$ DVR results are not variational!
- How to perform "DVR calculation" for mixed basis functions $Q_{n}(x)$ ?
$\checkmark$ Definition of grid points through zeroes of $Q_{N+M+1}(x)\left(\right.$ from $\eta_{N+1}(x)$ and $\left.P_{M}(\cos (x))\right)$ $\rightarrow$ explicit derivation of orthogonal mixed basis
$\checkmark$ Establishment of quadrature rule for mixed basis $\rightarrow$ derivation of $\Lambda$ matrix



## $Q_{m}(x)$ NDVR localized at $x_{k}, k=1,2,3,6$ from $\eta_{n}(x)(n=1-30), P_{l}(\cos (x))(I=1-2)$


$V_{1}=300+3000 \cos (x)+8000(\cos (x))^{2}-3000(\cos (x))^{3}-2000(\cos (x))^{4}$

$$
c_{\text {bend }}=10.0
$$



Legendre VBR for $\mathrm{V}_{1}$

$Q_{m}(x)$ VBR for $V_{1}$ from $\eta_{n}(x)(n=1-N-2), P_{1}(\cos (x))(l=1-2)$


Legendre VBR for $\mathrm{V}_{1}$


$$
Q_{m}(x) \text { VBR for } V_{1} \text { from } \eta_{n}(x)(n=1-N-2), P_{1}(\cos (x))(l=1-2)
$$



Legendre DVR for $\mathrm{V}_{1}$


## $Q_{m}(x)$ DVR for $V_{1}$ from $\eta_{n}(x)(n=1-N-2), P_{1}(\cos (x))(l=1-2)$



Legendre DVR for $\mathrm{V}_{1}$


## $Q_{m}(x)$ DVR for $V_{1}$ from $\eta_{n}(x)(n=1-N-2), P_{1}(\cos (x))(l=1-2)$




Legendre VBR for $\mathrm{V}_{2}$

$Q_{m}(x)$ VBR for $V_{2}$ from $\eta_{n}(x)(n=1-N-2), P_{1}(\cos (x))(I=1-2)$


Legendre VBR for $\mathrm{V}_{2}$

$Q_{m}(x)$ VBR for $V_{2}$ from $\eta_{n}(x)(n=1-N-2), P_{1}(\cos (x))(I=1-2)$


Legendre DVR for $\mathrm{V}_{2}$


## $Q_{m}(x)$ DVR for $V_{2}$ from $\eta_{n}(x)(n=1-N-2), P_{l}(\cos (x))(I=1-2)$



Legendre DVR for $\mathrm{V}_{2}$

$Q_{m}(x)$ DVR for $V_{2}$ from $\eta_{n}(x)(n=1-N-2), P_{1}(\cos (x))(I=1-2)$


## Part IV

## Conclusions

- Legendre functions form for many applications a good basis set for bending degrees of freedom
- However, they offer a limited flexibility in particular for the description of states with larger density close to the center of the interval
- The combination of $\eta_{n}(x)$ and $P_{1}(\cos (x))$ functions appears to be an interesting alternative to the pure Legendre basis
- For mixed basis VBR calculations:
$\checkmark$ all matrix elements can be evaluated analytically
$\checkmark$ computationally efficient because of Loewdin orthogonalization
$\checkmark$ more homogeneous accuracy distribution for different eigenstates
- For mixed basis DVR calculations:
$\checkmark$ explicit orthogonalization complicated $\rightarrow$ but needs to be performed only once since kinetic energy operator is always the same
$\checkmark$ quadrature rule for mixed basis set has been derived
$\checkmark$ accuracy similar to Legendre DVR can be reached, but further improvement necessary


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