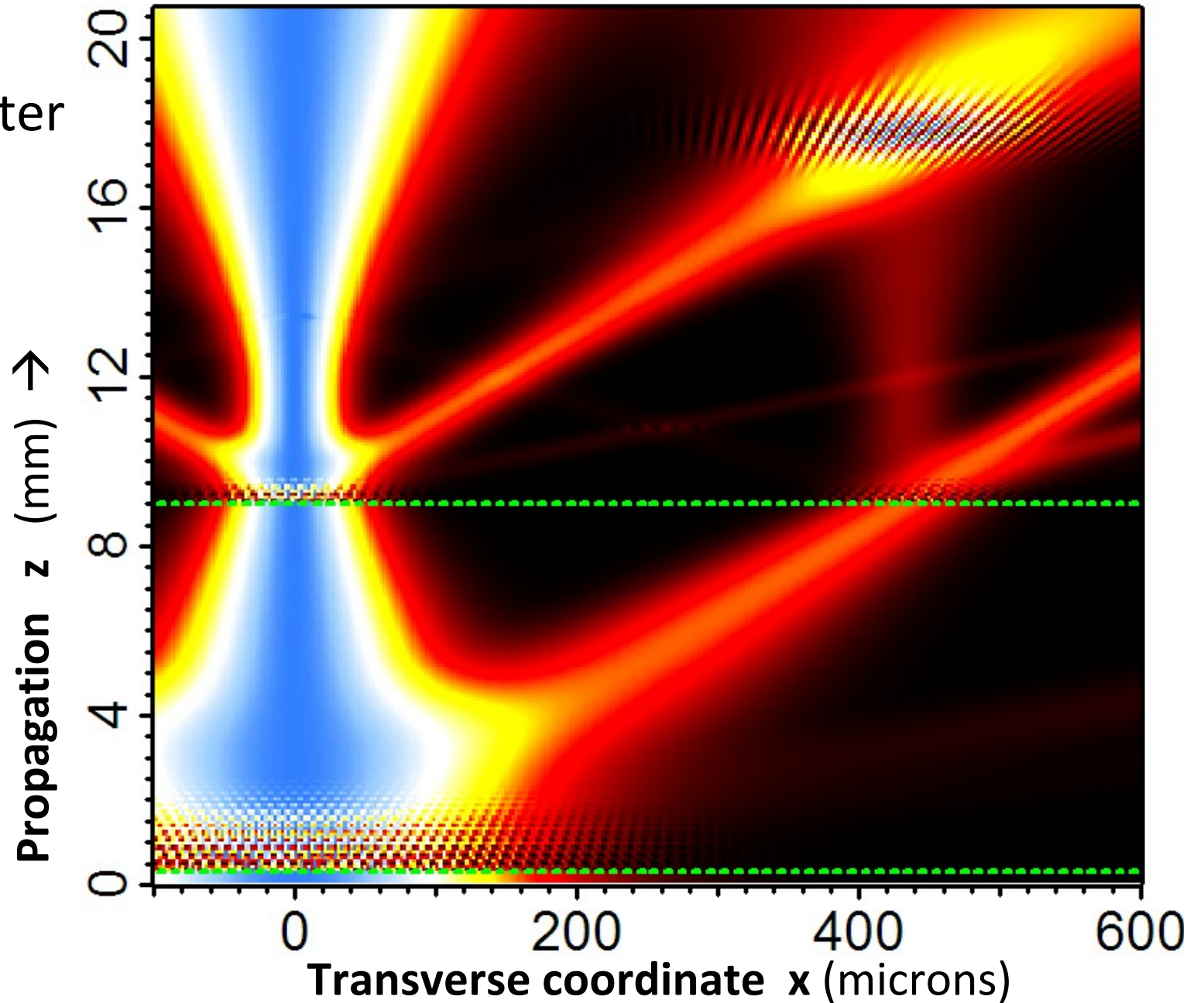


Grating  
Interferometer  
Simulation

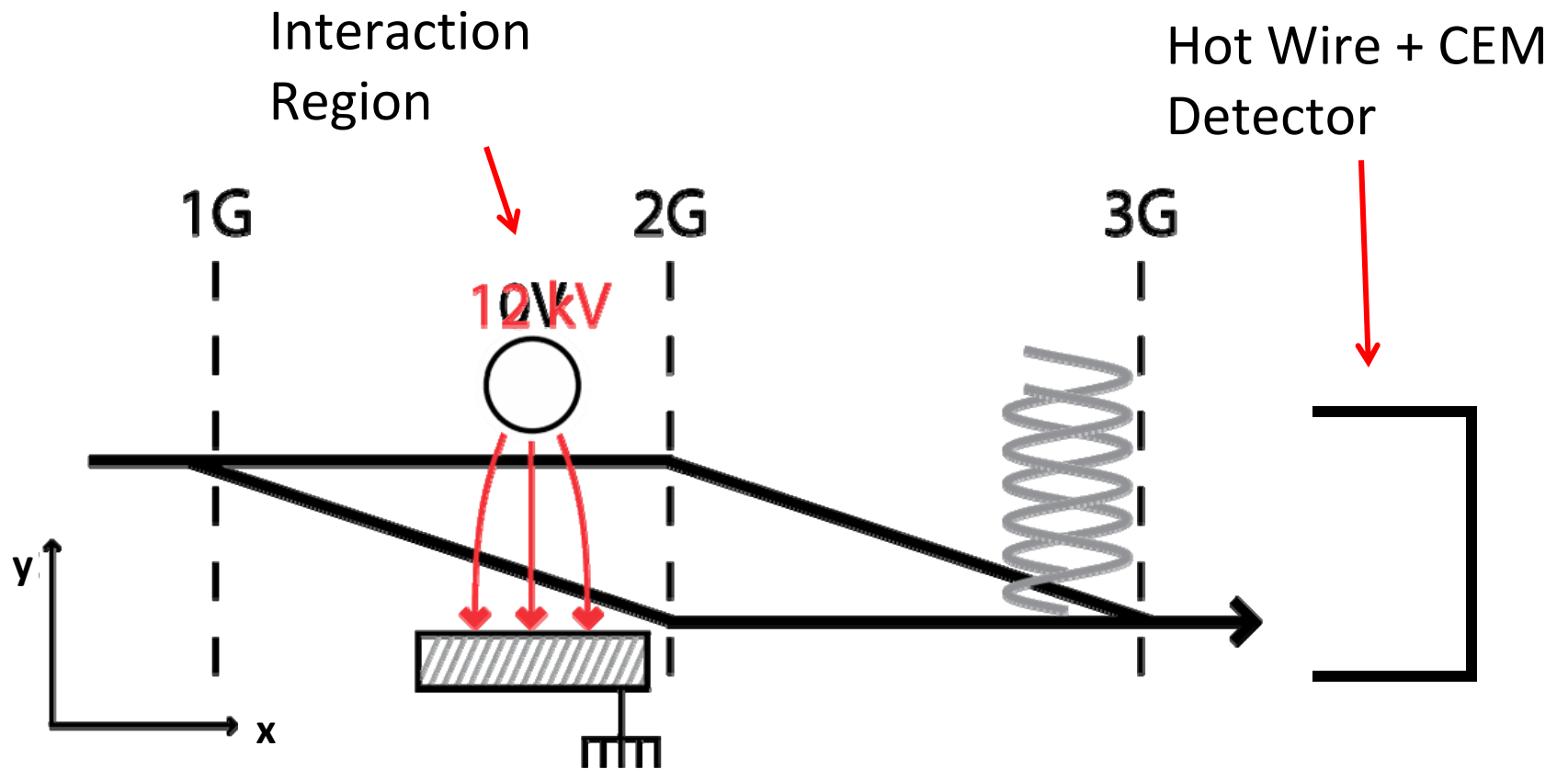


McMorran B. and Cronin, A.D., "Model for partial coherence and wavefront curvature in grating interferometers," *Phys. Rev. A.* **78**, 013601 (2008)

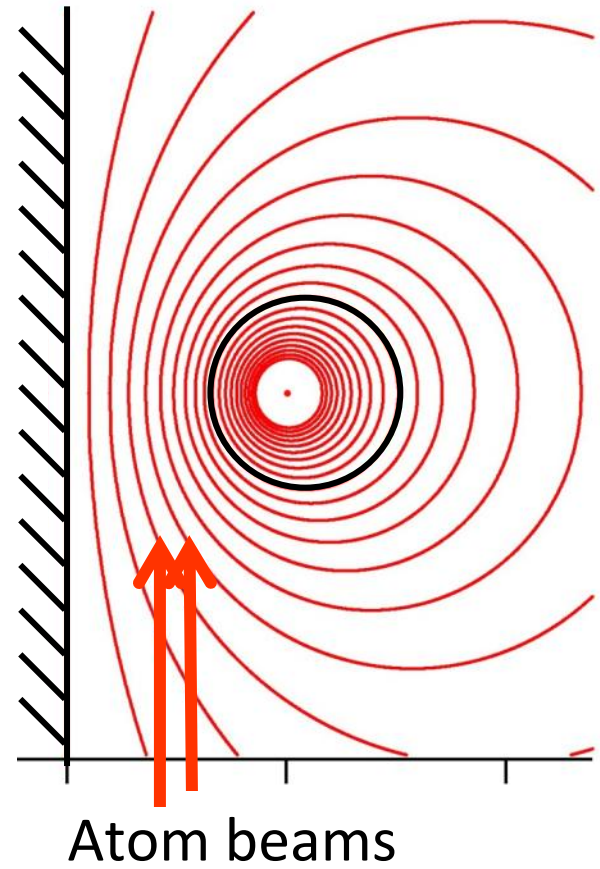
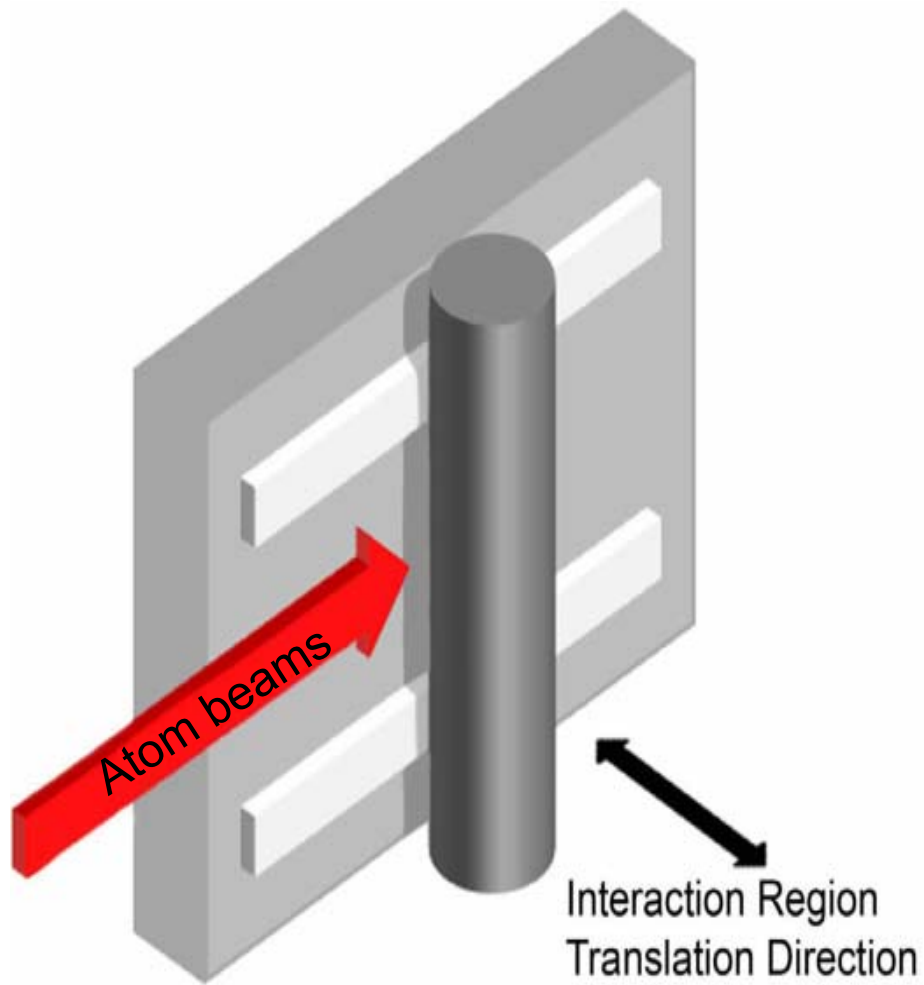
# Outline

- 1. New polarizability measurements**
- 2. Atom Interferometry C3 measurements**  
**tests of potential power law**  
**limits on non-Newtonian gravity**
- 3. More precise diffraction experiments**  
**Magic open fraction**  
**Effect of AuPd**  
**Ratios of C3 for Rb, K, Na**
- 4. Laser modified vdW interactions**

