

Energy Cascades in Granular Matter

Eli Ben-Naim

Theory Division

Los Alamos National Laboratory

with: Jon Machta (Massachusetts)

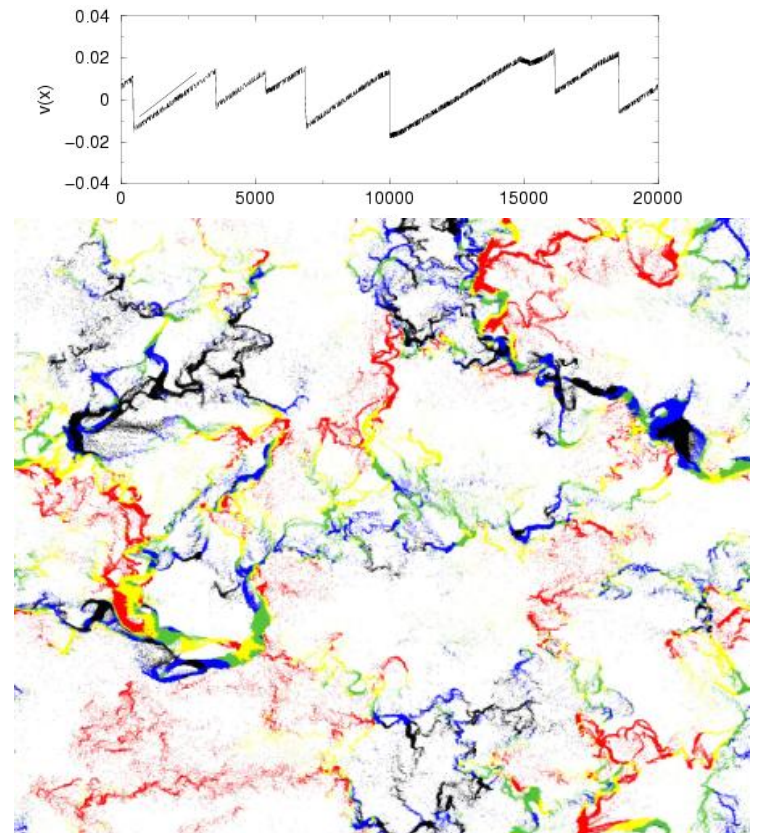
cond-mat/0411743

<http://cnls.lanl.gov/~ebn>



Energy dissipation in granular media

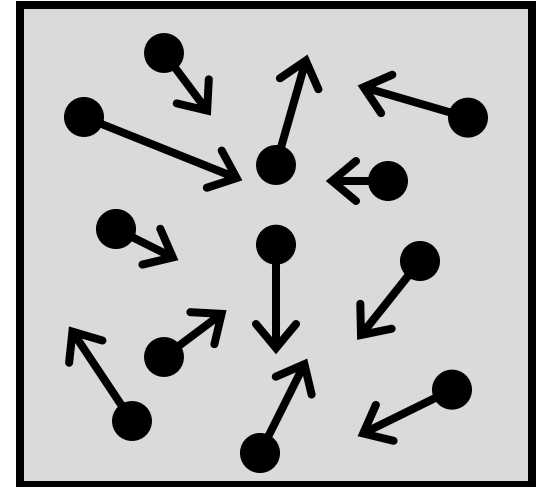
- ◆ Responsible for collective phenomena
 - » Clustering
 - » Hydrodynamic instabilities
 - » Shocks
 - » Pattern formation
- ◆ Anomalous statistical mechanics:
 - No energy equipartition
 - Nonequilibrium distributions



$$P(E) \neq \exp(-E / kT)$$

Inelastic gas

- ◆ Vigorous driving
- ◆ Spatially uniform system
- ◆ Particles undergo binary collisions
- ◆ Velocity changes due to
 1. Inelastic collisions (lose energy)
 2. Energy input (gain energy)



- ◆ What is the typical velocity (granular “temperature”)?

