



Precise Measurements of the Casimir Force: Experimental Details

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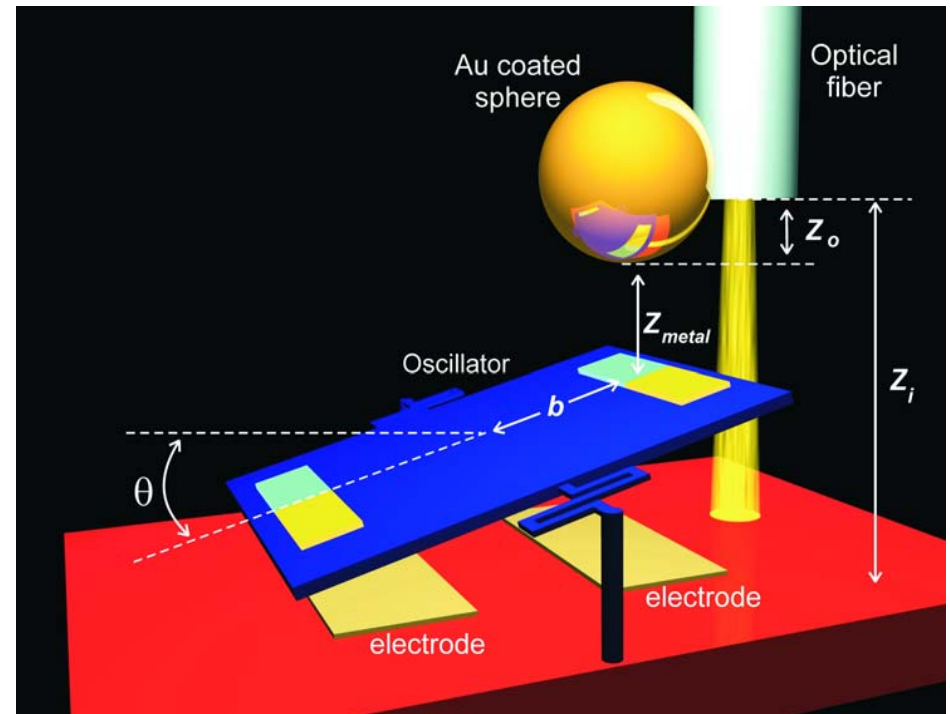
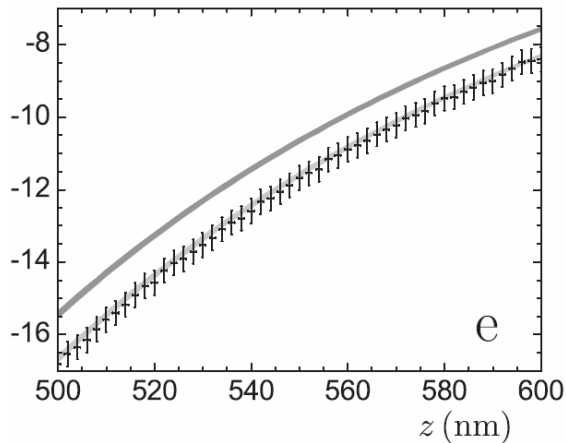
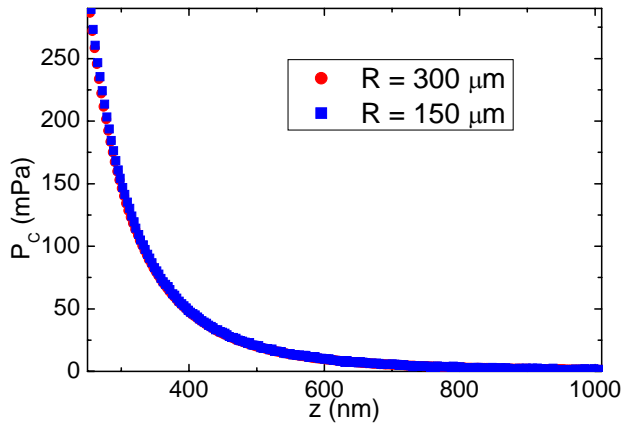
Los Alamos National Lab.

Funding

NSF, DOE, LANL, DARPA

Motivation

- Precise measurements of the Casimir force
 - Background for hypothetical forces
 - Allows for comparison with theory
 - Temperature dependence and effects on nanosystems



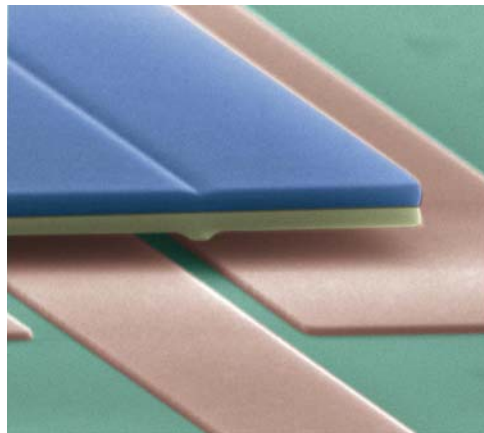


Outline

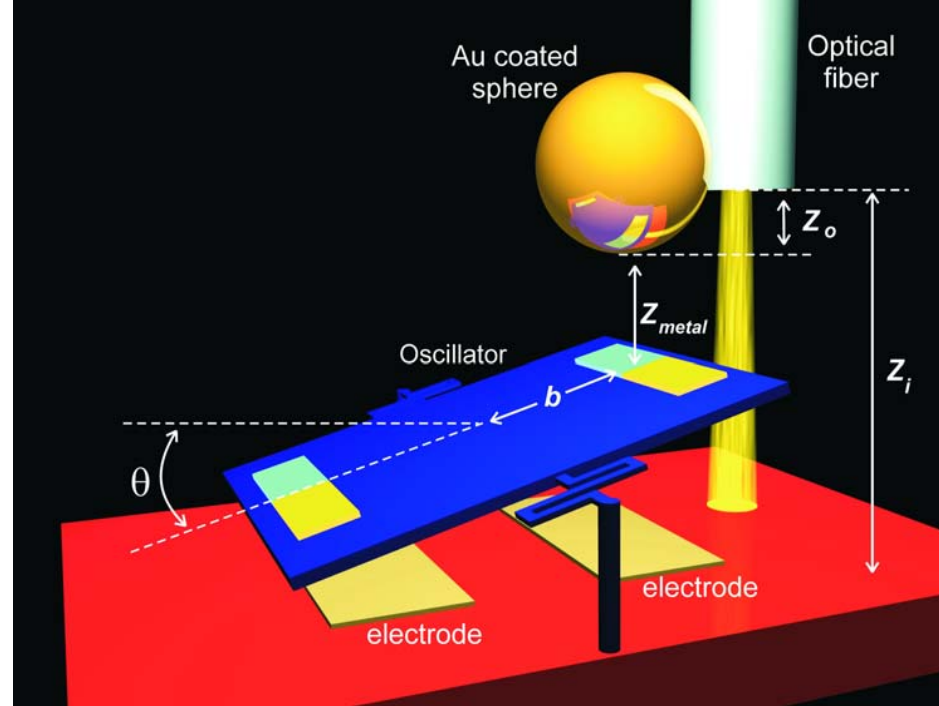
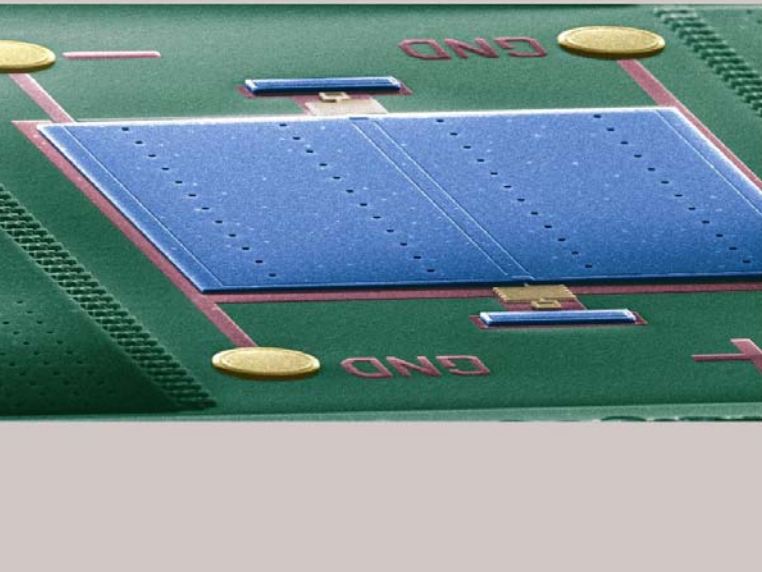
- Experimental setup, sample preparation, and characterization
- Measurement of the interaction
- Measurement of the separation dependence
- Comparison with theory
- Proposed measurements to see the effects of geometry
- Summary



Experimental setup



$$z_{metal} = z_i - z_o - z_g - b \Theta$$



In MEMS plate as 500 um x 500 um, have changed the setup, the sphere is on the oscillator, the plate is on top
Plate thickness: 3.5 um

Spring lengths: 500 um

Larger sample, requires different deposition.

$K_{torsion} \sim 10^{-10}$ Nm/rad

Sphere radii: 10 um - 150 um
Measurements done in vacuum at room temperature, in an oil-free chamber.