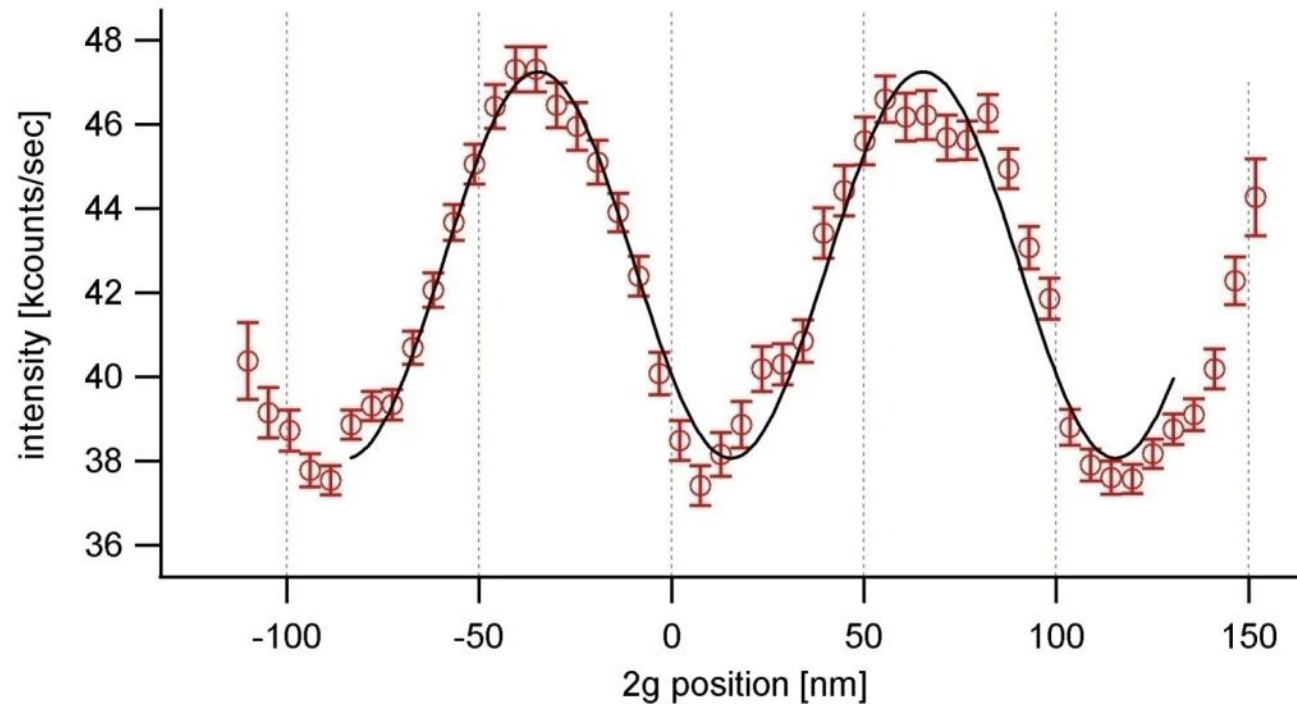
# Polarizability experiment



# Interaction Region

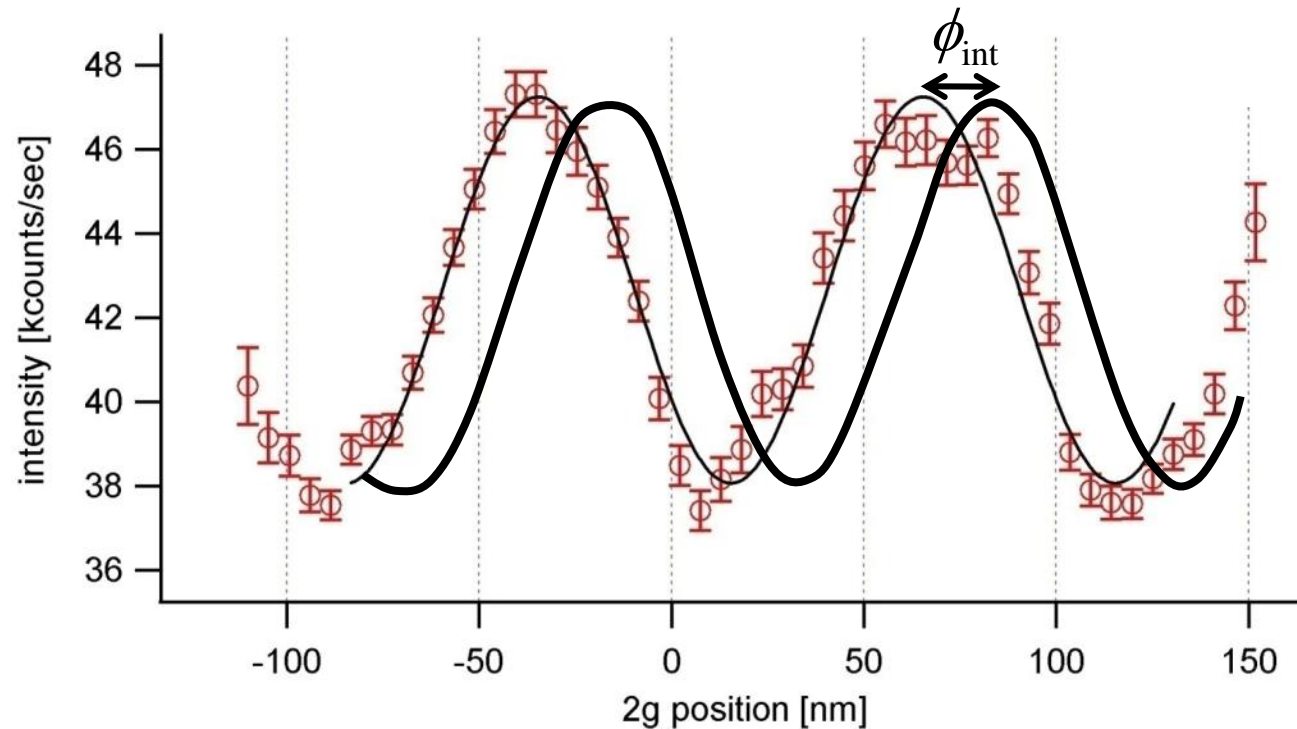


# Interference Fringe Data

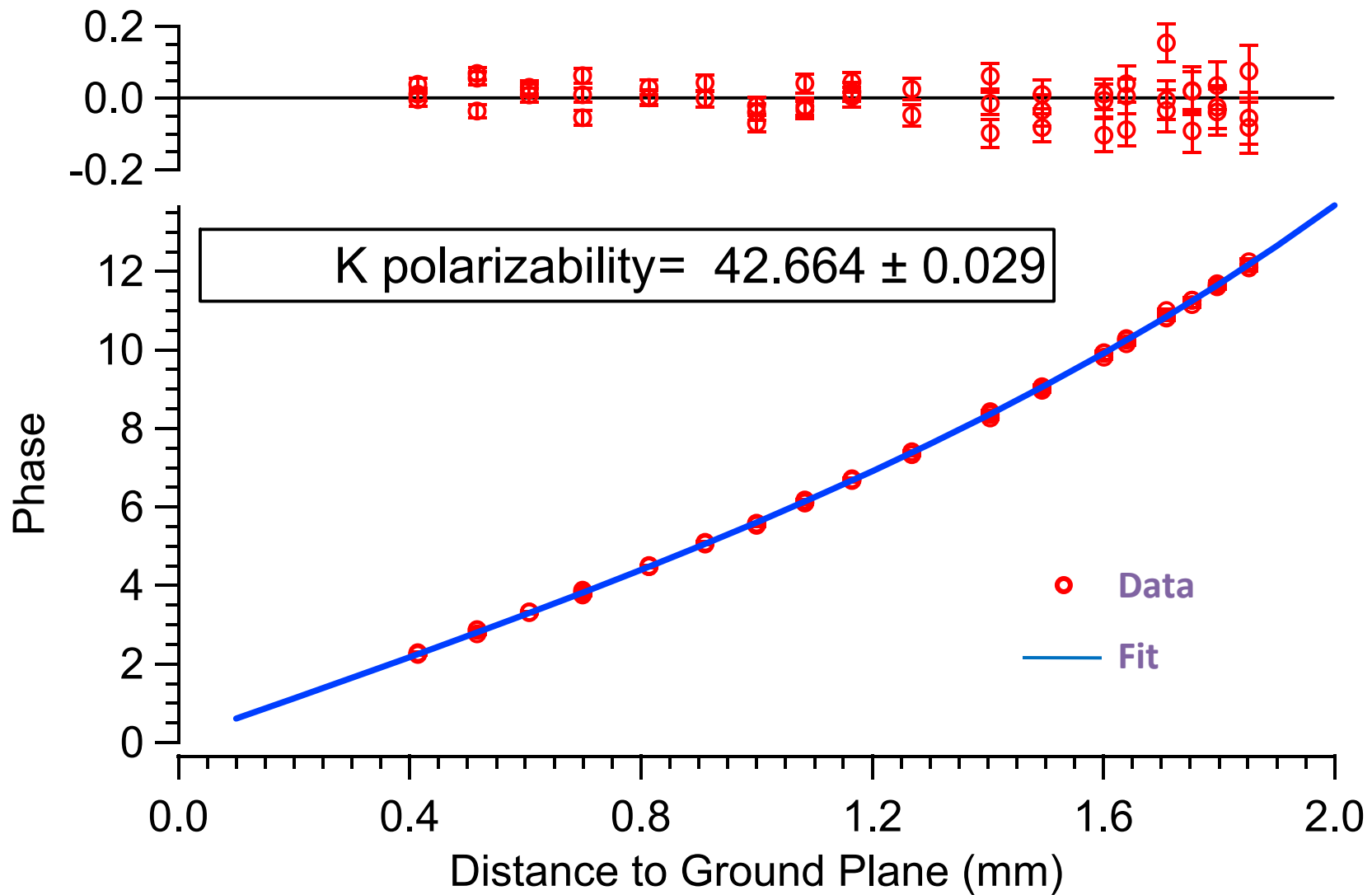


- We record atom flux vs grating position

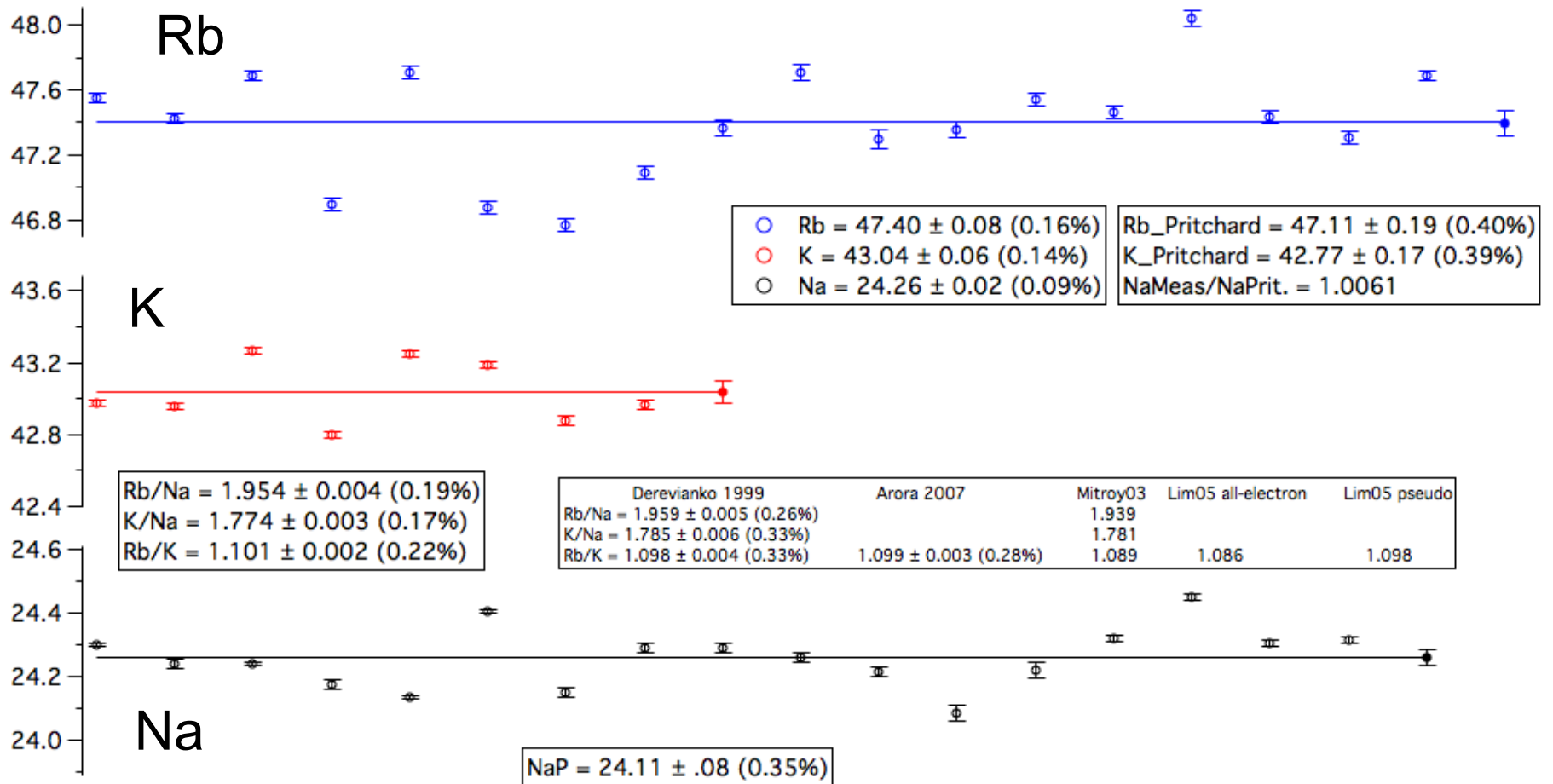
# Interference Fringe Data



- We record atom flux vs grating position
- Polarizability relates to this phase shift



# Preliminary Polarizability Experiment Results





## Preliminary Results from Polarizability Measurements

$$\alpha_{\text{Na}} = 24.26 \quad (0.09\% \text{ stat}) \quad (1\% \text{ syst})$$

$$\alpha_{\text{K}} = 43.04 \quad (0.14\% \text{ stat}) \quad (1\% \text{ syst})$$

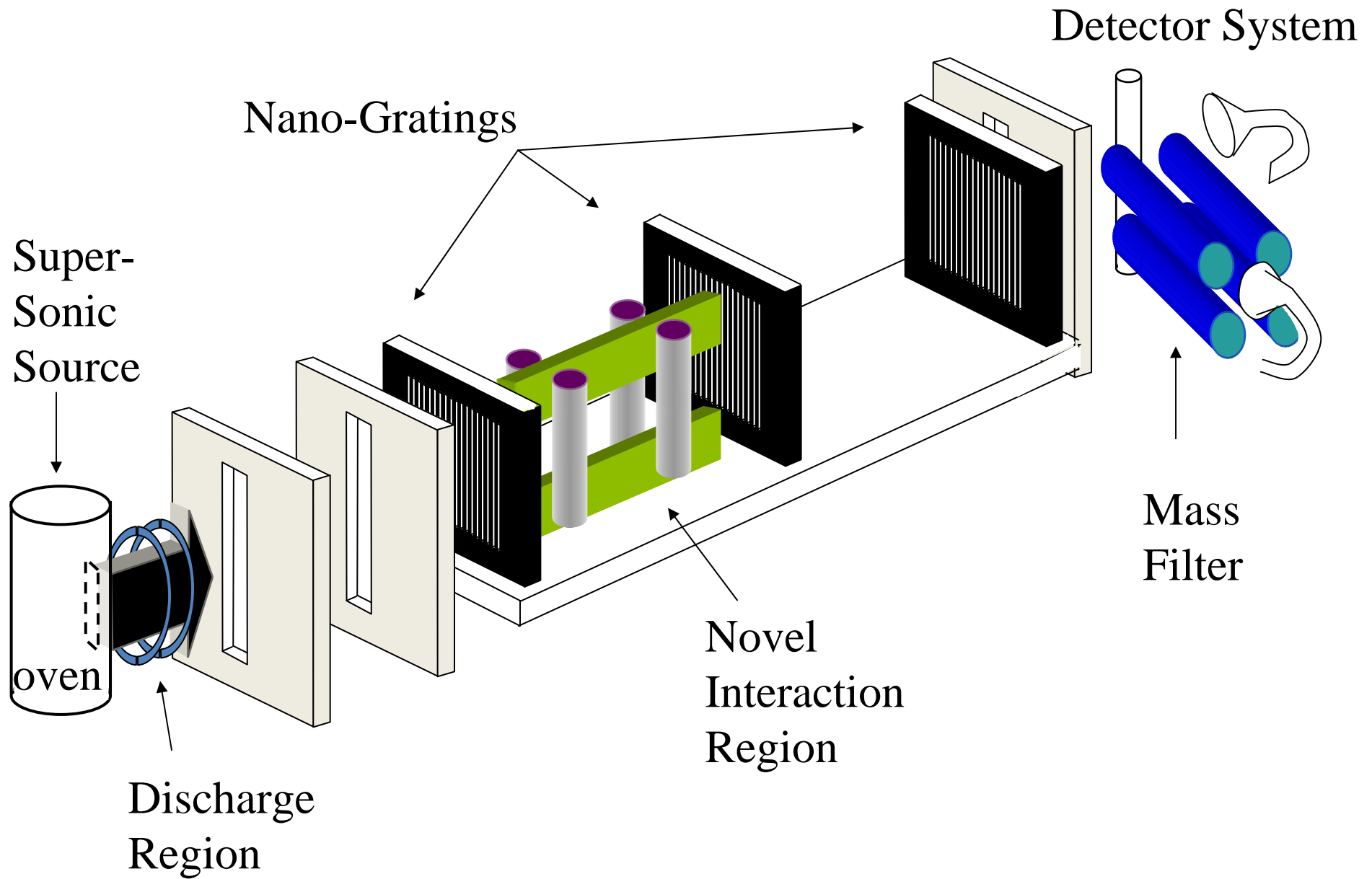
$$\alpha_{\text{Rb}} = 47.40 \quad (0.16\% \text{ stat}) \quad (1\% \text{ syst})$$

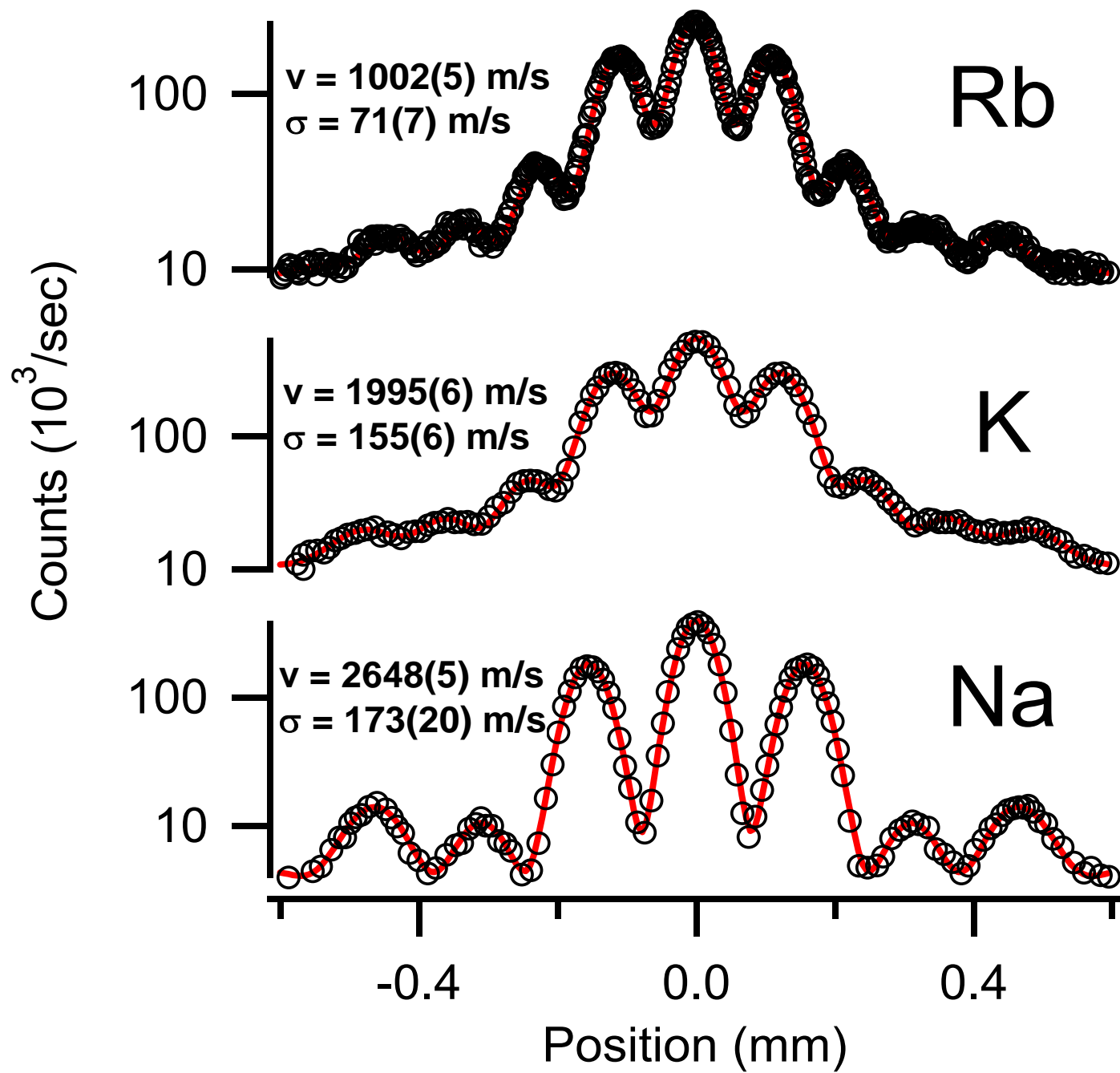
**U. Arizona (2009) Experiment:**  $\frac{\alpha_{\text{K}}}{\alpha_{\text{Na}}} = 1.774 \pm 0.003$