$$T = \langle v^2 \rangle$$

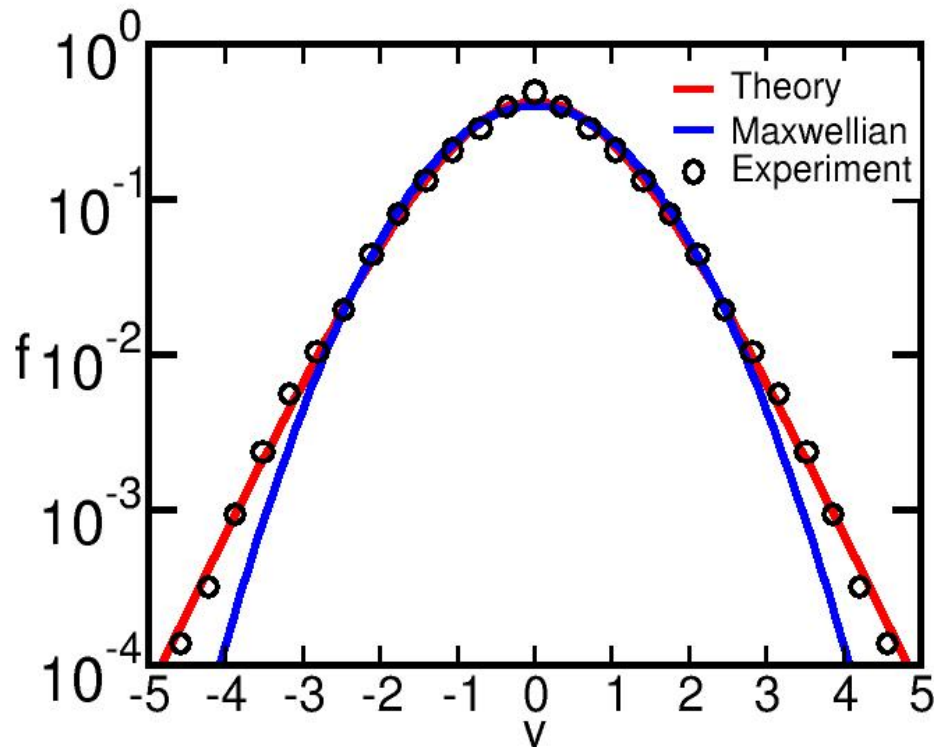
- ◆ What is the velocity distribution?

$$f(v)$$

Nonequilibrium velocity distributions

$$f(v) \sim \exp(-|v|^\delta) \quad 1 \leq \delta \leq 3/2$$

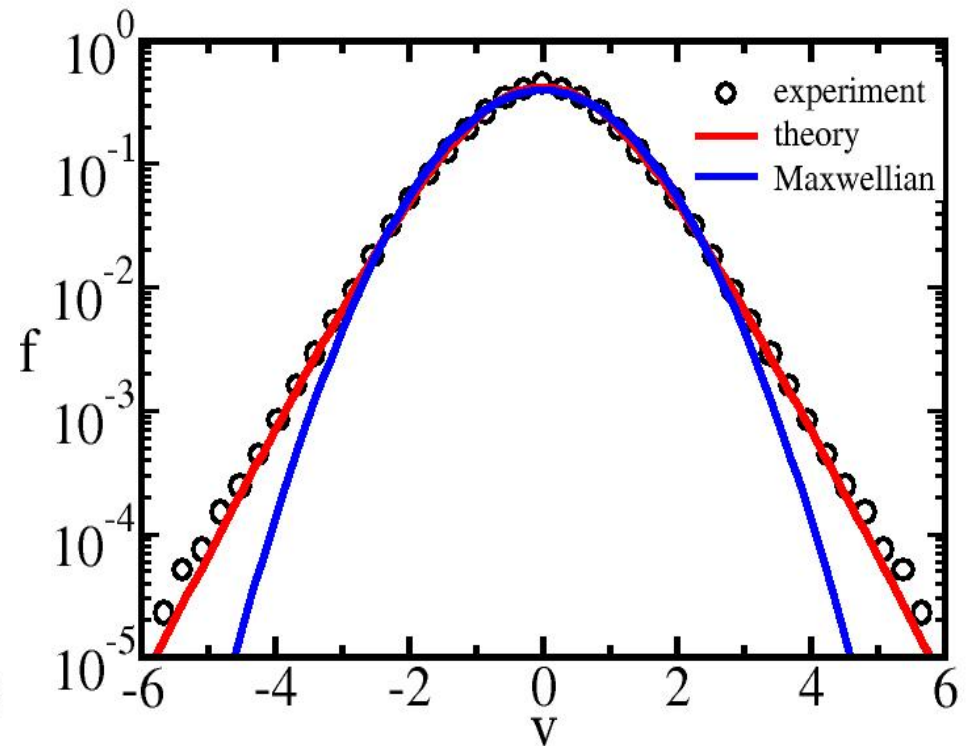
$$\langle v^4 \rangle / \langle v^2 \rangle^2 = 117/33 \cong 3.55$$