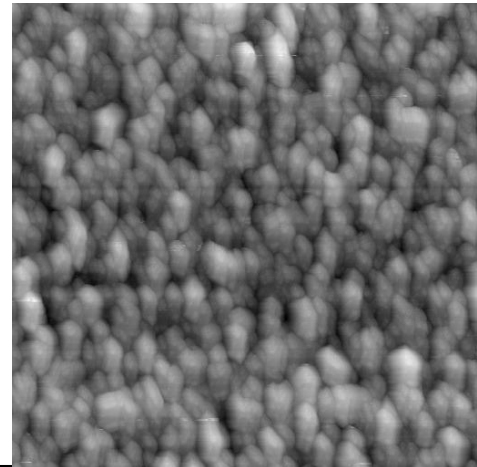
Resonance frequency ~ 1000 Hz

Quality factor ~ 10000 (@ 10^{-6} Torr)

Sample preparation and characterization

- Au on the sapphire sphere is deposited by thermal evaporation.
- Au on the oscillator is also deposited by thermal evaporation
- In new, larger samples it is deposited by electroplating (on Si[111])
- Samples are characterized by measuring resistance as a function of temperature, AFM measurements and also ellipsometry in the electrodeposited sample.
- The sample is mounted into the system, baked to ~ 60 °C for $\frac{1}{2}$ hour.

($10 \times 10 \mu\text{m}^2$)
 $\sim 20 \text{ nm}_{\text{pp}}$

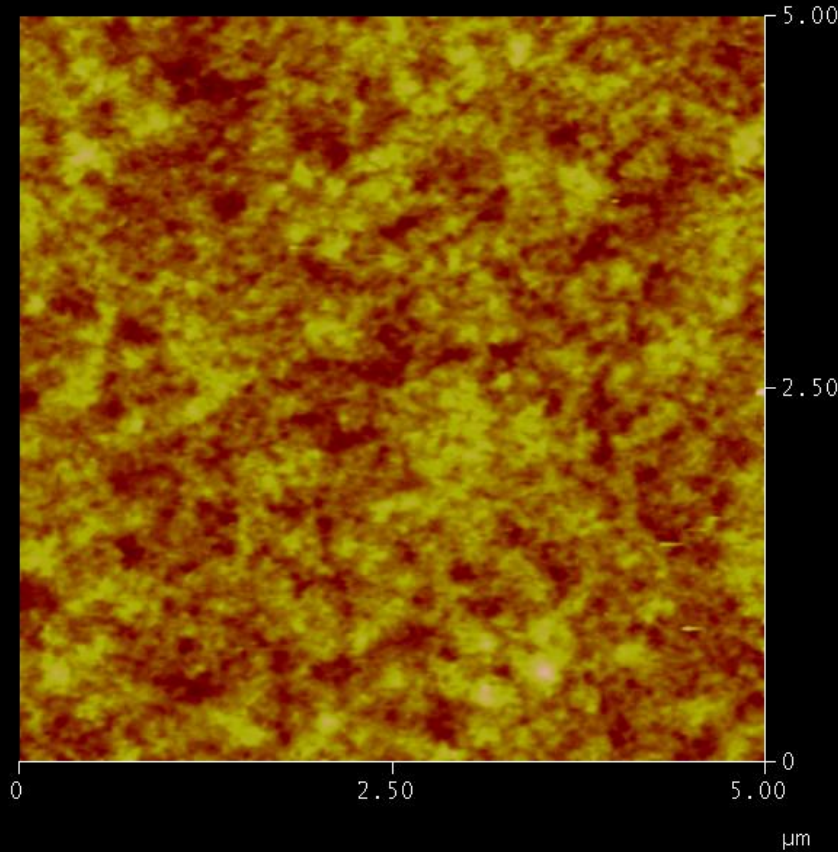




AFM measurements

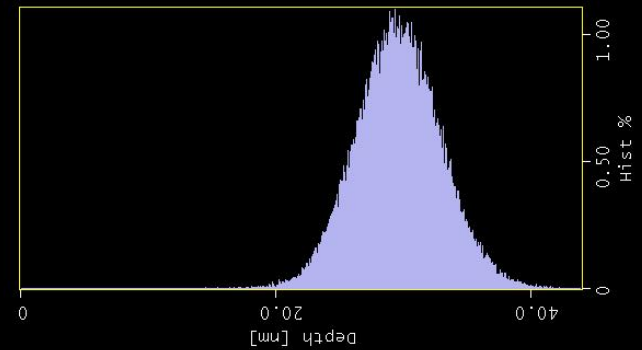
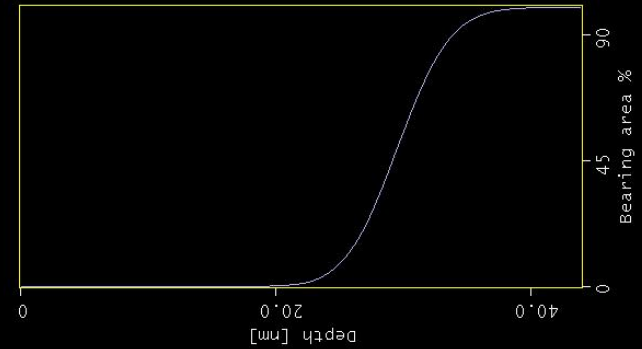
Clear Execute Undo

Flatten



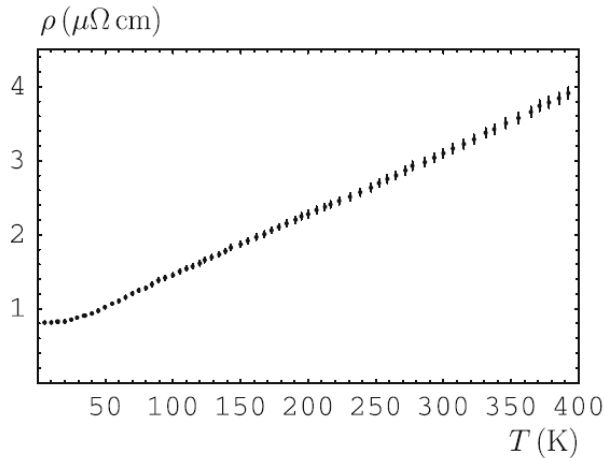
electro-au-si.000

Histogram





Resistivity and spectroscopic ellipsometry



-R vs T and spectroscopic ellipsometry (190 nm to 830 nm) used to determine ω_p .

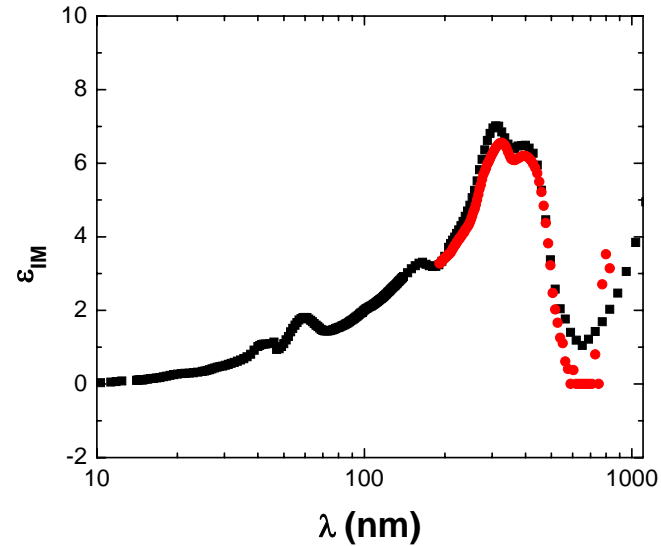
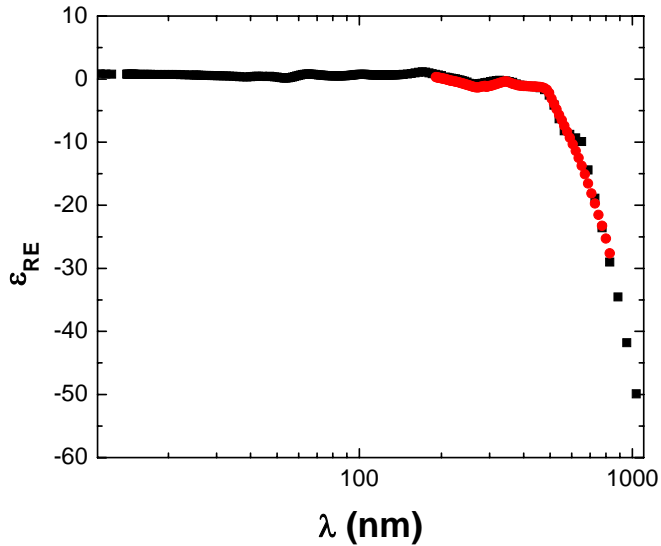
-Both methods indicate a rather good Au sample

$$\rho(T) = \frac{4\pi}{\omega_p^2 \cdot \tau(T)} \propto \frac{T}{\omega_p^{3/2}}$$