**Derevianko (1999) Calculation:**  $\frac{\alpha_{\text{K}}}{\alpha_{\text{Na}}} = 1.785 \pm 0.006$

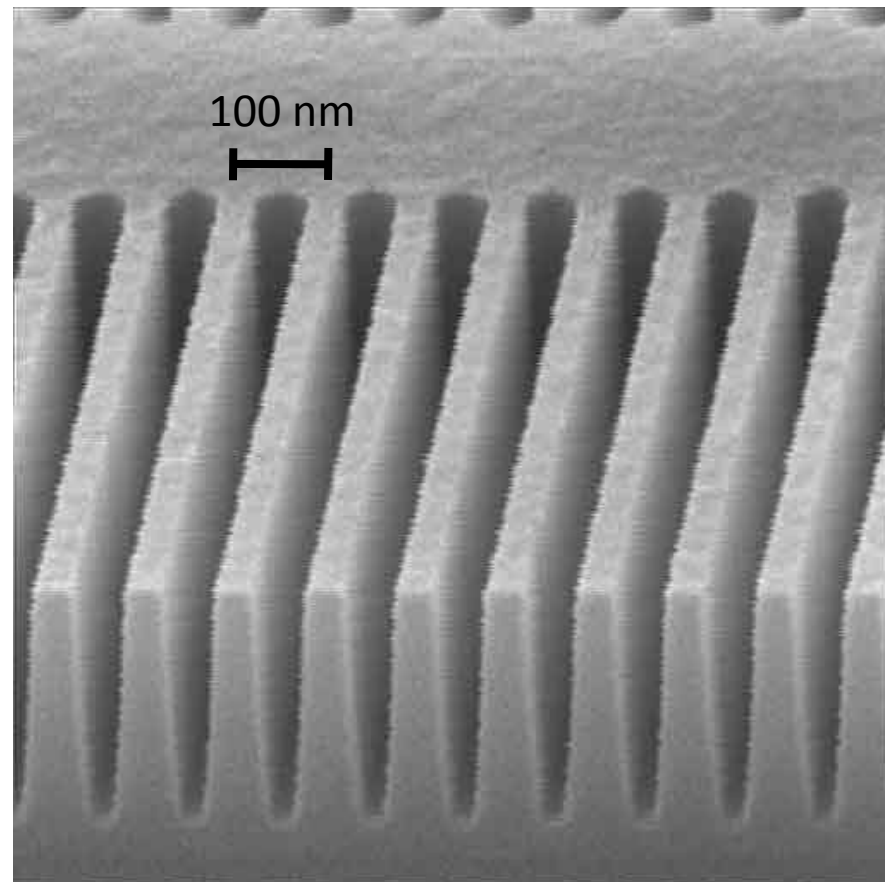
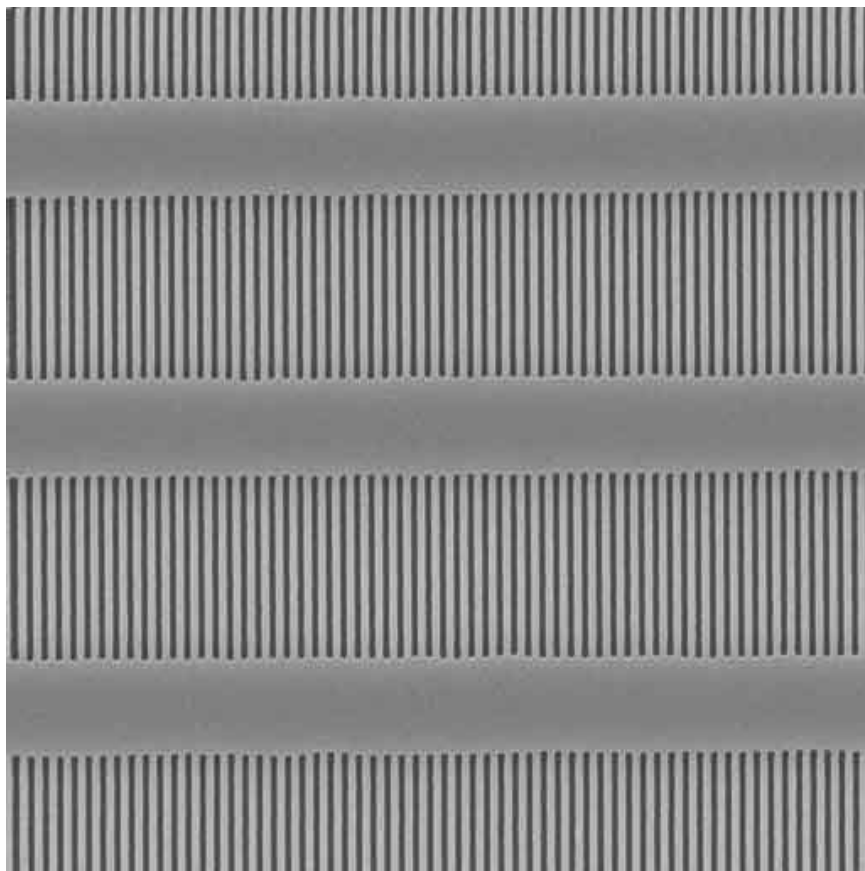
**Experiment ratio precision: 0.26%    Theory: 0.33%.    Disagreement: 0.6%**

# Plan for $\alpha_{Sr^*}$





# Nano-Grating beam splitters

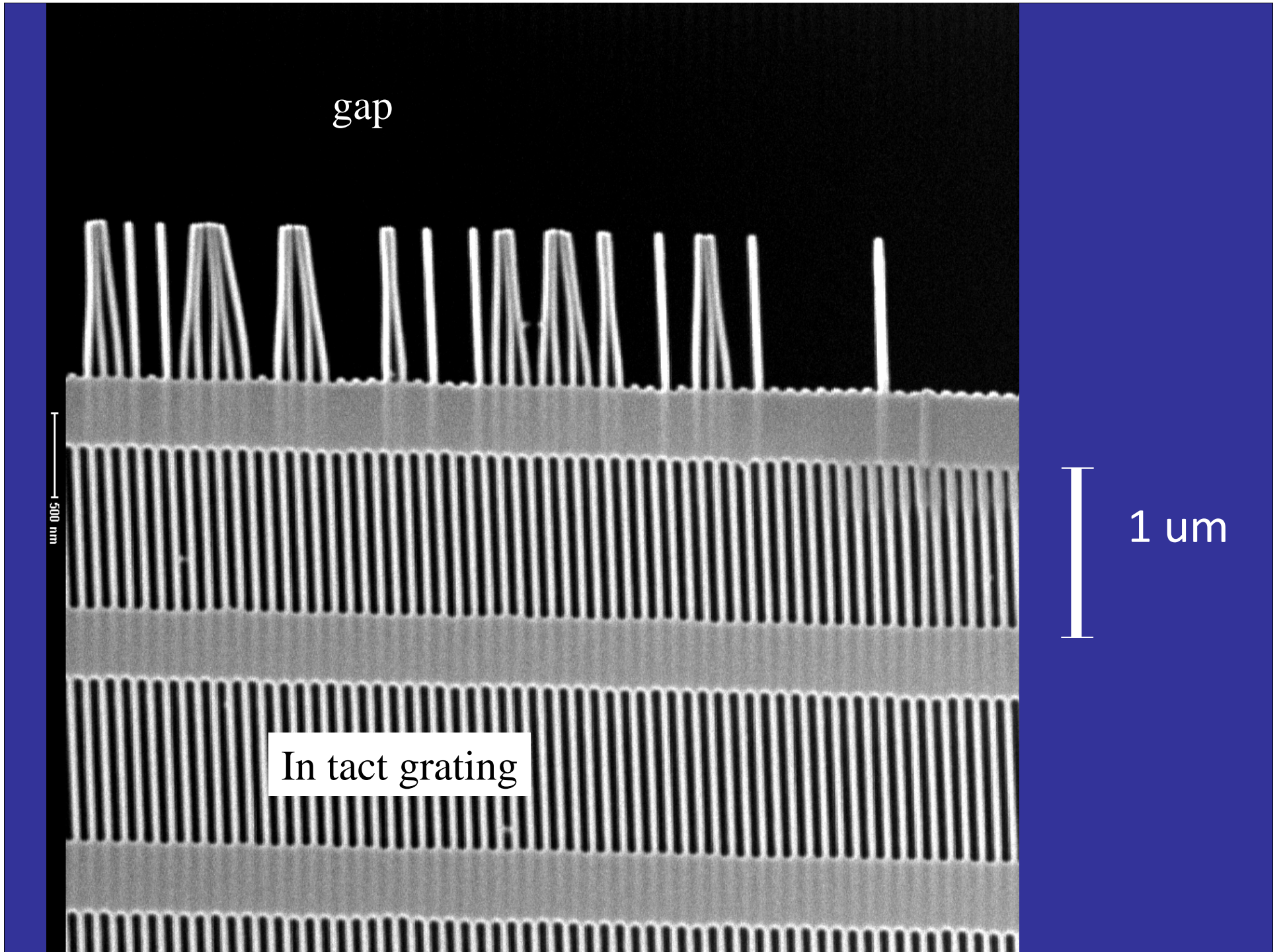


gap

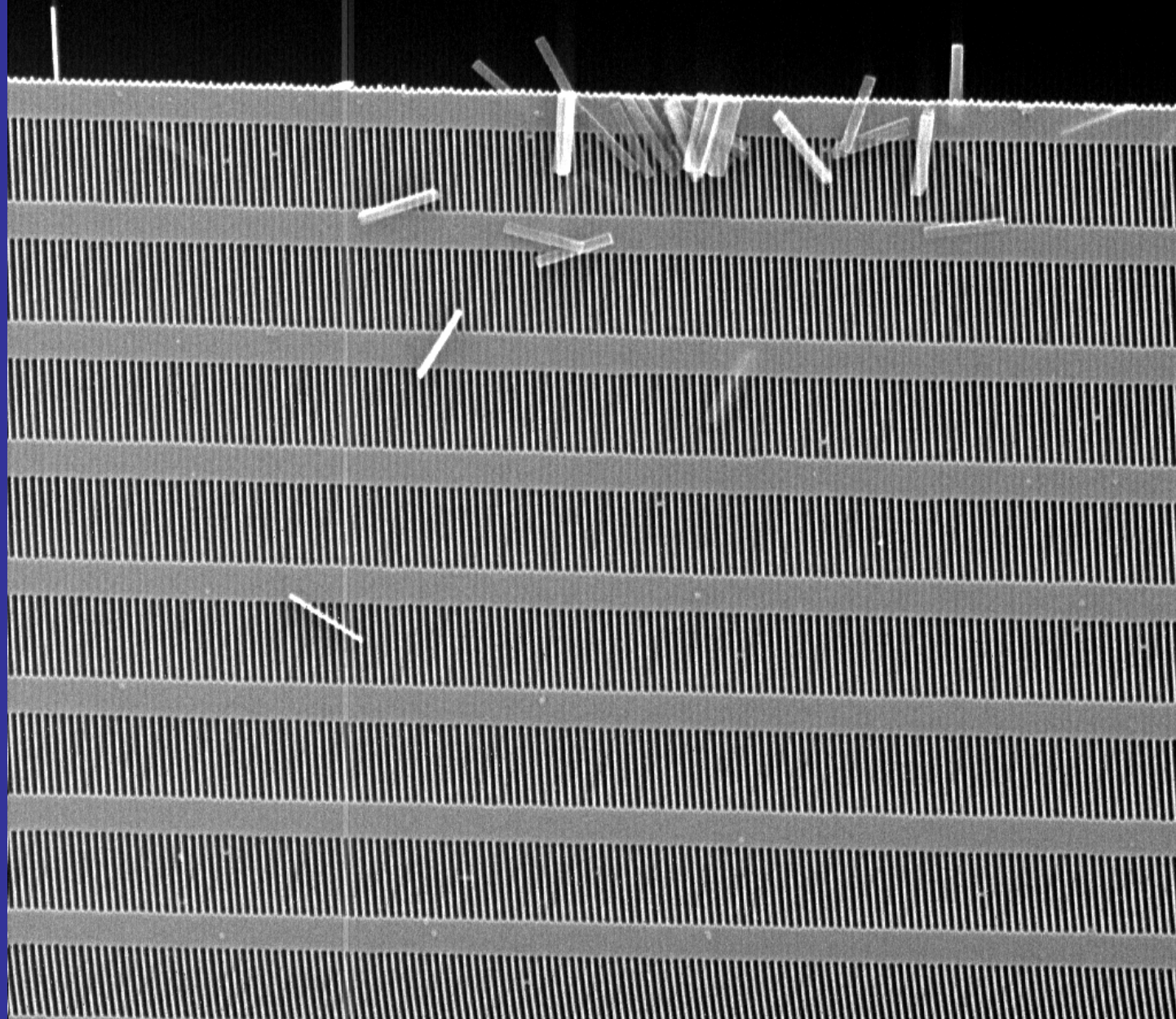
500 nm

In tact grating

1  $\mu\text{m}$



# Nano-Structure Grating



1 um

1 um

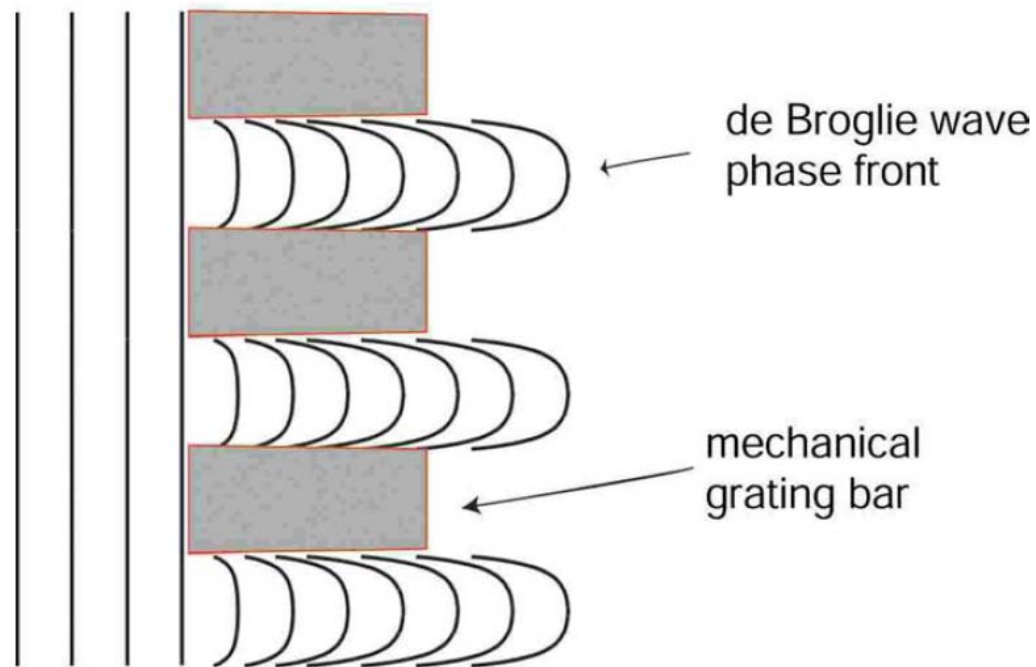


FIG. 63. (Color online) Distorted de Broglie waves. van der Waals interactions with mechanical grating bars cause near-field phase shifts. This view is exaggerated: in beam experiments there are typically  $10^4$  wave fronts in the 100-nm thickness of a nanograting slot (Grisenti *et al.*, 1999; Perreault *et al.*, 2005).