$$\omega_p = (8.9 \pm 0.1) \text{ eV}$$



Resistivity and spectroscopic ellipsometry



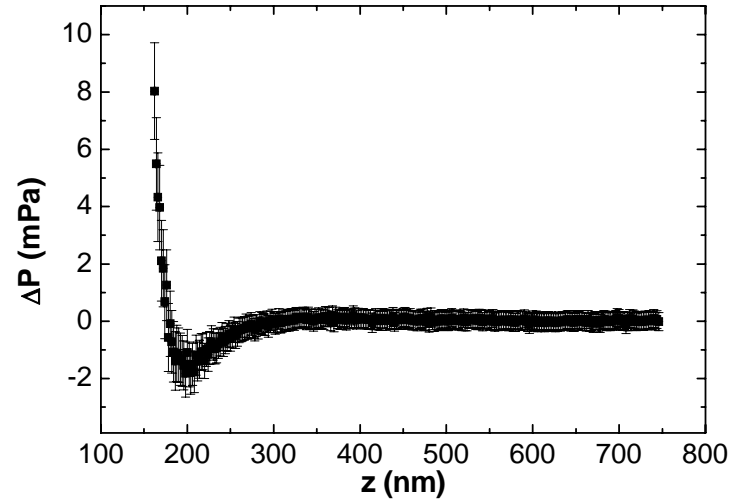
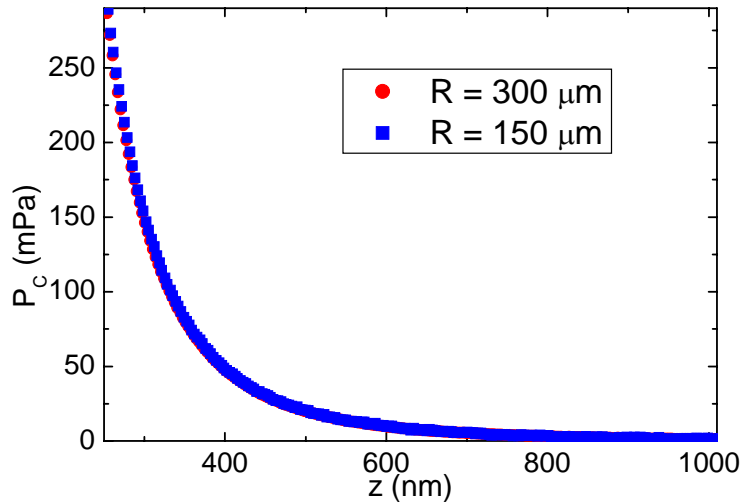
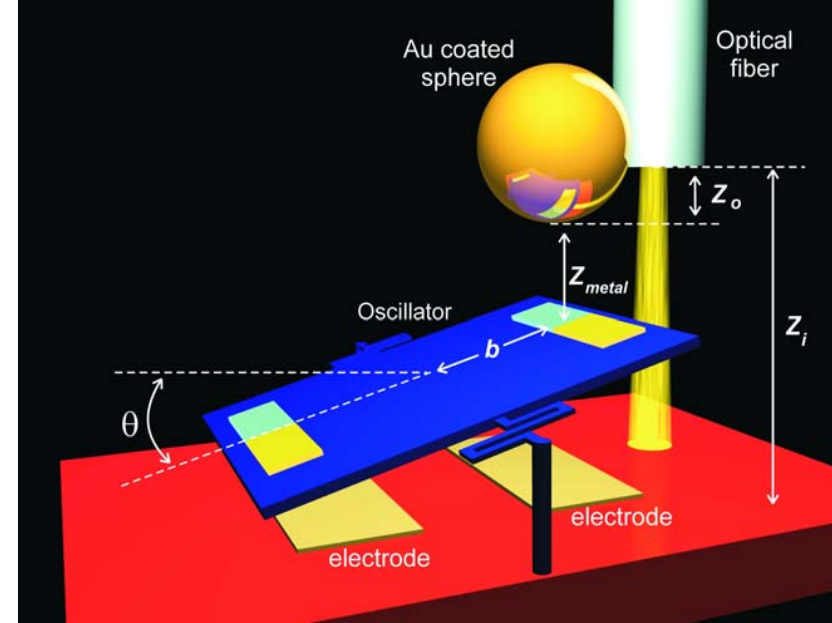
-Measured real and imaginary parts of the dielectric functions (red circles) are similar to published values (Palik, black squares)

-It was checked that either can be used, giving similar results. Palik values are used on the rest of this presentation.

Pressure measurements

$$\omega_r^2 = \omega_o^2 \left(1 - \frac{b^2}{I\omega_o^2} \frac{\partial F_C}{\partial z} \right)$$

$$F_C = 2\pi R \times E_C \Rightarrow \frac{\partial F_C}{\partial z} = 2\pi R \times P_C$$





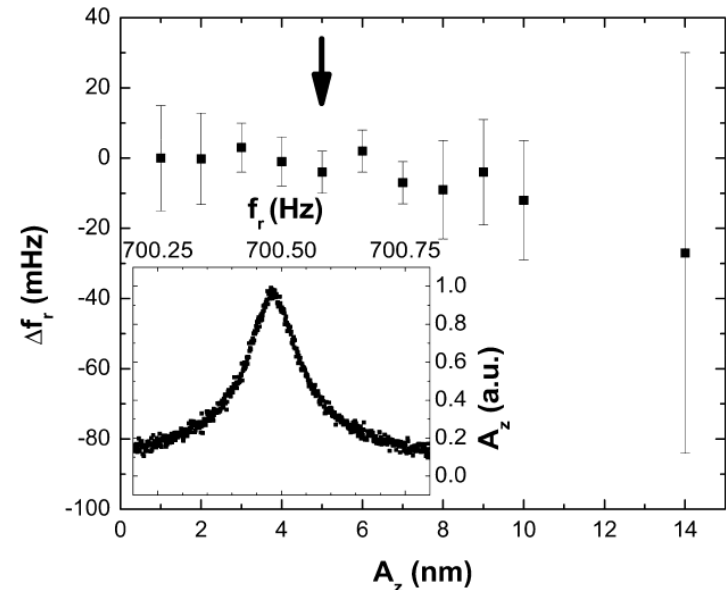
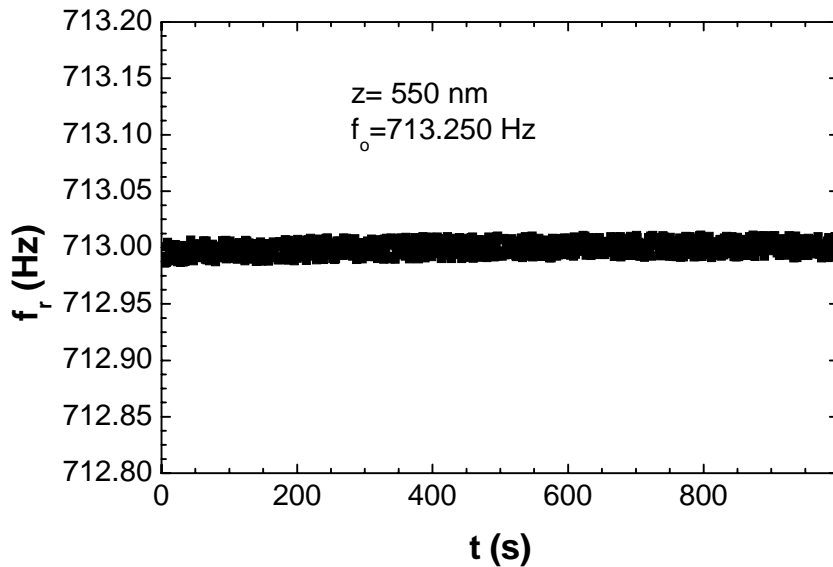
Pressure determination

$$\omega_r^2 = \omega_o^2 \left(1 - \frac{b^2}{I\omega_o^2} \frac{\partial F_C}{\partial z} \right)$$

$$F_C = 2\pi R \times E_C \Rightarrow \frac{\partial F_C}{\partial z} = 2\pi R \times P_C$$

Determined by:

- Looking into the response of the oscillator in the thermal bath
- Inducing a time dependent separation between the plate and the sphere





Pressure measurements

$$\omega_r^2 = \omega_o^2 \left(1 - \frac{b^2}{I\omega_o^2} \frac{\partial F_C}{\partial z} \right)$$

$$F_C = 2\pi R \times E_C \Rightarrow \frac{\partial F_C}{\partial z} = 2\pi R \times P_C$$

Errors

Minimum values

Frequency:

6mHz

~28 mHz (at 750 nm)

R:

0.2 μm

150 μm

b^2/I :

0.0005 μg^{-1}

1.2432 μg^{-1}

Errors:

Random: 0.46 mPa (162 nm)

0.11 mPa (300 nm)

Systematic: 2.12 mPa (162 nm)

0.44 mPa (300 nm)



Separation measurement

z_o is determined using a known interaction

z_i , Θ are measured for each position

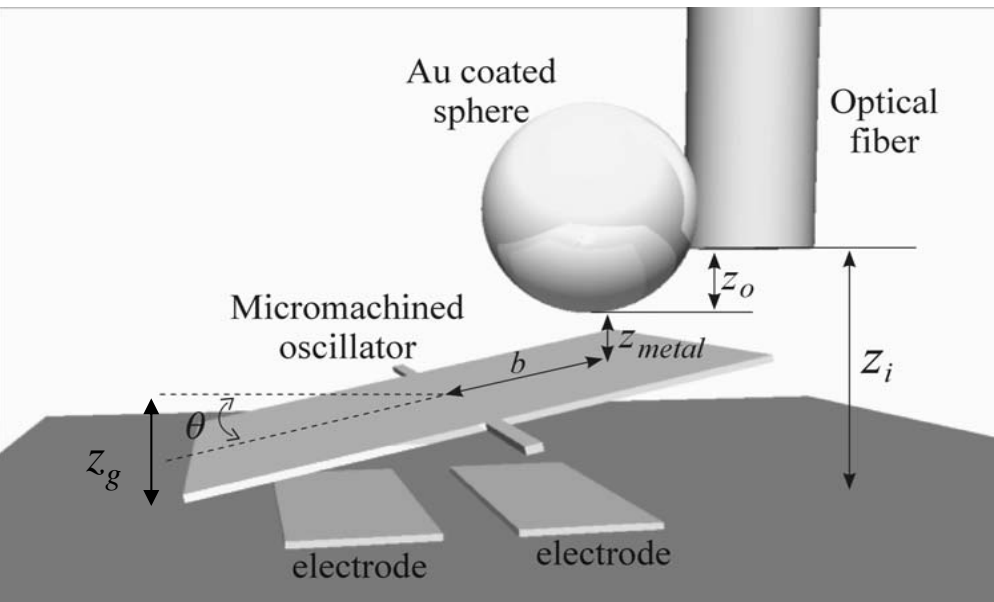
$z_g = (2172.8 \pm 0.1)$ nm, interferometer

$z_i = \sim(12000.0 \pm 0.2)$ absolute interferometer

$z_o = (8162.3 \pm 0.5)$ nm, electrostatic calibration

$b = (207 \pm 2)$ μm , optical microscope

$\Theta = \sim(1.000 \pm 0.001)$ μrad





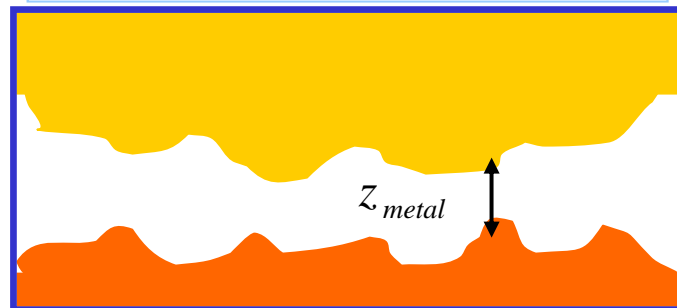
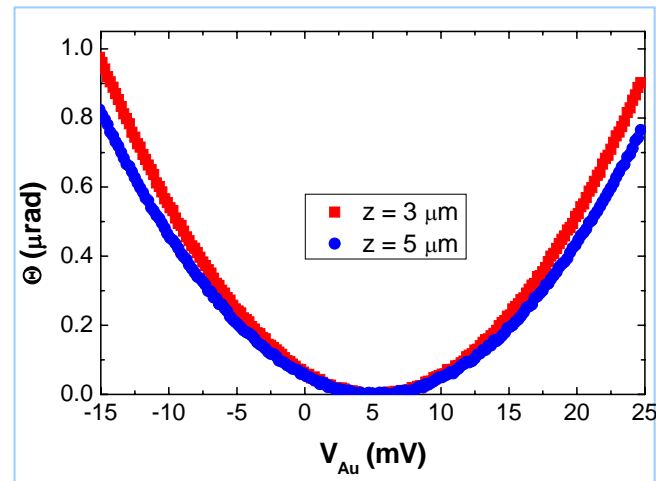
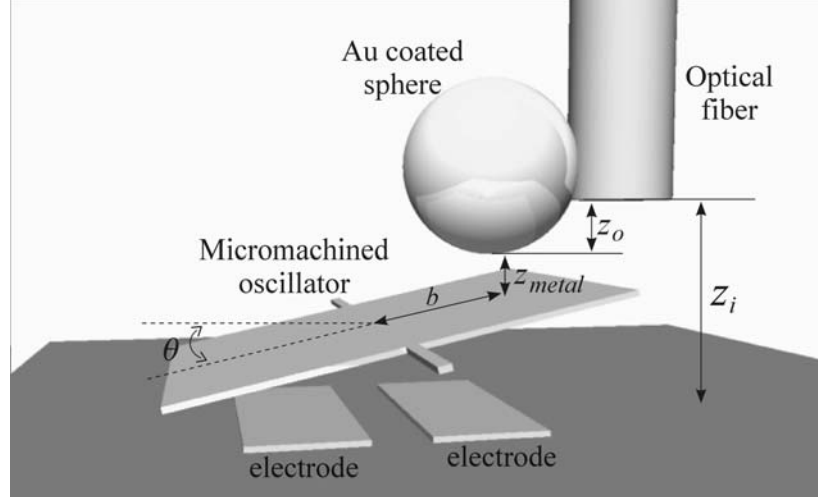
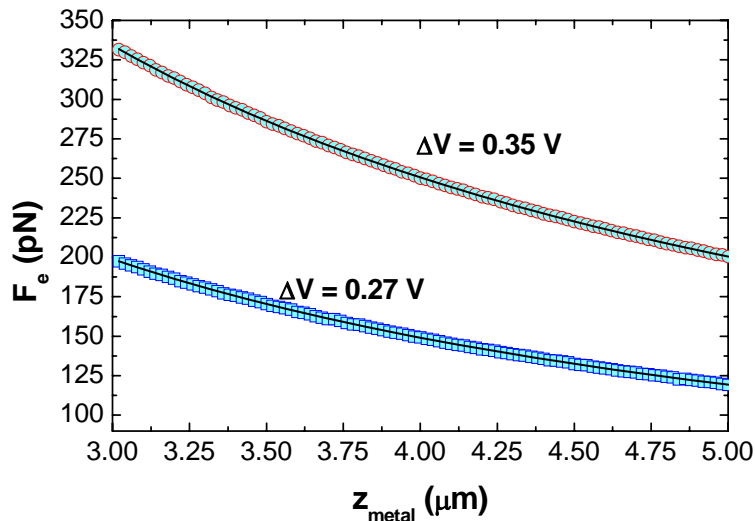
Separation measurement

Electrostatic force calibration

$$F_e = -2\pi\epsilon_o(V - V_{Au})^2 \sum_n \frac{\coth u - n \coth nu}{\sinh nu} \sim$$

$$-2\pi\epsilon_o(V - V_{Au})^2 \sum_0^7 A_i \left(\frac{R}{(z_{metal} + \bar{z})} \right)^{1-i}$$

$$u = 1 + \frac{(z_{metal} + \bar{z})}{R}$$

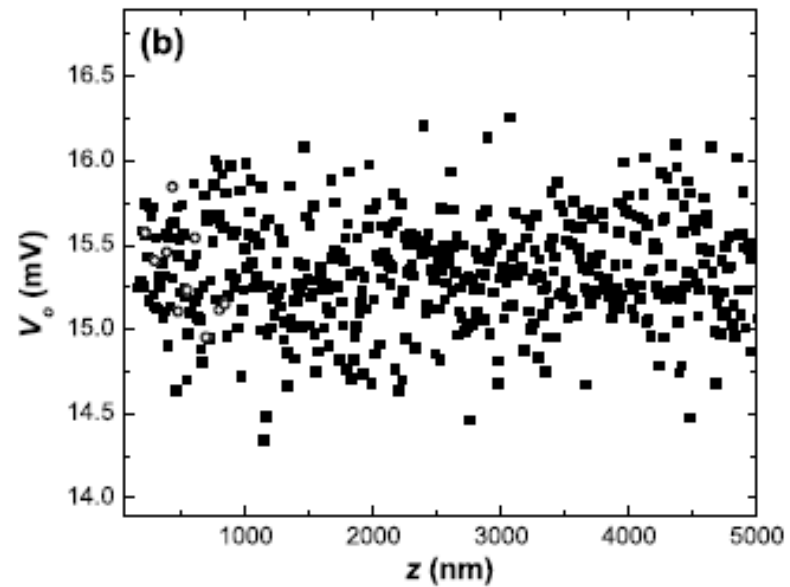
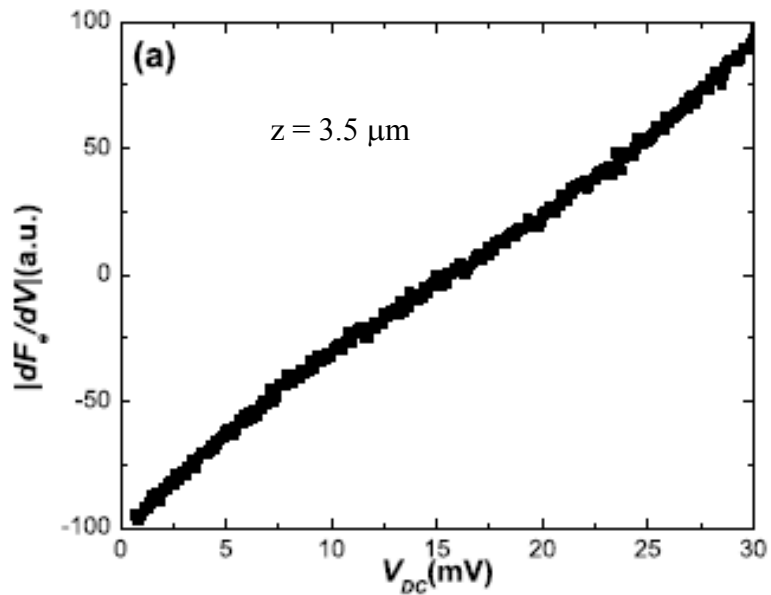
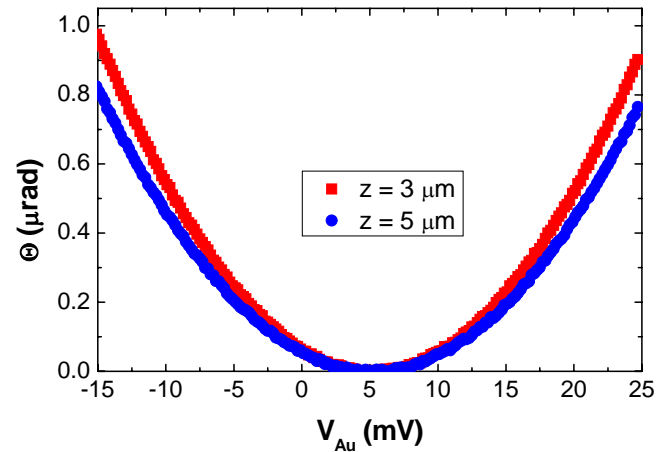