Cronin, A.D., Schmiedmayer J., and Pritchard D.E., "Optics and Interferometry with atoms and molecules." *Rev. Mod. Phys* **81**, 1051 (2009)

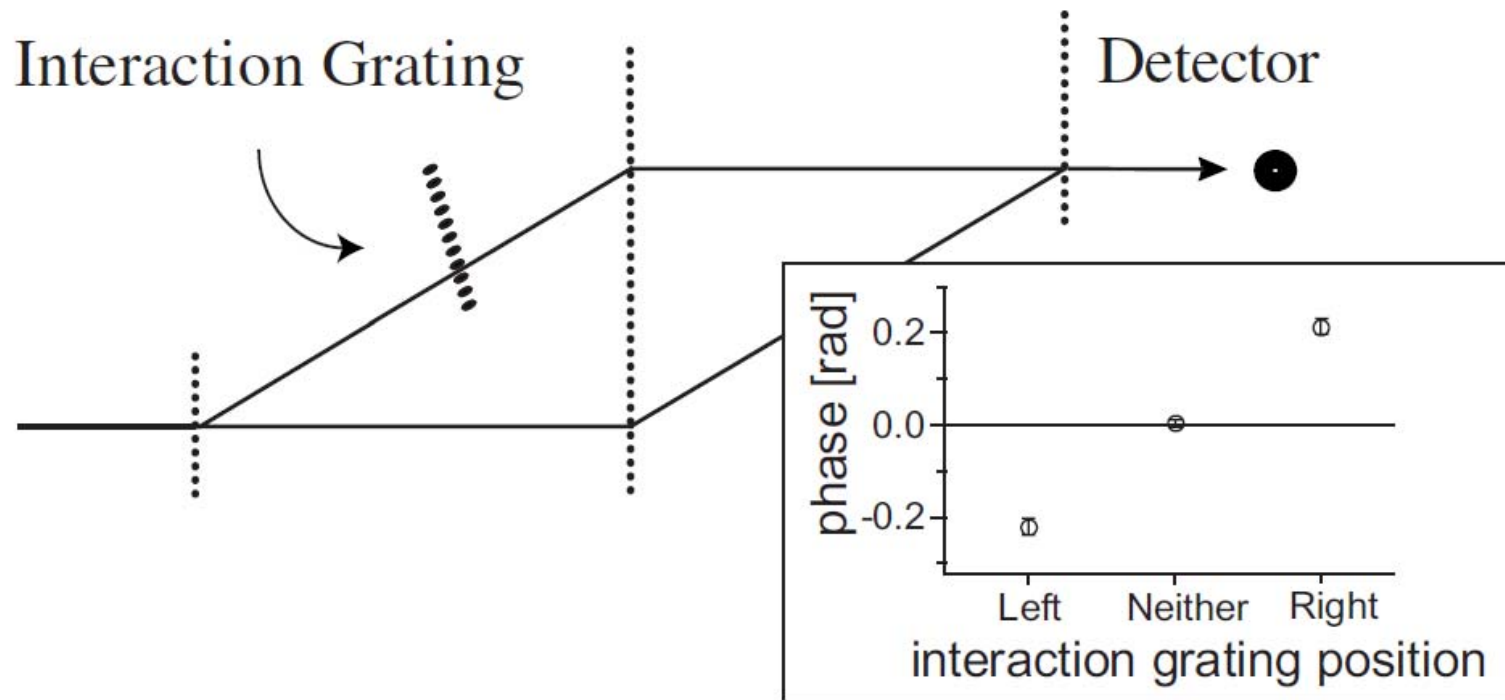


FIG. 68. An “interaction grating” was inserted and removed from each path of an interferometer to measure the phase shift  $\Phi_0$  due to van der Waals interactions (Perreault and Cronin, 2005).

Cronin, A.D., Schmiedmayer J., and Pritchard D.E., “Optics and Interferometry with atoms and molecules.” *Rev. Mod. Phys* **81**, 1051 (2009)



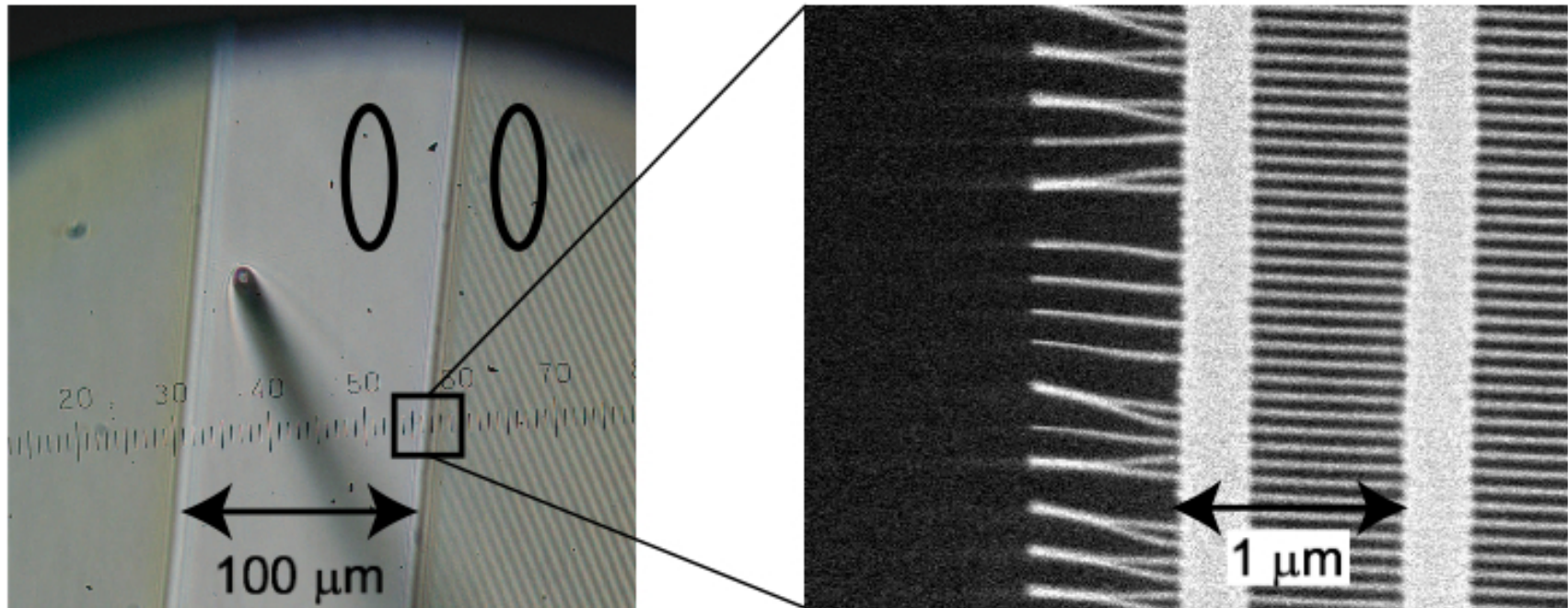
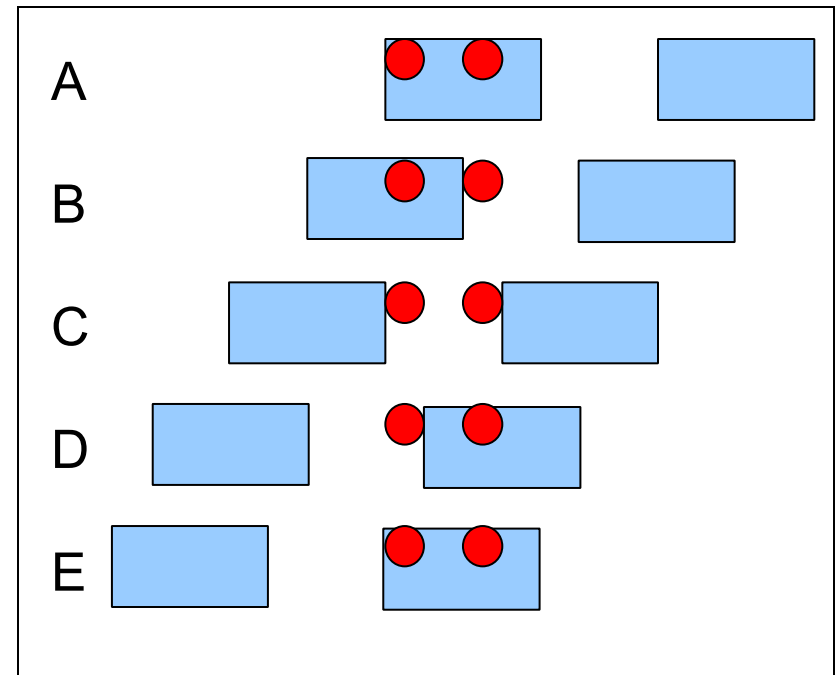
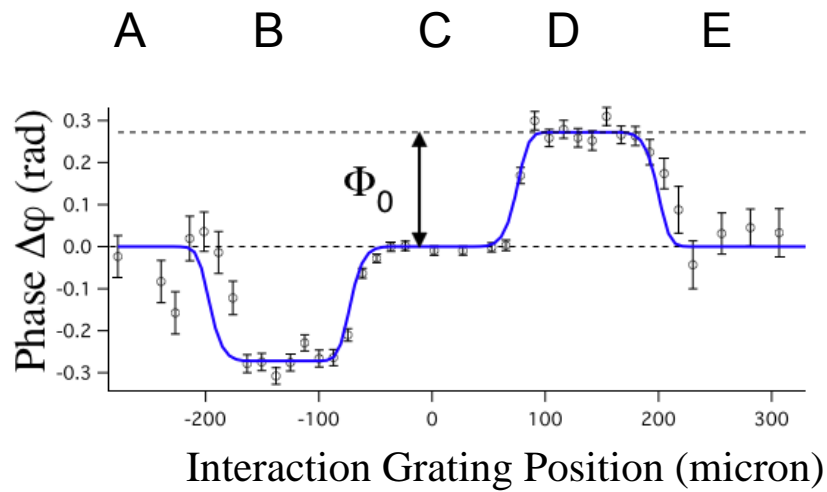


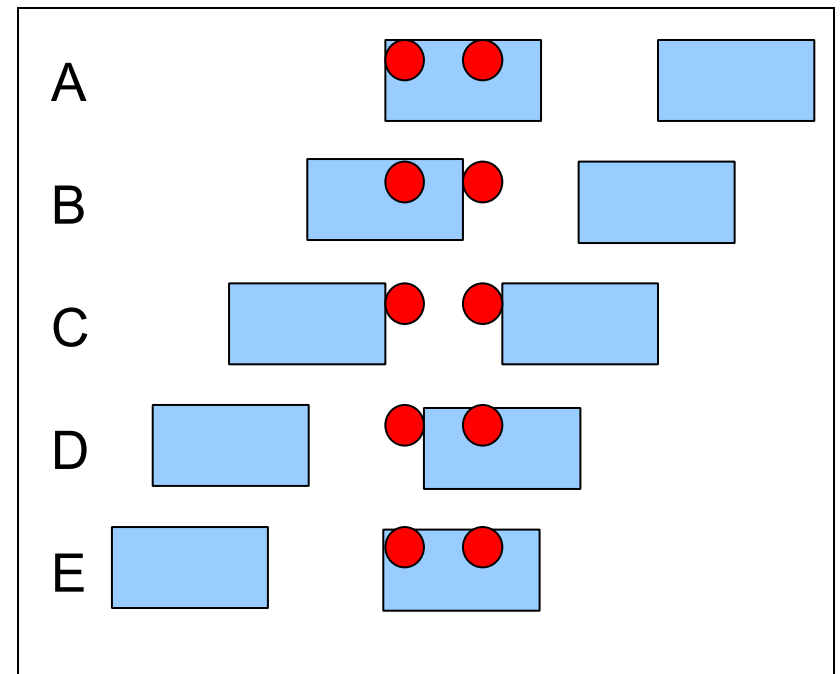
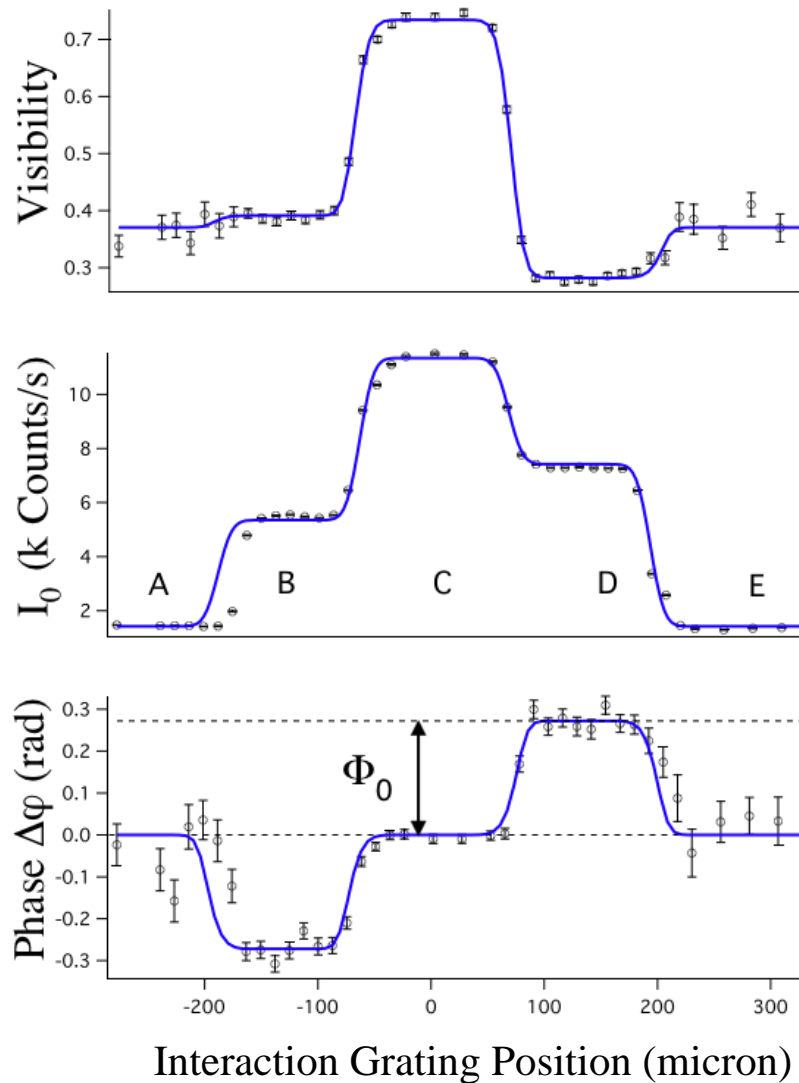
Figure 5.4: Preparation of interaction grating. The left image shows an optical microscope image with the glass capillary tube in the perforation. The ellipses represent the test and reference arms (not to scale) of the interferometer going into the page. The scanning electron microscope image on the right shows a further magnified view of the transition from intact grating to gap.

# Phase vs Interaction Grating Position



Lepoutre, S., Haikel J., Trenec G., Bruchner M., Vigue J., Lonij V., and Cronin A.D.,  
“Dispersive atom interferometry phase shifts due to atom-surface interactions,”  
submitted to Europhysics Letters, (2009)

# Visibility, Intensity, and Phase



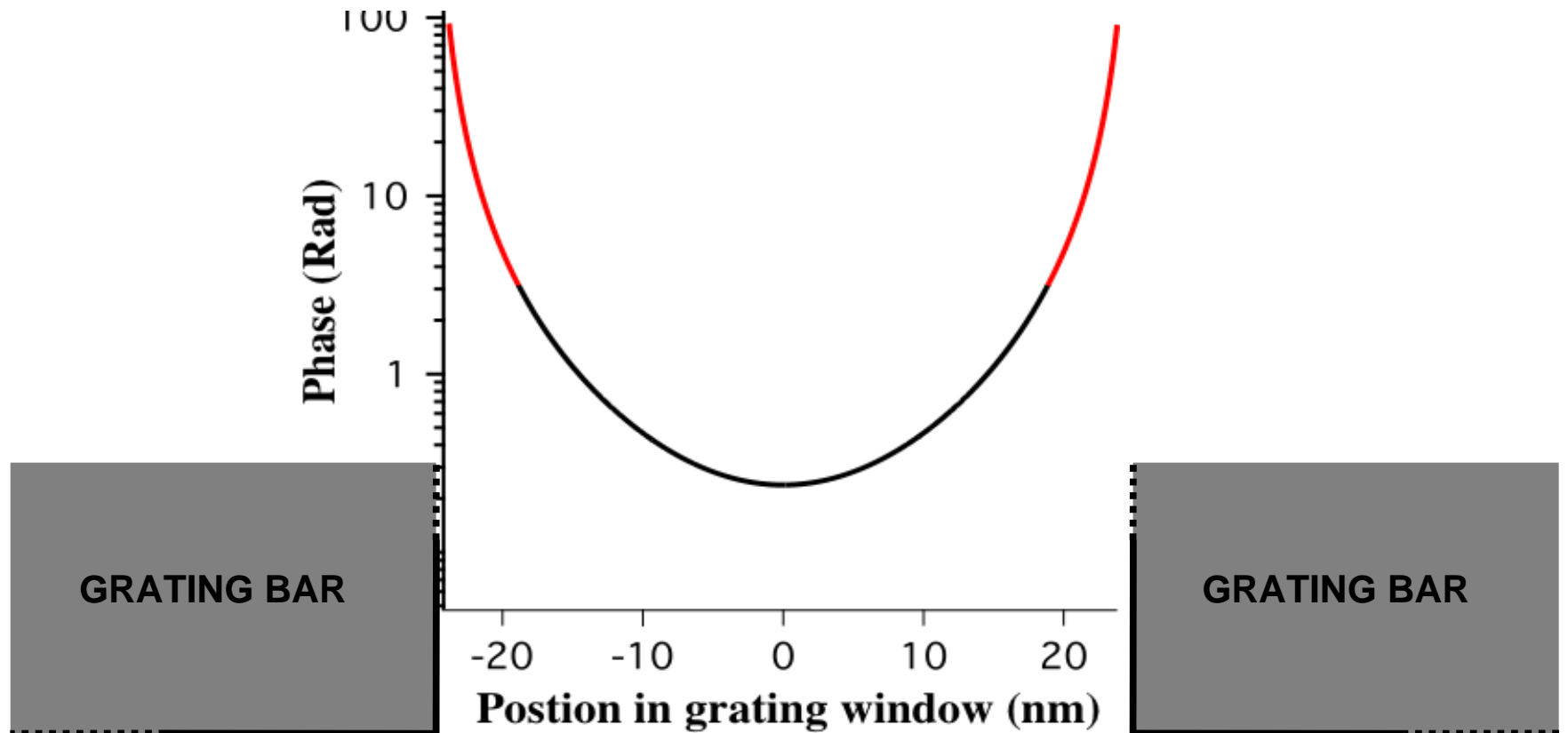
# Acquired phase

near field

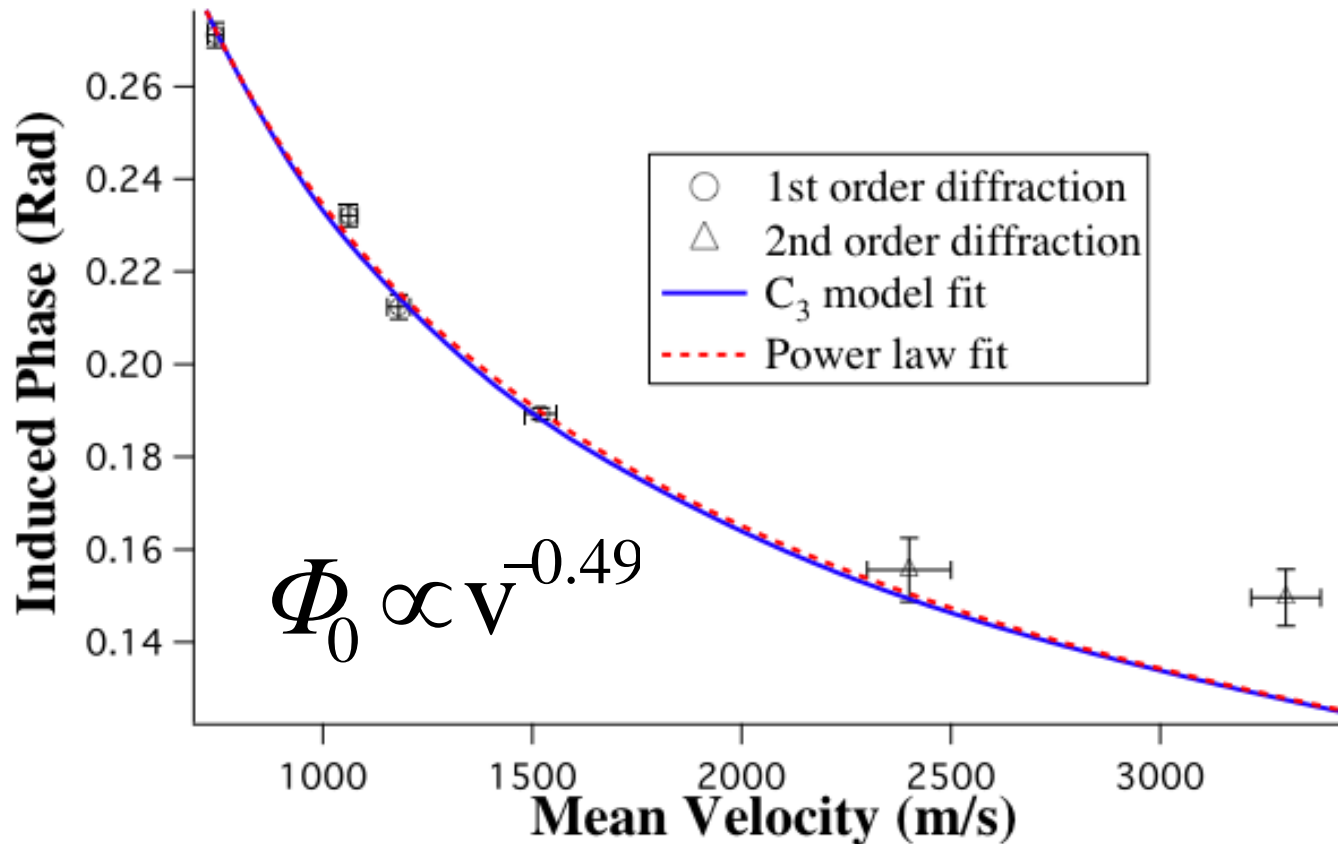
$$\phi(x) \propto \frac{C_3}{\hbar} \left[ x_1^{-3} + x_2^{-3} \right] \frac{L}{v}$$

far field

$$\Phi_0 = \arctan \left( \frac{\int_{-w/2}^{w/2} \sin[\phi(x)] dx}{\int_{-w/2}^{w/2} \cos[\phi(x)] dx} \right)$$

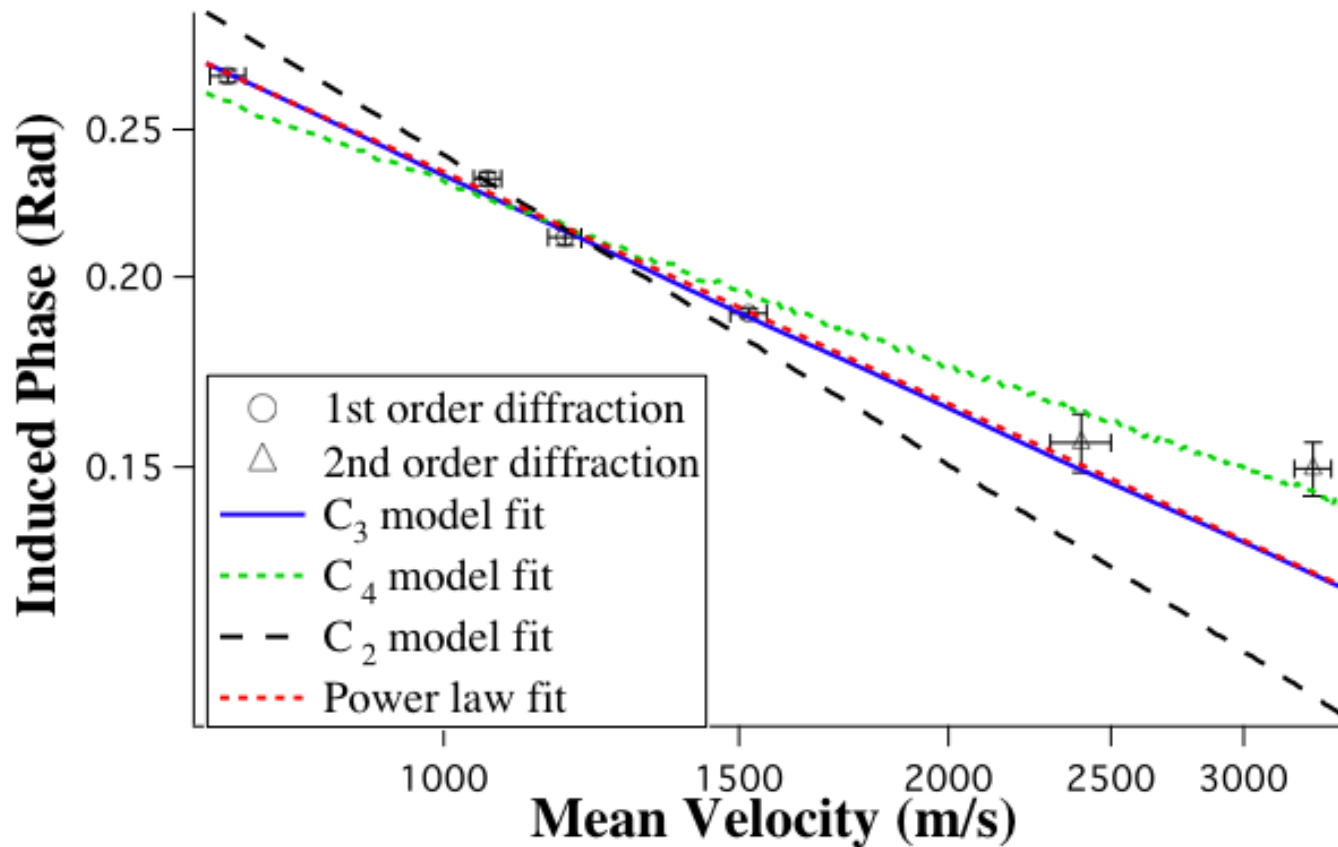


# Velocity Dependence



Lepoutre, S., Haikel J., Trenec G., Bruchner M., Vigue J., Lonij V., and Cronin A.D.,  
“Dispersive atom interferometry phase shifts due to atom-surface interactions,”  
submitted to Europhysics Letters, (2009)

# Testing the Potential Power Law

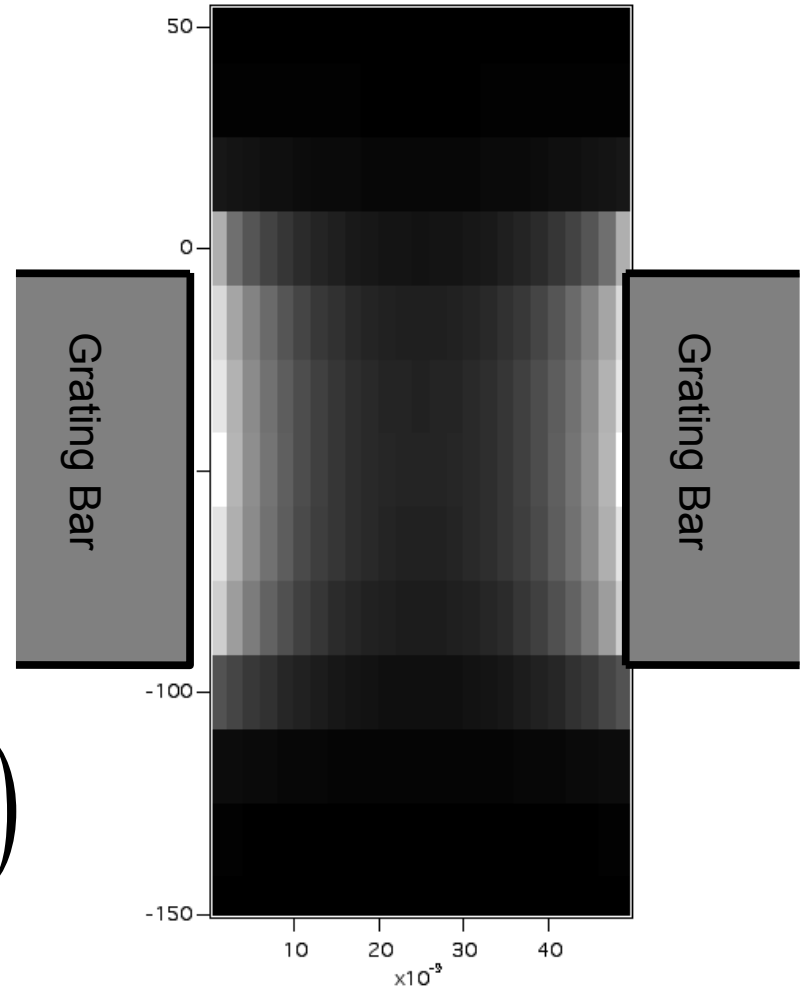


$$U = - \frac{C_p}{r^p} \quad p = 2.9 \pm .2$$

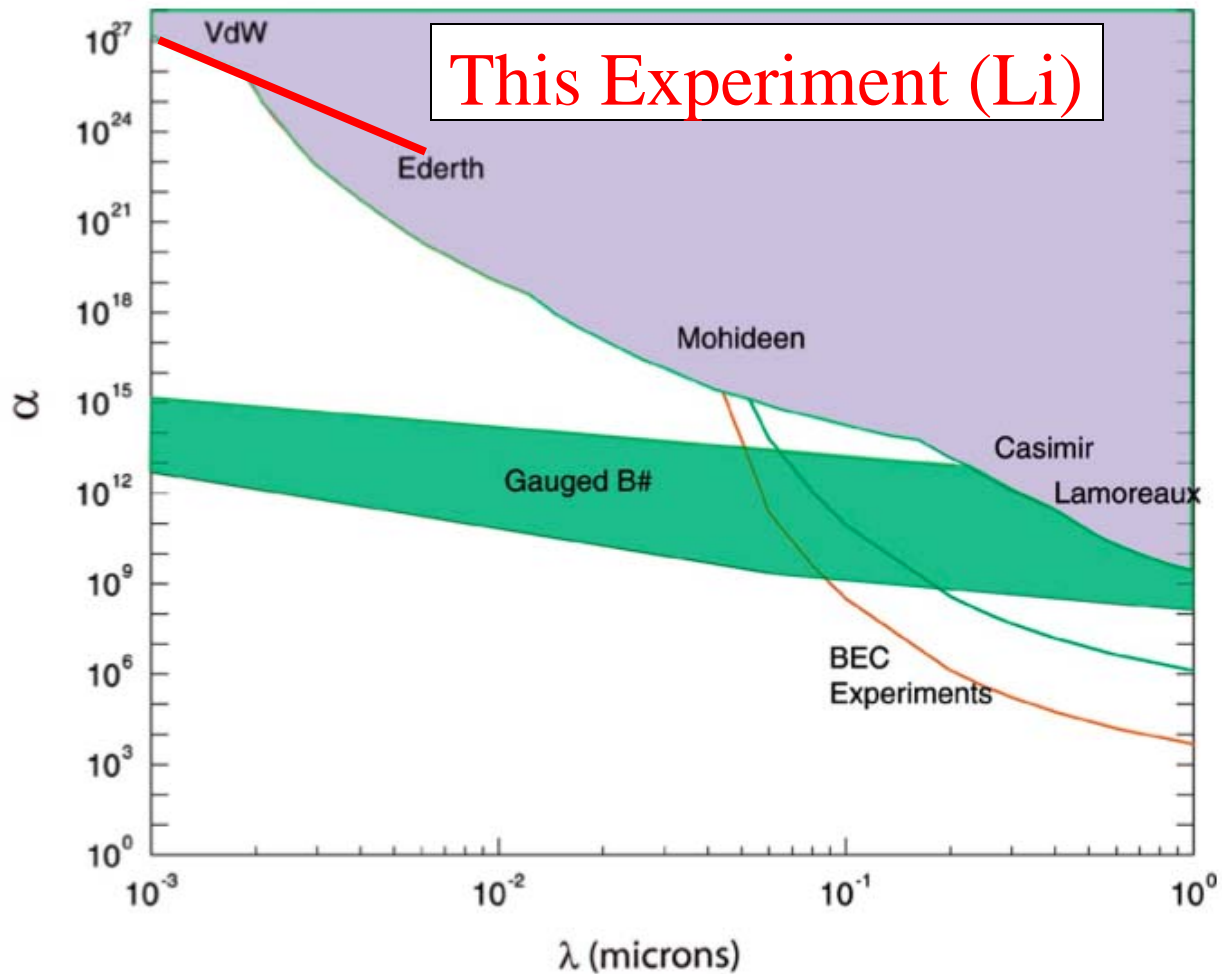
# Non-Newtonian Gravity

Potential inside the grating due to modified gravitational interaction on a log scale