Separation measurement

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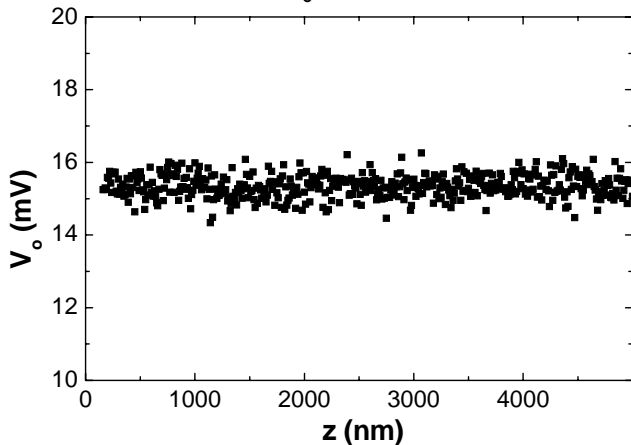
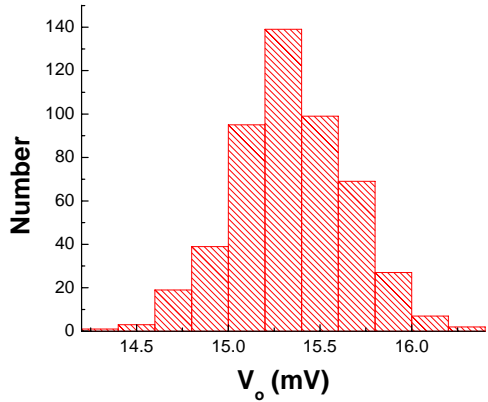
Electrostatic force calibration





Separation measurement

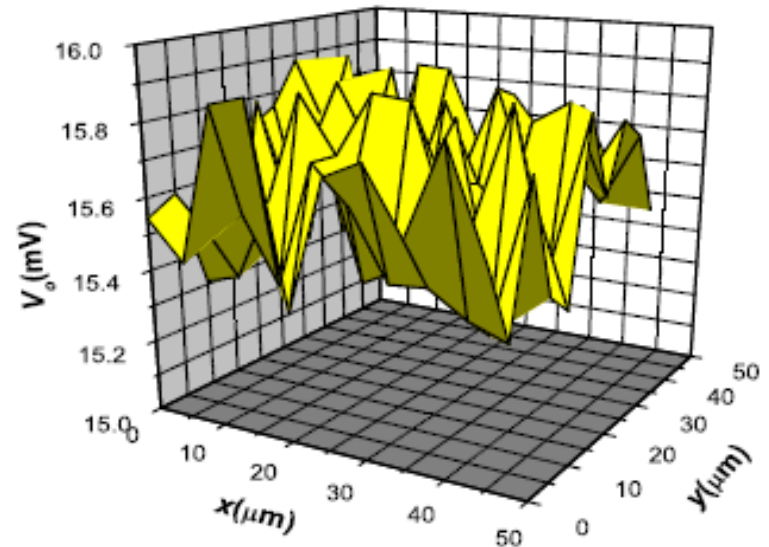
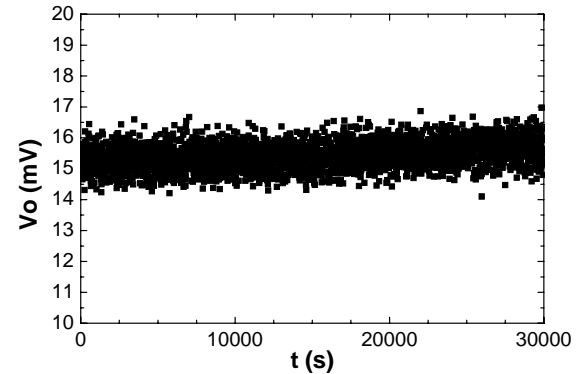
Electrostatic force calibration



V_o **must** be constant as a function of separation...

Otherwise, V_o needs to be determined at each point

... and time

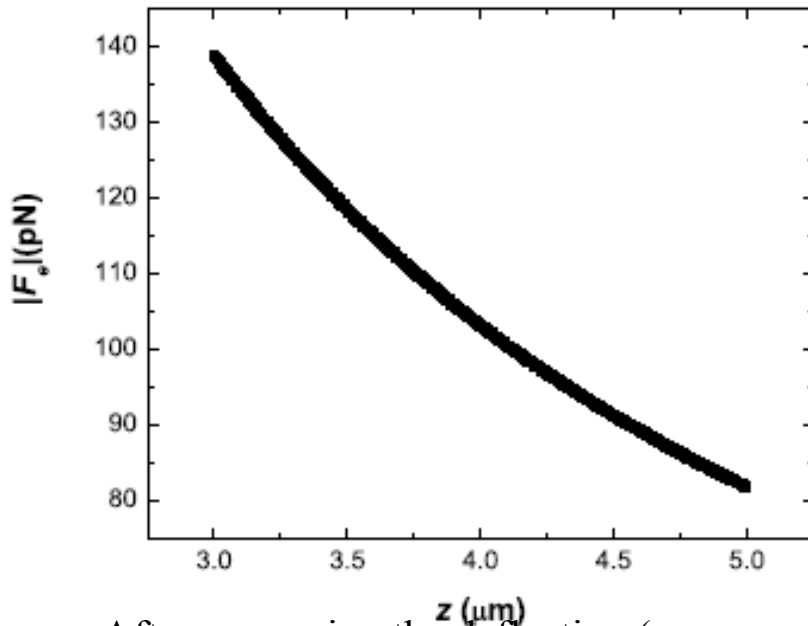


10 x 10 grid, 5 μm pitch



Separation measurement

Electrostatic force calibration

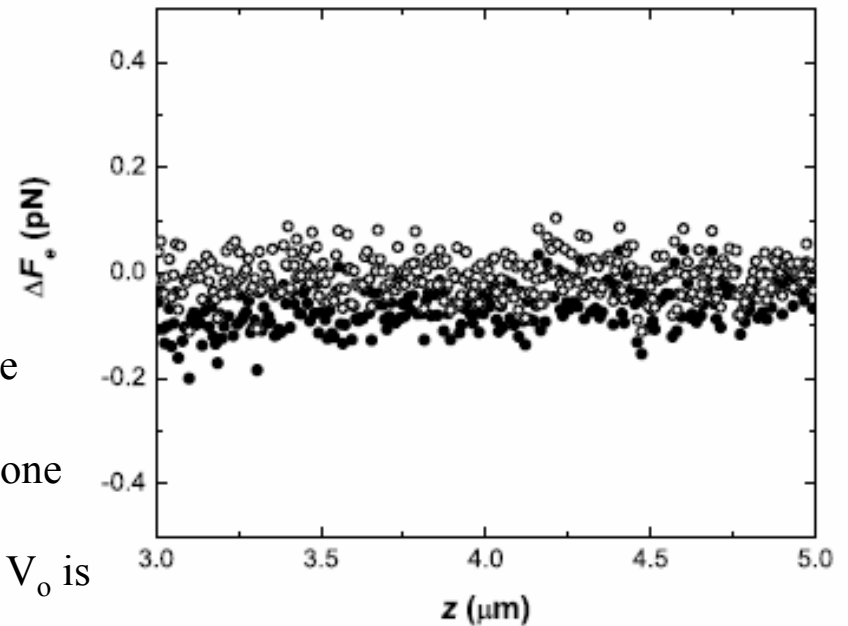


- After measuring the deflection (expressed as force here), we adjust for the unknown separation.
- The figure shows the ΔF_e for the optimal and one off by 1.5 nm
- The error in the force associated with the error in V_0 is < 1 fN.