$$U_{\infty} = G \frac{m_1 m_2}{r} \left( 1 + a e^{-r/\lambda} \right)$$



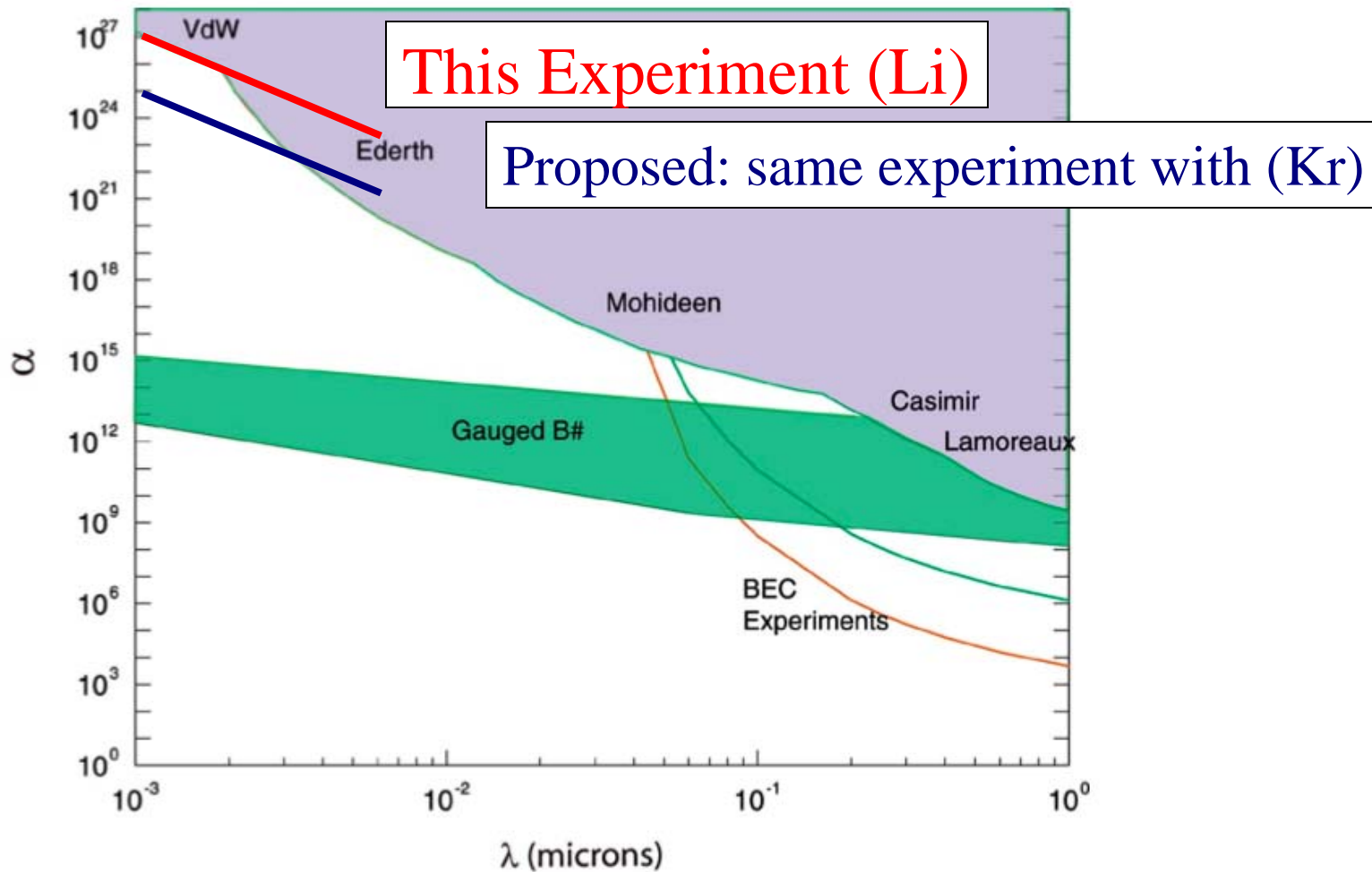
# Non-Newtonian Gravity



“Probing submicron forces by interferometry of Bose-Einstein condensed atoms”  
Savas Dimopolous and Andrew Geraci PRD 68, 124021 (2003)

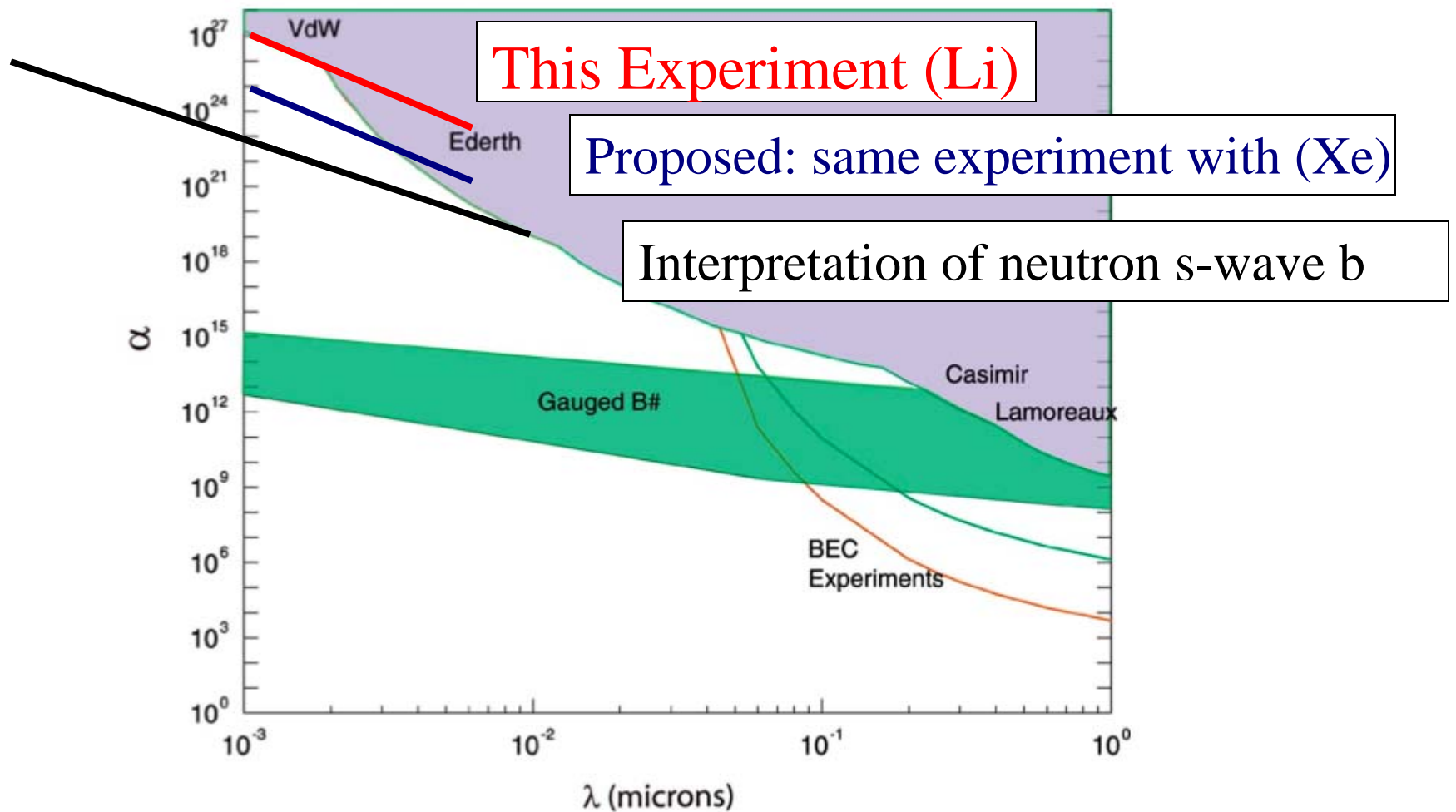


# Non-Newtonian Gravity



“Probing submicron forces by interferometry of Bose-Einstein condensed atoms”  
Savas Dimopolous and Andrew Geraci PRD 68, 124021 (2003)

# Non-Newtonian Gravity



“Probing submicron forces by interferometry of Bose-Einstein condensed atoms”  
Savas Dimopolous and Andrew Geraci PRD 68, 124021 (2003)

S-wave  
scattering  
length

nN Gravity

$$\Phi = N\lambda_{dB}bL = \frac{UL}{\hbar v}$$

$$\frac{N\hbar bL}{m_n v} = \frac{2\pi G m_n (N m_n A) \alpha \lambda^2 L}{\hbar v}$$

$$\alpha \lambda^2 < 10^7 m^2$$

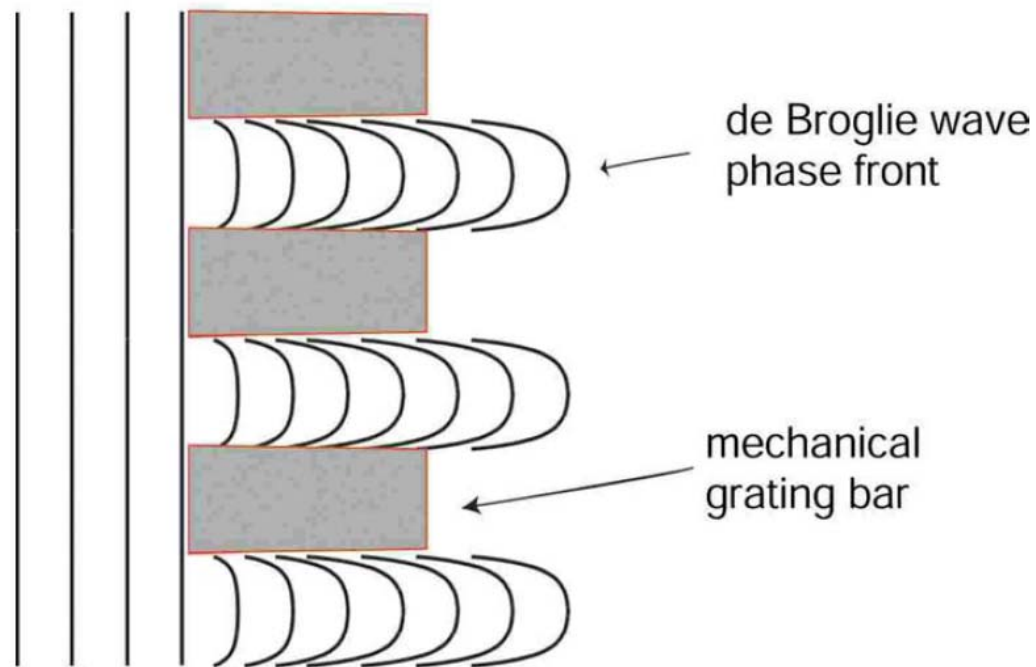


FIG. 63. (Color online) Distorted de Broglie waves. van der Waals interactions with mechanical grating bars cause near-field phase shifts. This view is exaggerated: in beam experiments there are typically  $10^4$  wave fronts in the 100-nm thickness of a nanograting slot (Grisenti *et al.*, 1999; Perreault *et al.*, 2005).

Cronin, A.D., Schmiedmayer J., and Pritchard D.E., "Optics and Interferometry with atoms and molecules." *Rev. Mod. Phys* **81**, 1051 (2009)

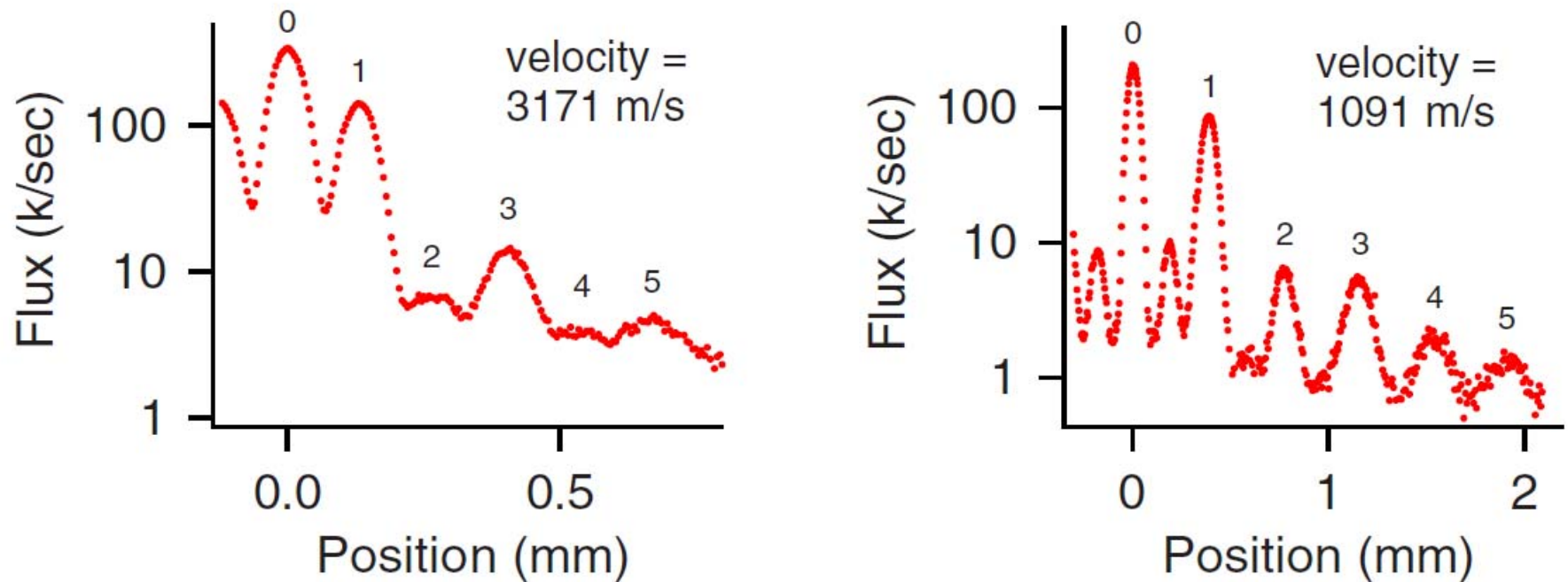
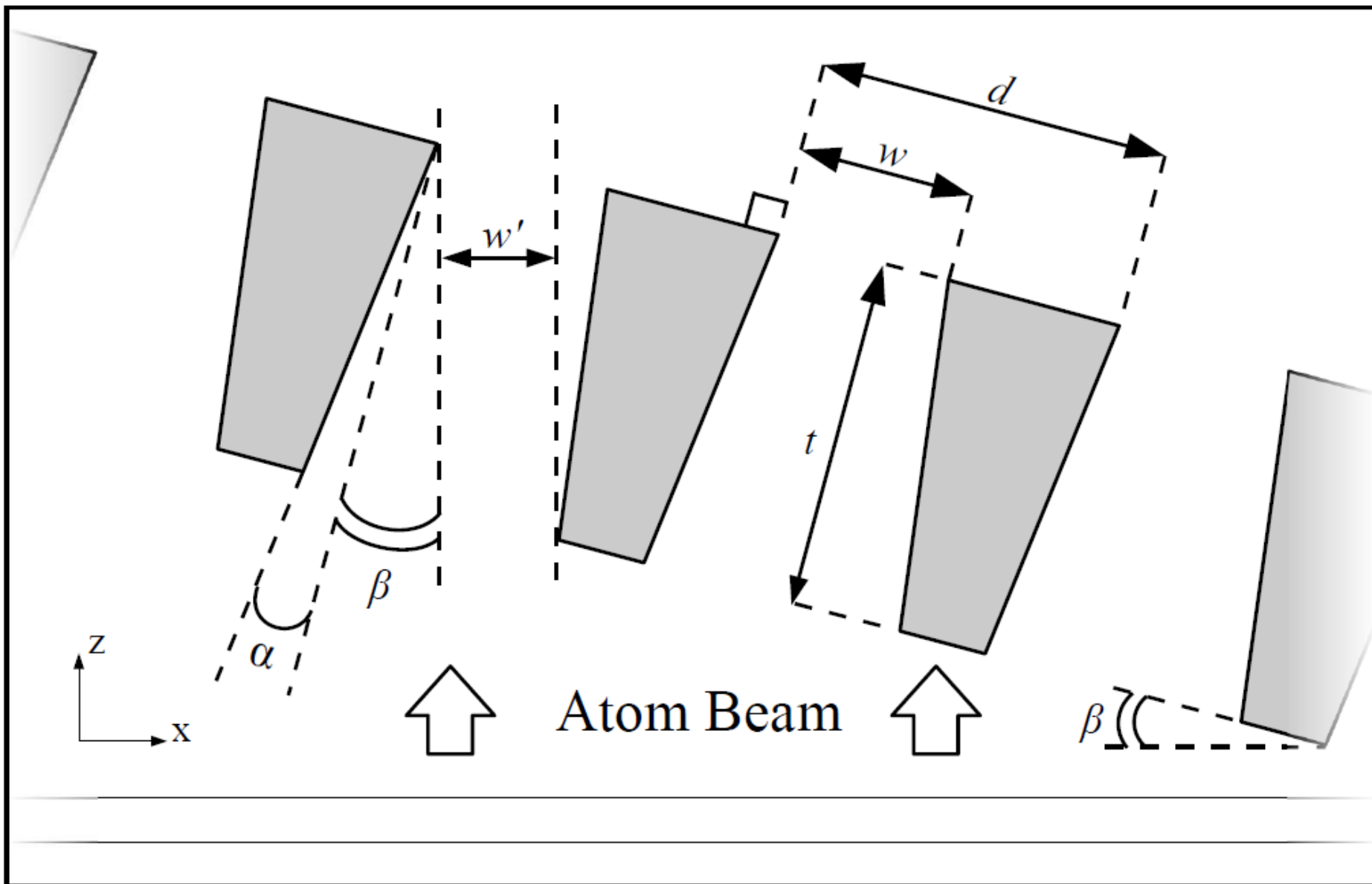


FIG. 65. (Color online) Diffraction intensities used to measure the strength of  $C_3$  for Na-silicon nitride (Perreault *et al.*, 2005). Data for two different velocities show how the second- and third-order change their relative intensity (as predicted in Fig. 64). Diffraction of  $\text{Na}_2$  molecules is also visible.