$$F_e = -2\pi\epsilon_0(V - V_{Au})^2 \sum_n \frac{\coth u - n \coth nu}{\sinh nu} \sim$$

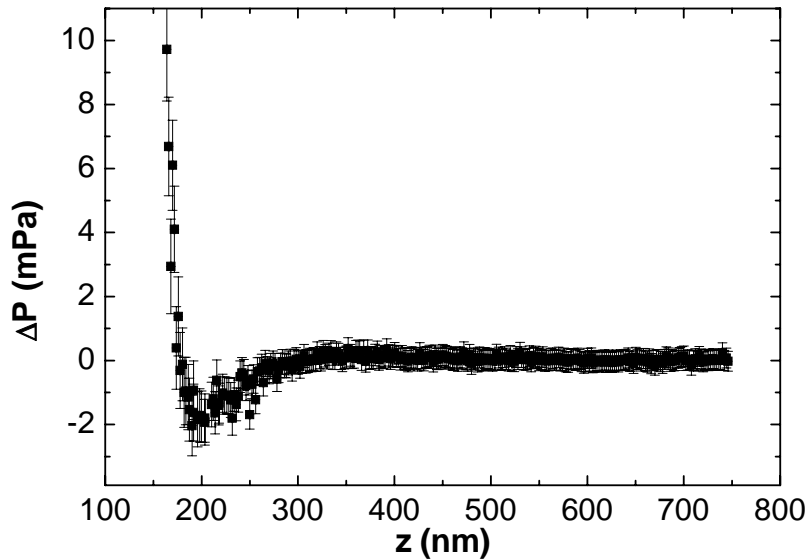
$$-2\pi\epsilon_0(V - V_{Au})^2 \sum_0^7 A_i \left(\frac{R}{(z_{metal} + \bar{z})} \right)^{1-i}$$

$$u = 1 + \frac{(z_{metal} + \bar{z})}{R}$$

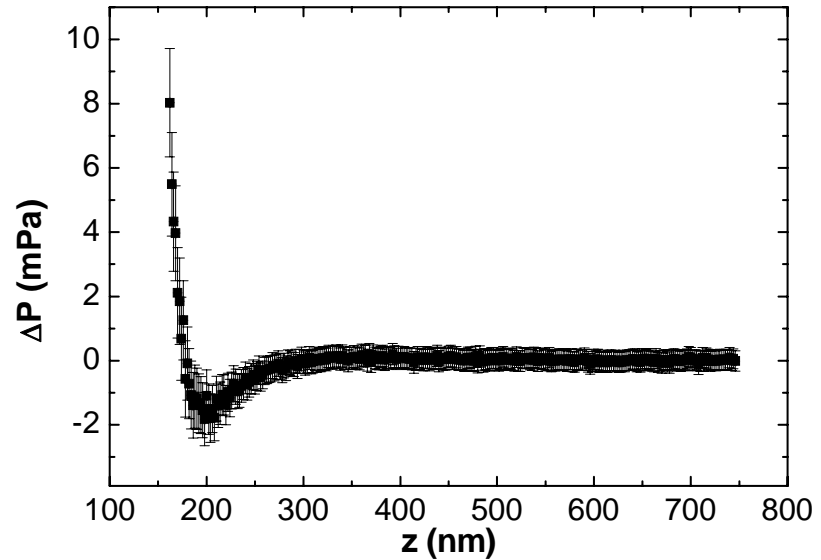




Equivalent P_C measurement



New sample (September 2009)

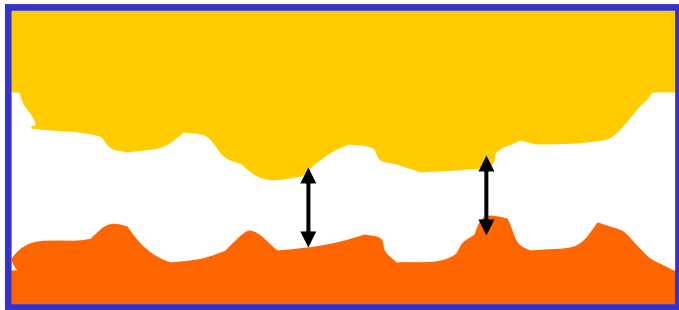


Old data (2005)



Comparison with theory

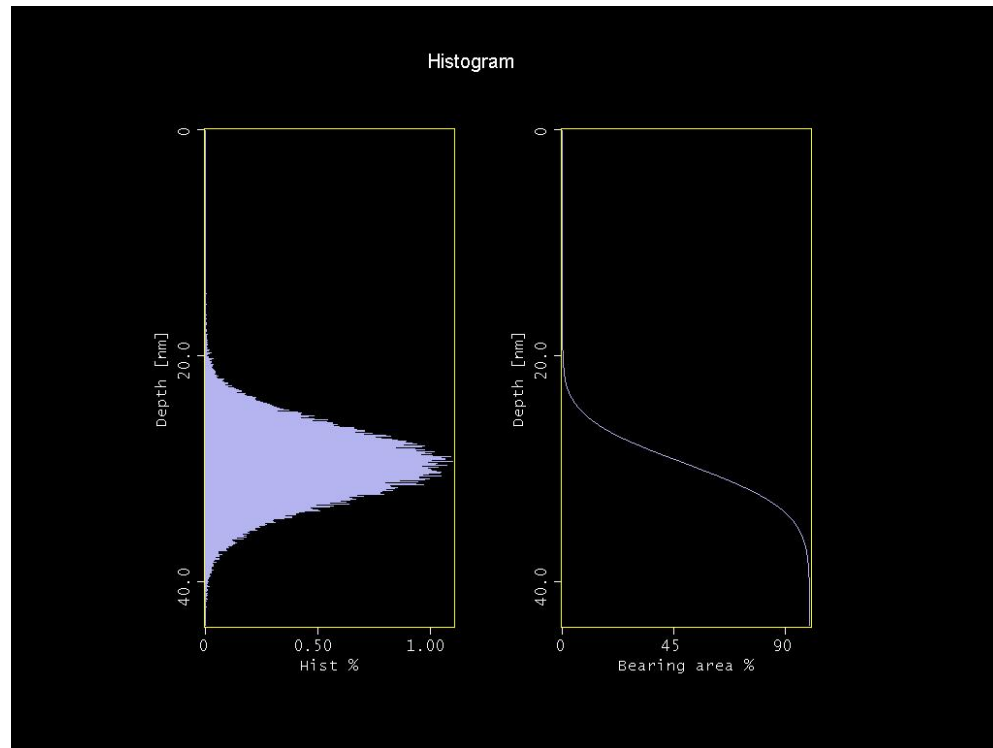
Roughness corrections



$$F_C = \sum_i v_i F_{CS}(z_i)$$

v_i : Fraction of the sample at separation z_i

Roughness corrections are ~0.5% to the Casimir force at 160 nm





Comparison with theory

IUPUI

Finite conductivity and finite temperature

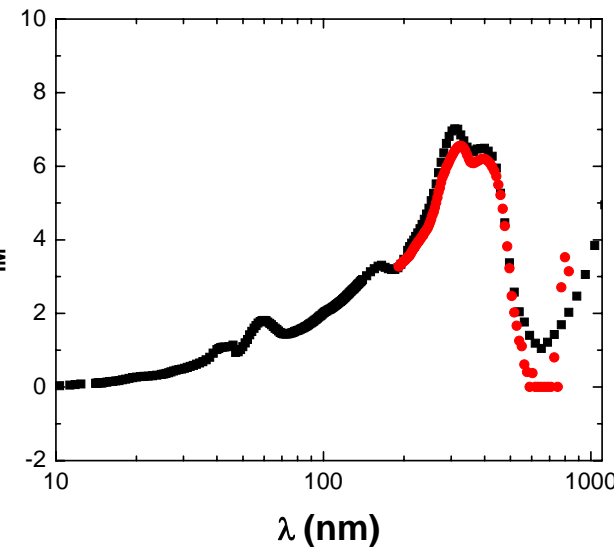
$$P(z) = -\frac{k_B T}{\pi} \sum_{l=0}^{\infty} \int_0^{\infty} k_{\perp} dk_{\perp} q_l \times \left\{ [r_{\parallel}^{-2}(\xi_l, k_{\perp}) e^{2q_l z} - 1]^{-1} + [r_{\perp}^{-2}(\xi_l, k_{\perp}) e^{2q_l z} - 1]^{-1} \right\}$$

$$r_{\parallel, L}^{-2}(\xi_l, k_{\perp}) = \left[\frac{k_l + \varepsilon(i\xi_l) q_l}{k_l - \varepsilon(i\xi_l) q_l} \right]^2, \quad r_{\perp, L}^{-2}(\xi_l, k_{\perp}) = \left[\frac{k_l + q_l}{k_l - q_l} \right]_{\varepsilon_{\text{IM}}}^2$$

$$q_l^2 = k_{\perp}^2 + \left(\frac{4\pi^2 k_B T l}{hc} \right)^2$$

$$\varepsilon(i\xi) = 1 + \frac{2}{\pi} \int_0^{\infty} \frac{\omega \text{Im} \varepsilon(\omega)}{\omega^2 + \xi^2} d\omega$$