Cronin, A.D., Schmiedmayer J., and Pritchard D.E., "Optics and Interferometry with atoms and molecules." *Rev. Mod. Phys* **81**, 1051 (2009)



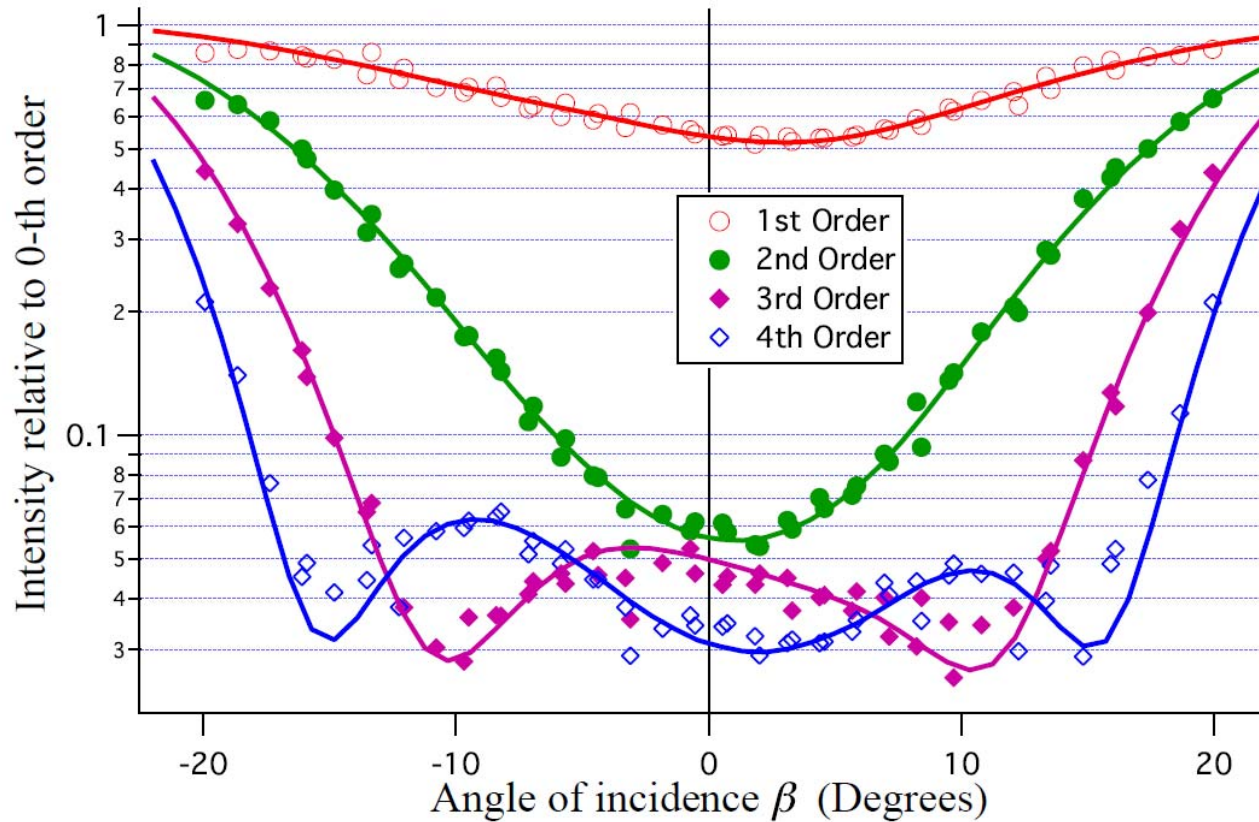


FIG. 6: The intensity of orders 1, 2, 3, and 4 relative to the 0-th order as a function of the grating rotation  $\beta$  (see Fig. 1). Each set of points at one angle represents one diffraction scan like the one shown in Fig. 2. A least squares fit to these data allows us to determine  $\alpha$  and  $t$ . Although  $w$  and  $C_3$  are left as free parameters in this fit, they are not well constrained. The data represented here are for a grating determined to have  $\alpha = 5 \pm 1$  degrees and  $t = 110 \pm 10$  nm.

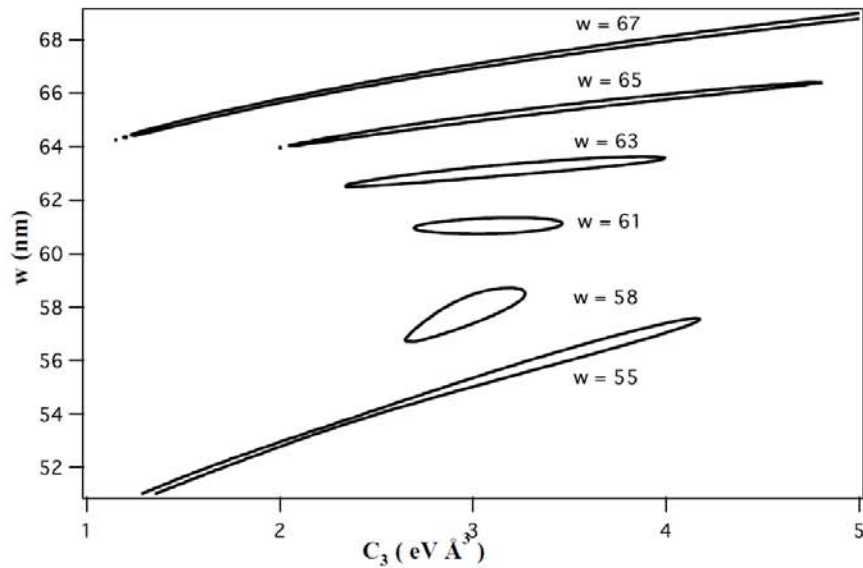


FIG. 4: The figure shows contours of  $\chi^2 = \chi_{min}^2 + 1$  in the  $C_3$ - $w$  plane corresponding to the  $1 - \sigma$  uncertainty of the parameters. Contours are shown for simulated data sets with  $C_3 = 3 \text{ eV \AA}^3$ ,  $d = 100 \text{ nm}$ , and different values of  $w$ . The contours describe extended valleys at most open-fractions but narrow to a well defined minimum for one particular magic open-fraction.

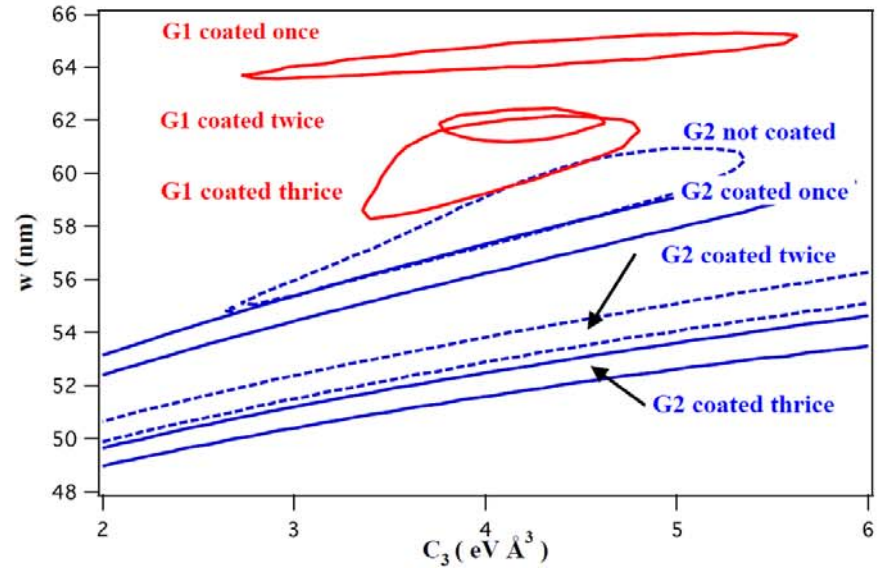


FIG. 5: Contours of  $\chi^2 = \chi_{min}^2 + 1$  in  $C_3$ - $w$  space for two gratings (labeled G1 and G2) after repeated coating with Au/Pd. The small contour for “G1 coated twice” shows the impact of the magic open-fraction.

Lonij, V., W. Holmgren, and A. Cronin,  
 “Magic ratio of window width to grating period for Van der Waals potential  
 measurements using material gratings,” submitted to Phys. Rev. A (2009)



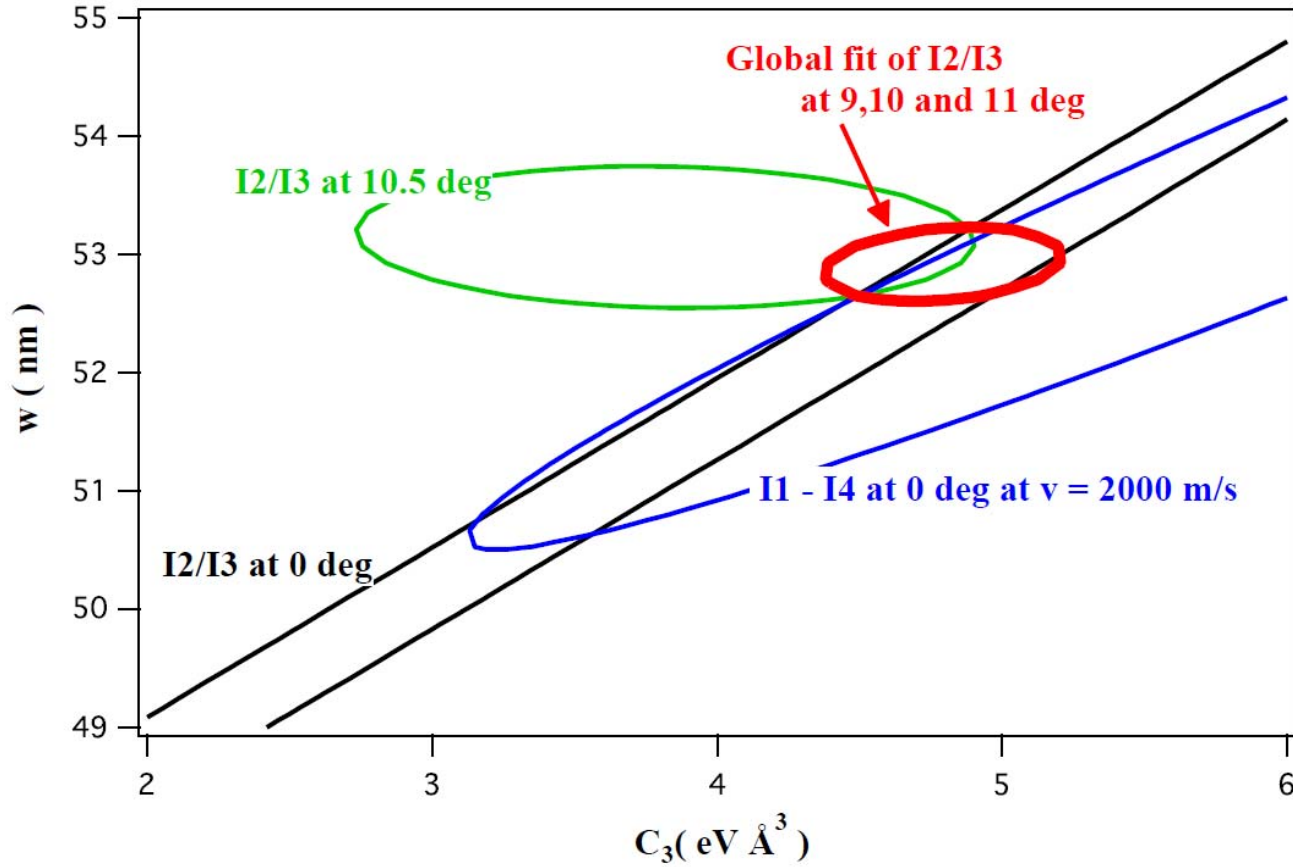


FIG. 8: Contours of  $\chi^2 = \chi_{min}^2 + 1$  in  $C_3$ - $w$  space for different angles. The blue ellipse corresponds to a fit of data like Fig. 2 at 2000 m/s at normal incidence. The black and green ellipses correspond to data similar to Fig. 7 at a single angle  $\beta = 0$  and  $\beta = 10.5$  degrees respectively. The red ellipse corresponds to the fit of the three angles shown in Fig. 7.

TABLE I: Experimental results and theoretical predictions

Experiment	$C_3(eV\text{\AA}^3)$	
This work:		
Na and SiN <sub>x</sub> fit parameter	3.26	$\pm 0.16^a$
Na and Au/Pd fit parameter G1	4.30	$\pm 0.5$
Na and Au/Pd fit parameter IG	4.80	$\pm 0.5$
Na and SiN <sub>x</sub> with PWI correction	3.42	$\pm 0.19^a$
Na and Au/Pd with PWI correction G1	4.51	$\pm 0.5$
Na and Au/Pd with PWI correction IG	5.04	$\pm 0.5$
Previous work:		
Na and SiN <sub>x</sub> [3]	2.70	$\pm 0.8$
Theory		
Na and perfect conductor	7.6	[13]
Na and SiN <sub>x</sub> using single oscillator $\alpha_{\text{pol}}(i\omega)$ and tabulated $\epsilon(i\omega)$	3.3	[12], eqn. (5)
Na and SiN <sub>x</sub> using tabulated $\alpha_{\text{pol}}(i\omega)$ and $\epsilon(i\omega)$	3.48	[12, 15], eqn. (5)
Na and Bulk Au	5.11	[16]
Na and 1 nm Au at $r = 10$ nm	4.3	[17]
Na and 2 nm Au at $r = 10$ nm	4.5	[17]
Na and 3 nm Au at $r = 10$ nm	4.6	[17]

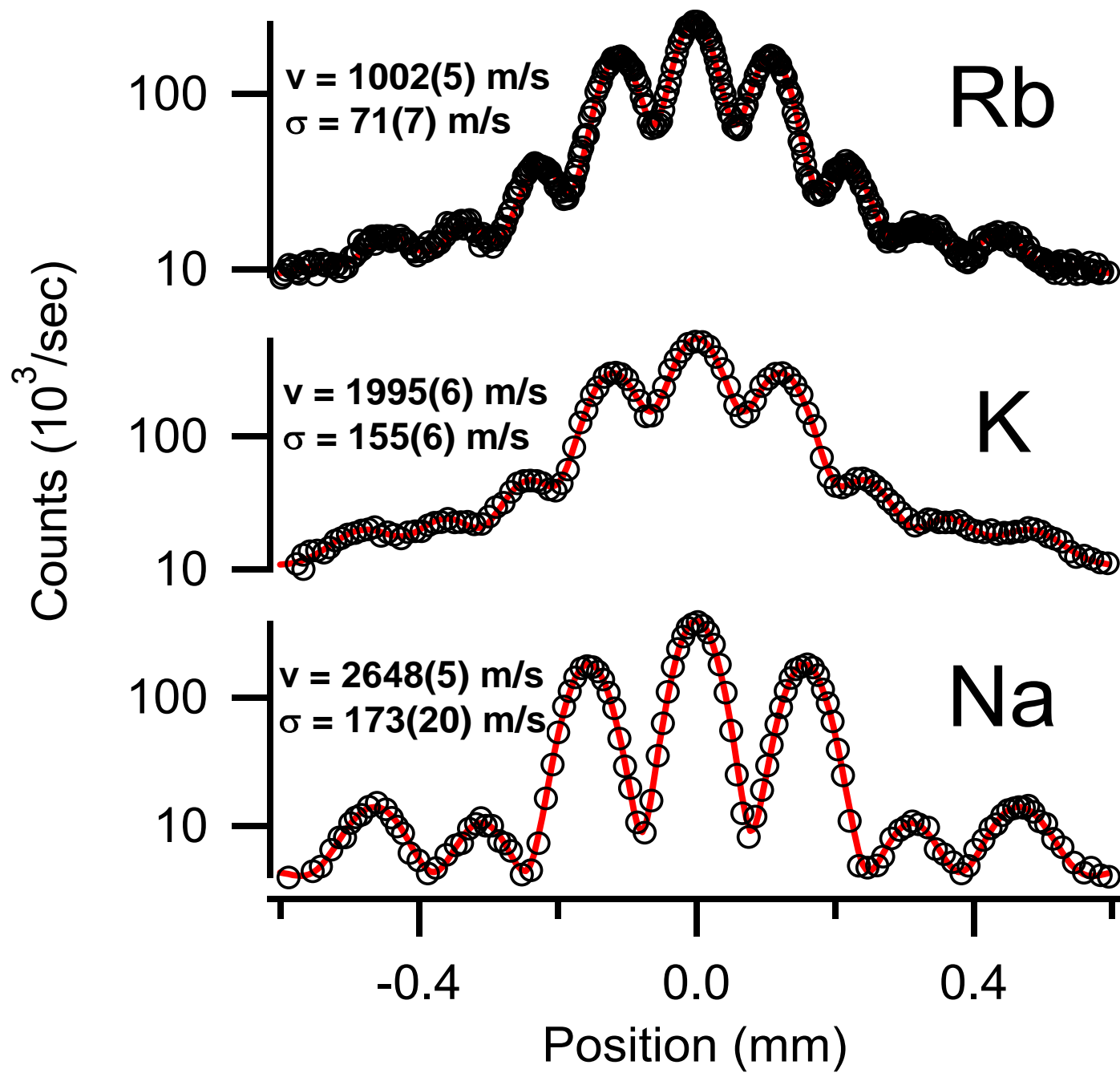
<sup>a</sup>The source of the reported error is discussed in appendix B.