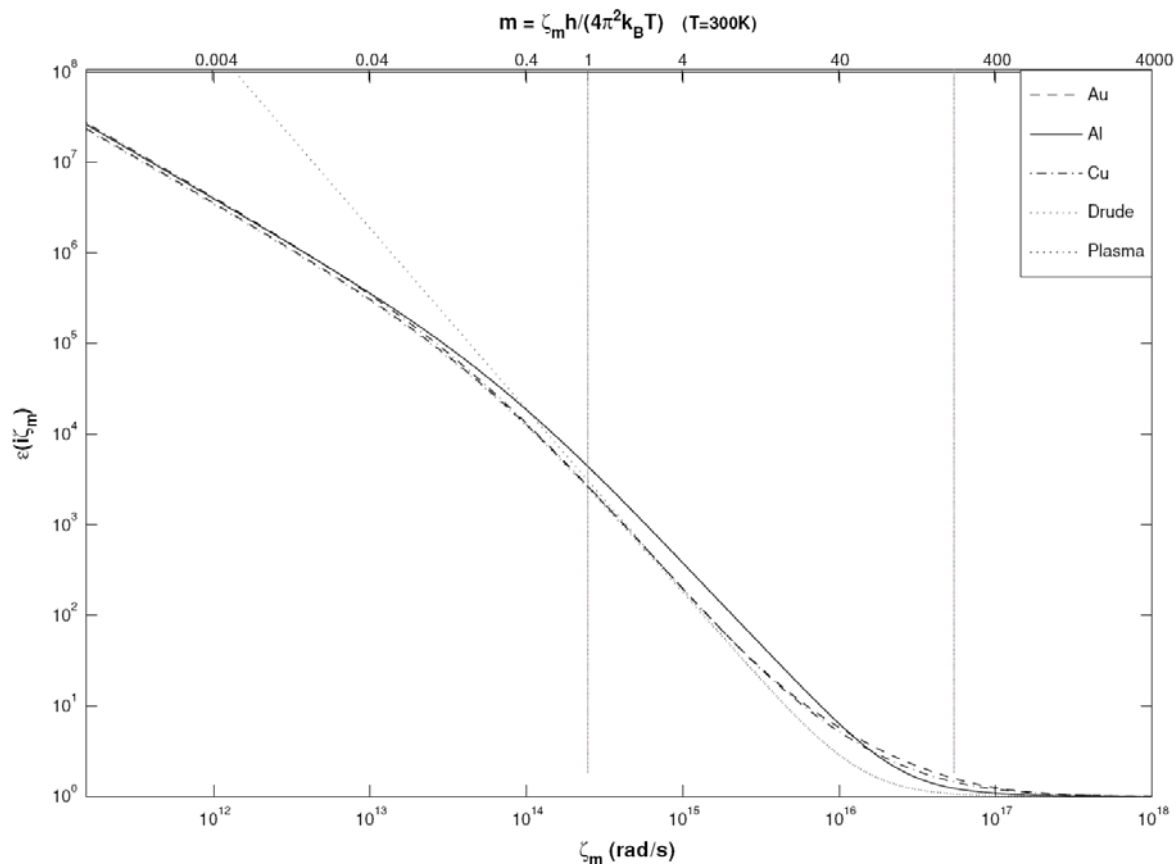
$$k_l^2 = k_{\perp}^2 + \varepsilon(i\xi_l) \xi_l^2 / c^2$$





Comparison with theory

$$\varepsilon(i\xi) = 1 + \frac{2}{\pi} \int_0^\infty \frac{\omega \operatorname{Im} \varepsilon(\omega)}{\omega^2 + \xi^2} d\omega$$



$$\varepsilon(\omega) = 1 - \frac{\omega_p^2}{\omega(\omega + i\gamma)}$$

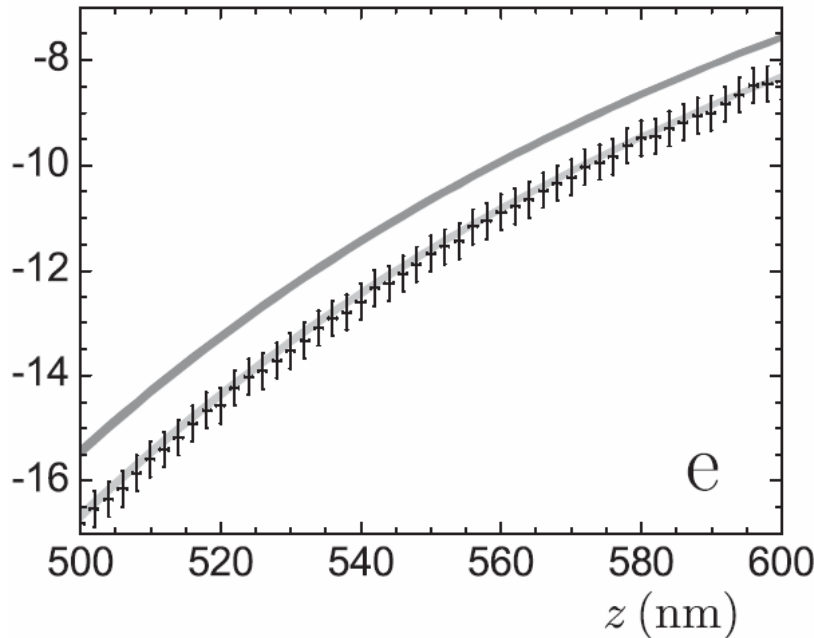
$$\varepsilon(\omega) = 1 - \frac{\omega_p^2}{\omega^2}$$

Bentsen et al., J. Phys. A (2005)



Pressure determination

PRD 75, 077101



There is a significant issue: Drude does not agree with the data

- Experimental problem?
- Theoretical problem?
- Theory not applied to the right experiment?

- Dark grey, Drude model approach
- Light grey, impedance approach

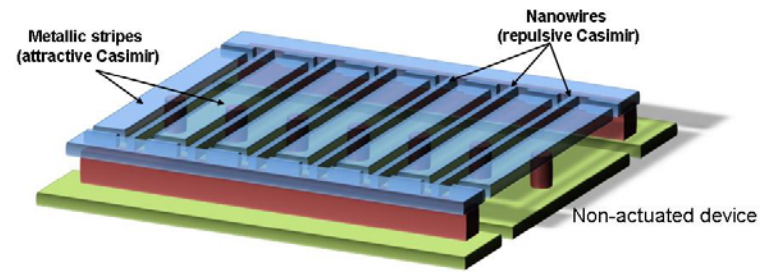
Theoretical errors:

- Sample dependence: 0.5%
- Separation dependence: 1.5% (162 nm)
0.32%(750 nm)

~19 mPa @162 nm (Exp: ~2.5 mPa @162 nm)

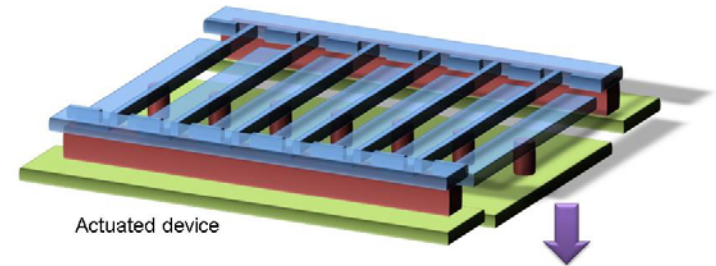
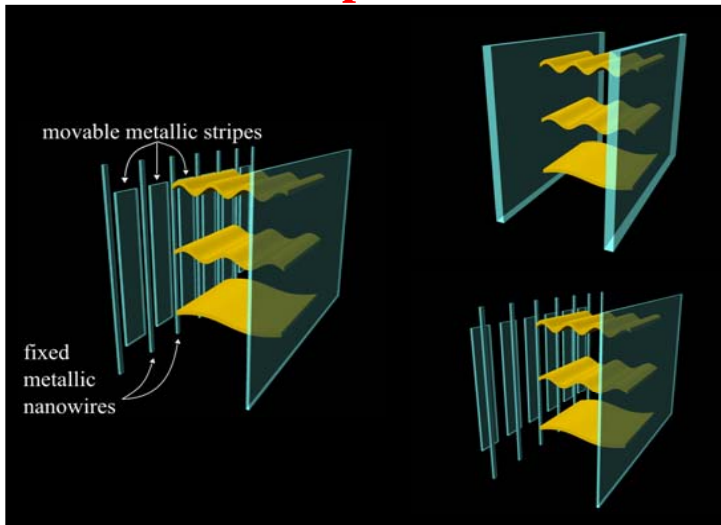


Geometrical effects



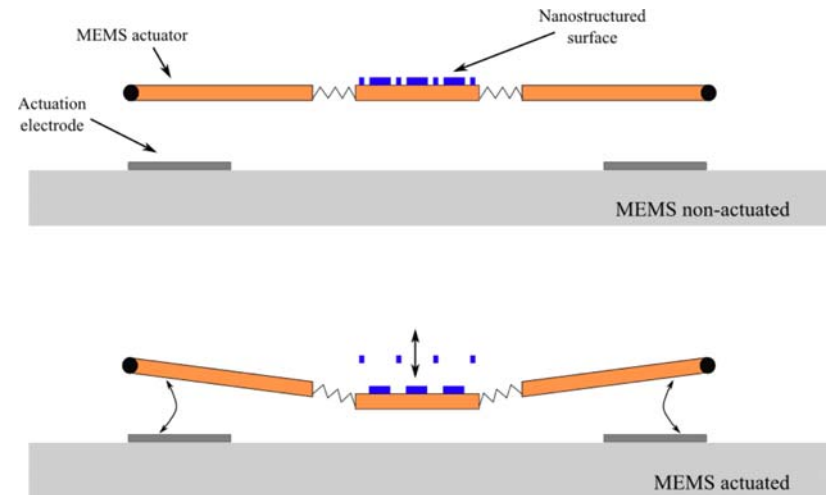
“Role of surface plasmons in the Casimir effect”, F. Intravaia et al., PRA 76 (2007)