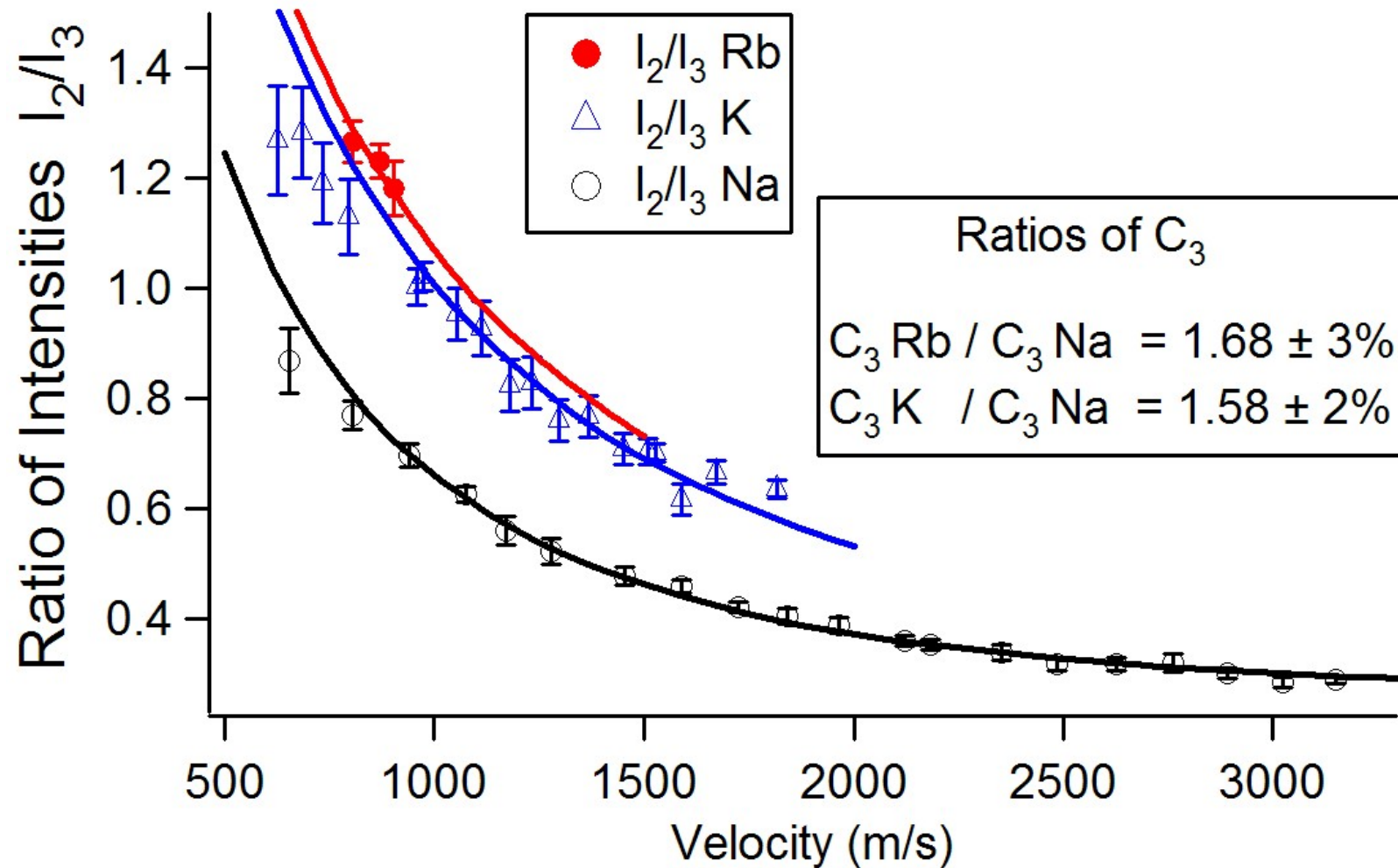
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Na and SiN <sub>x</sub> using tabulated $\alpha_{\text{pol}}(i\omega)$ and $\epsilon(i\omega)$	3.48	[12, 15], eqn. (5)
Na and Bulk Au	5.11	[16]
Na and 1 nm Au at $r = 10$ nm	4.3	[17]
Na and 2 nm Au at $r = 10$ nm	4.5	[17]
Na and 3 nm Au at $r = 10$ nm	4.6	[17]

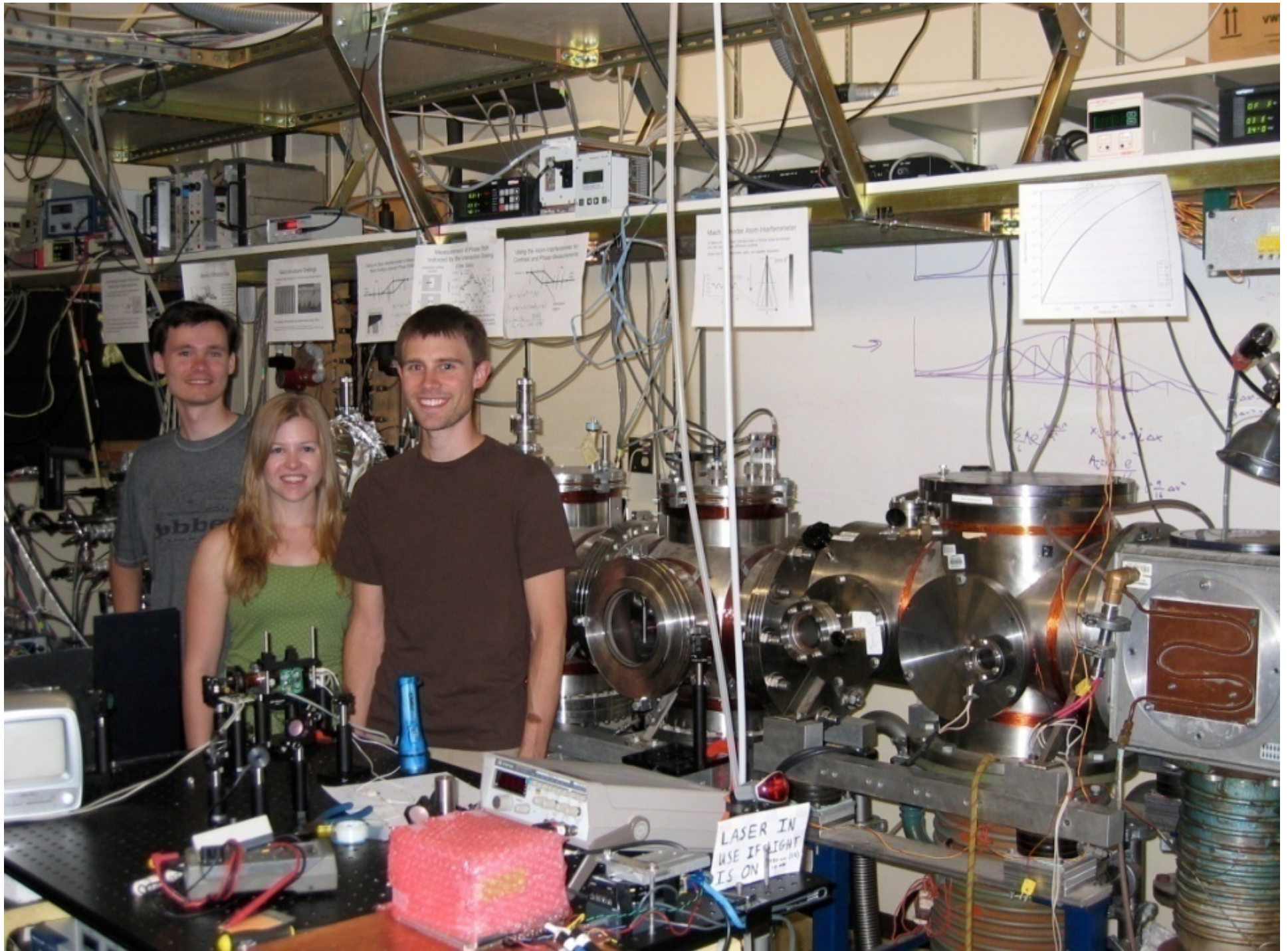
[12] H. R. Philipp. Journal of the Electrochemical Society  
**120**, 295 (1973).

[15] A. Derevianko, S. G. Porsev, and J. F. Babb,  
arXiv:0902.3929v1 (2009).



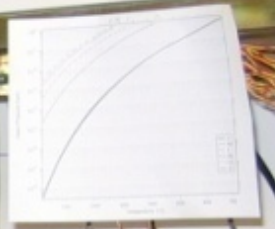


C <sub>3</sub>	Li – SiNx:	3.25 ± 0.20	(6 %)
C <sub>3</sub>	Na – SiNx:	3.42 ± 0.19	(5 %)
C <sub>3</sub>	Na – AuPd:	5.04 ± 0.50	(10 %)
C <sub>3</sub>	K – SiNx:	5.40 ± 0.39	(8 %)
C <sub>3</sub>	Rb – SiNx:	5.74 ± 0.40	(8 %)



Measurement of Phase Shift Induced by the Laser Light  
Using the Magneto-optical Effect  
Magneto-optical Effect

Magneto-optical Effect



Handwritten notes on a whiteboard, including the equation  $\epsilon_{xy} = \frac{1}{2} \chi \sin \theta$  and other mathematical expressions.

LASER IN USE IF LIGHT IS ON

# Outline

- 1. New polarizability measurements**
- 2. Atom Interferometry C3 measurements**  
**tests of potential power law**  
**limits on non-Newtonian gravity**
- 3. More precise diffraction experiments**  
**Magic open fraction**  
**Effect of AuPd**  
**Ratios of C3 for Rb, K, Na**
- 4. Laser modified vdW interactions**



PHYSICAL REVIEW A 77, 043406 (2008)

## Modifying atom-surface interactions with optical fields

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M. Bhattacharya, Vincent P. A. Lonij, and Alexander D. Cronin

*Department of Physics, University of Arizona, Tucson, Arizona 85721, USA*

- Intensity of  $5 \text{ W/cm}^2$  will double the vdW attractive force with  $\Delta = 100 \Gamma$ .
- Light polarization will change the enhancement by a factor of 2
- Enhancement depends on laser detuning squared  $(\omega - \omega_0)^2$
- Dielectric and conducting surfaces have the same laser-induced enhancement factor.

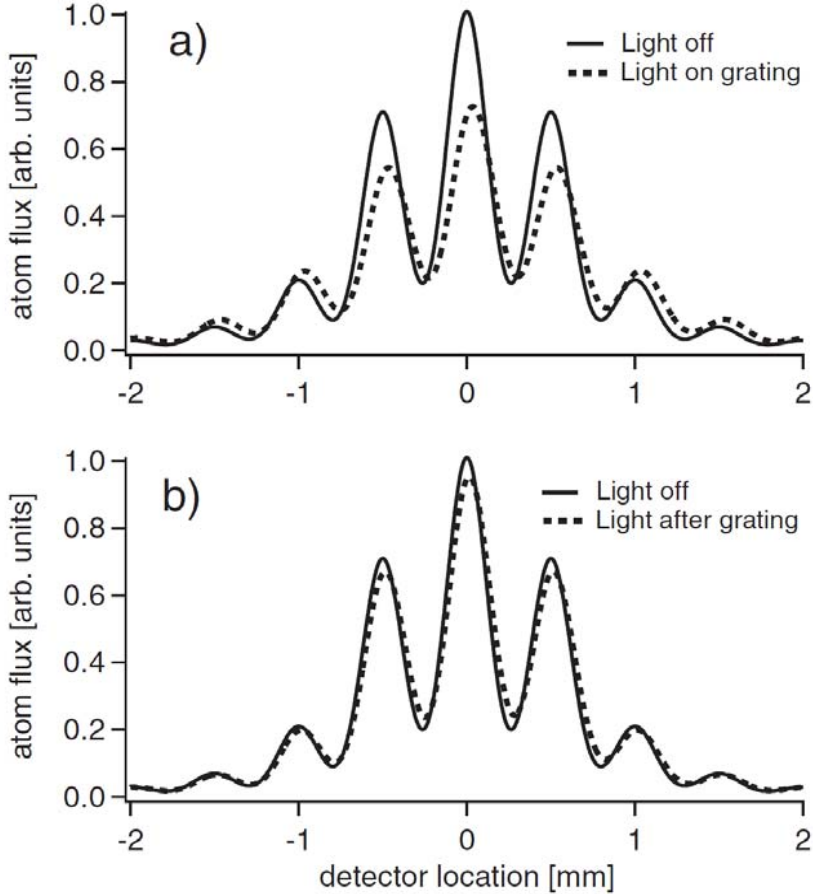
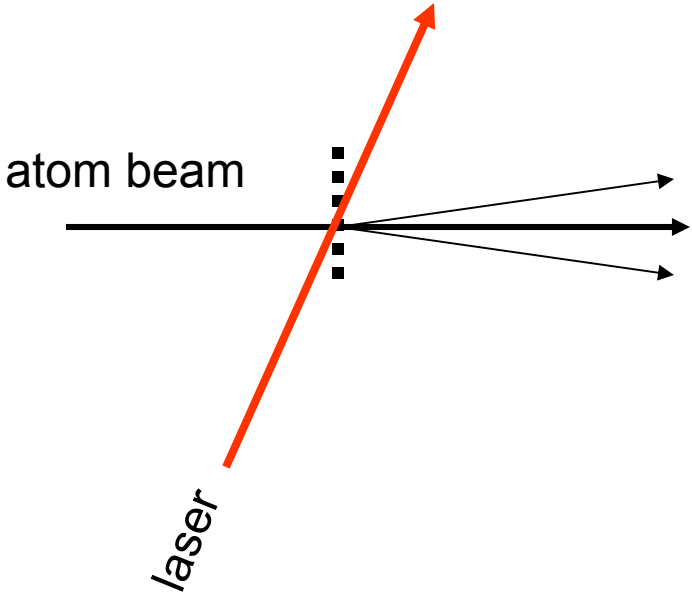


FIG. 3. Simulated laser-modified atom diffraction patterns. Light is predicted to change the diffraction patterns because of modified atom-surface interactions and also photon scattering. Plot (a) shows the effect of a laser beam, with intensity  $I=5 \text{ W/cm}^2$  and detuning  $\Delta=100 \text{ MHz}$ , illuminating the atoms inside a nanograting. Plot (b) shows the effect of the same laser illuminating atoms after passing through a nanograting. In each simulation atoms exposed to the laser have an equal probability of spontaneously scattering 1 or 0 photons.

# Thank You

- Vincent Lonij, Will Holmgren, Melissa Revelle
- NSF, Research Corporation, U. of Arizona

Cronin, A.D., Schmiedmayer J., and Pritchard D.E., “Optics and Interferometry with atoms and molecules.” *Rev. Mod. Phys* **81**, 1051 (2009)

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Lonij, V., W. Holmgren, and A. Cronin, “Magic ratio of window width to grating period for Van der Waals potential measurements using material gratings,” submitted to *Phys. Rev. A* (2009)

Holmgren, W., Lonij V., Revelle M., and Cronin AD., “Measurements of Na, K, and Rb polarizabilities with an atom interferometer,” in prep. for *Phys. Rev. A*.