Real-time manipulation



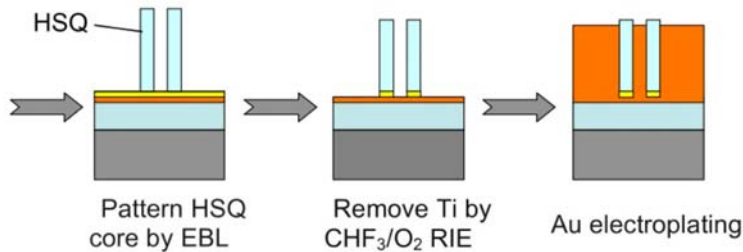
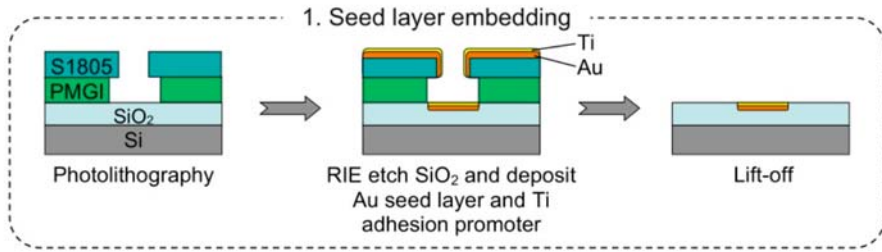
Dynamically deformable nanostructure

- Integration of nanostructure with MEMS
- Displacement $\sim 500\text{nm}$
- Precise control of motion ($\pm 1\text{nm}$)
- Shielded surfaces (fringe fields)

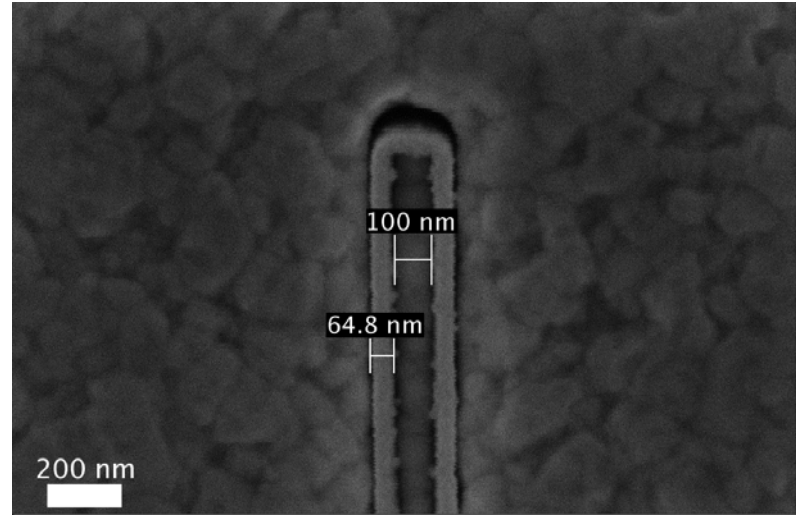
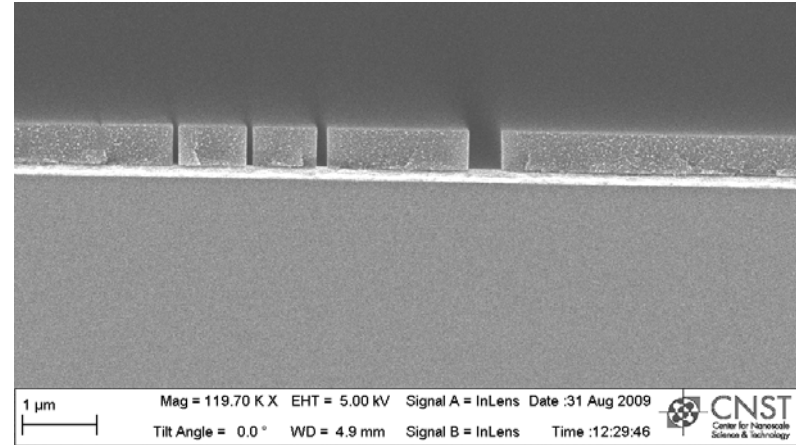


Geometrical effects

Metallic nanostructures

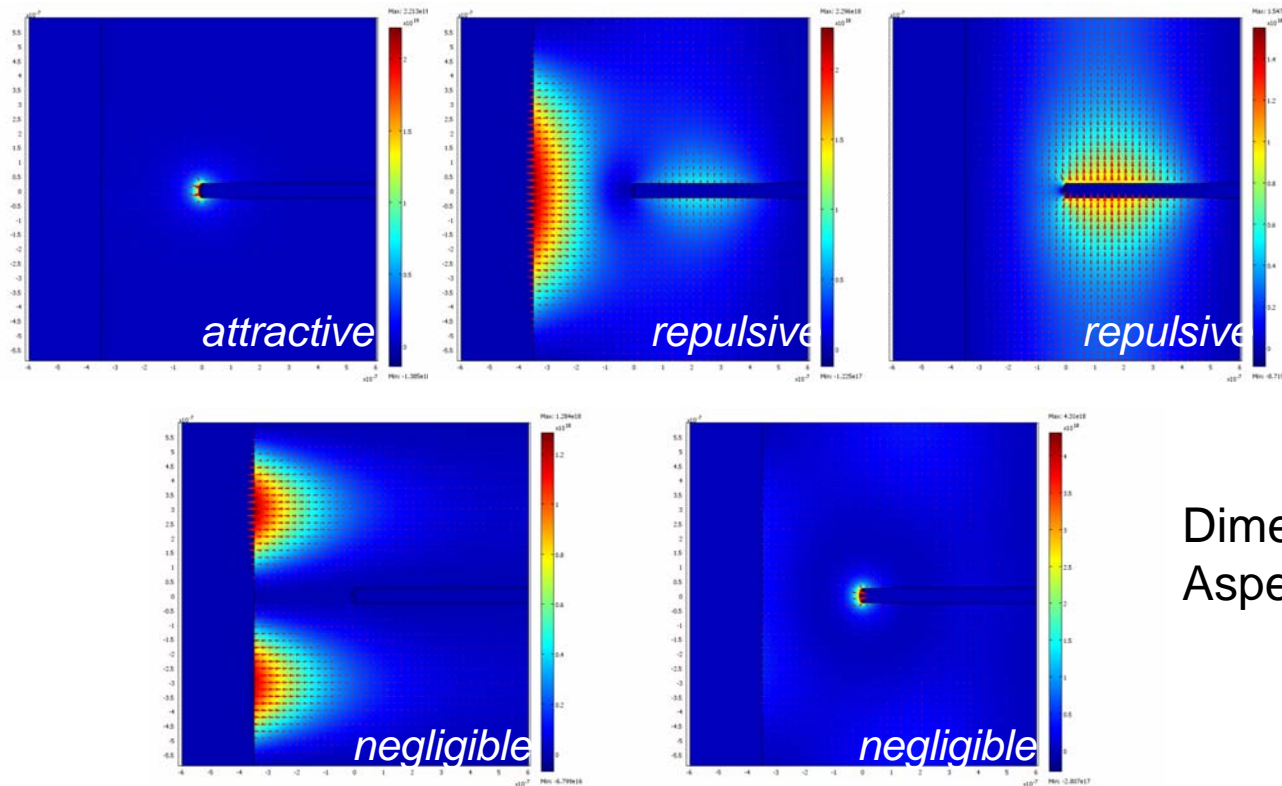


- Electroplating process
- HSQ molds (highest resolution resist)
- 100keV electron beam lithography tool
 - pattern thick resist (1 μm)
 - large depth of focus
 - small electron scattering



Geometrical effects

- “Role of surface plasmons in the Casimir effect”, F. Intravaia et al., PRA 76 (2007)
- Metallic nanowire ($w < \lambda_p$) close to a flat metallic surface



Dimension $< 100\text{nm}$
Aspect ratio $> 5:1$

Net contribution from the first 5 plasmonic modes is repulsive for $d \geq 200\text{nm}$



Summary

- **Precise experiments of the Casimir force between Au-Au surfaces**
- **Good agreement with plasma model**

Differences with Drude model **cannot** be explained as a problem in the separation measurement, or the Au layer properties. It appears that any model with a finite relaxation time will give discrepancies when comparing with the Casimir force.
Why do Casimir modes decouple from the dissipative part?

- **Geometrical effects**

An innovative MEMS that allows to modify the geometry *in situ* is being designed and tested. This system will be used to investigate the plasmonic contributions to the Casimir effect.