Mechanically vibrated beads

Rouyer & Menon 2000

Theory: ebn & krapivsky 2002



Electrostatically driven powders

Aronson & Olafsen 2002

Inelastic Collisions

- ◆ Relative velocity reduced by $0 < r < 1$

$$v_1 - v_2 = -r(u_1 - u_2)$$

- ◆ Momentum is conserved

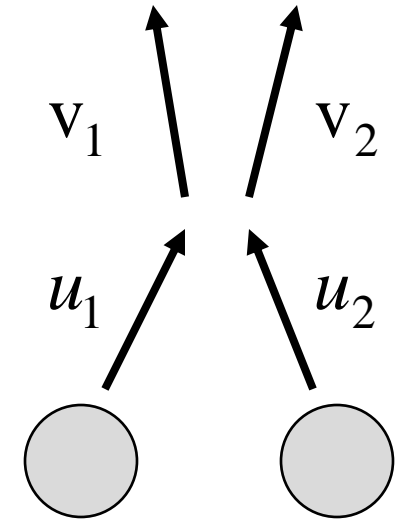
$$v_1 + v_2 = u_1 + u_2$$

- ◆ Energy is dissipated

$$\Delta E \propto -(1 - r)(\Delta v)^2$$

- ◆ Limiting cases

$$r = \begin{cases} 0 & \text{completely inelastic } (\Delta E = \max) \\ 1 & \text{elastic } (\Delta E = 0) \end{cases}$$



Freely decaying states

- ◆ Energy loss in a collision $\Delta T = \langle v^2 \rangle$
- ◆ Collision rate $\Delta t \sim 1/(\Delta v)^\lambda$
- ◆ Energy balance equation

$$\frac{\Delta T}{\Delta t} \sim -(\Delta v)^{2+\lambda} \quad \Rightarrow \quad \frac{dT}{dt} \sim -T^{1+\lambda/2}$$

- ◆ Temperature decays, system comes to rest

$$T \sim t^{-2/\lambda} \quad \Rightarrow \quad P(v) \rightarrow \delta(v)$$

Trivial steady-state

Kinetic Theory

◆ Collision rule (linear)

$$(u_1, u_2) \rightarrow (pu_1 + qu_2, qu_1 + pu_2) \quad r = 1 - 2p$$

◆ Boltzmann equation

$$\frac{\partial P(v)}{\partial t} = \iint du_1 du_2 P(u_1) P(u_2) |u_1 - u_2|^\lambda \left[\underbrace{\delta(v - pu_1 - qu_2)}_{\text{collision rate}} - \underbrace{\delta(v - u_2)}_{\text{gain}} \right]_{\text{loss}}$$

◆ Collision rate related to interaction potential

$$U(r) \sim r^{-\gamma} \quad \lambda = 1 - 2 \frac{d-1}{\gamma} = \begin{cases} 0 & \text{Maxwell molecules } (\alpha = 2, 2D) \\ 1 & \text{Hard spheres } (\alpha = \infty) \end{cases}$$

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collision rate gain loss

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Are there nontrivial steady states?

An exact solution

- ◆ One-dimensional Maxwell molecules
- ◆ Fourier transform obeys a closed equation

$$F(k) = F(pk)F(k - pk) \quad F(k) = \int dv e^{ikv} f(v)$$

- ◆ Exponential solution

$$F(k) = \exp(-|k|v_0)$$

- ◆ Lorentzian velocity distribution

$$f(v) = \frac{1}{\pi v_0} \frac{1}{1 + (v/v_0)^2}$$

Nontrivial steady states do exist

Properties of stationary state

- ◆ Perfect balance between collisional loss and gain
- ◆ Power-law high-energy tail

$$P(v) \sim v^{-\sigma} \quad \sigma = 2$$

- ◆ Infinite energy, infinite dissipation!

Is this stationary state physical?

Cascade Dynamics (1D)

- ◆ Collision rule: arbitrary velocities

$$(u_1, u_2) \rightarrow (pu_1 + qu_2, pu_2 + qu_1)$$



- ◆ Large velocities cascade

$$v \rightarrow (pv, qv)$$

- ◆ High-energies: linearized equation

$$\frac{1}{p^{1+\lambda}} f\left(\frac{v}{p}\right) + \frac{1}{q^{1+\lambda}} f\left(\frac{v}{q}\right) - f(v) = 0$$

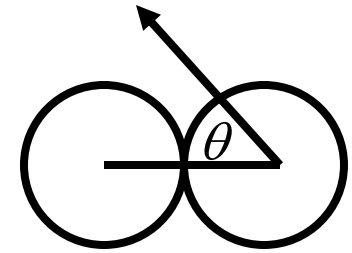
- ◆ Power-law tail

$$f(v) \sim v^{-\sigma} \quad \sigma = 2 + \lambda$$

Cascade Dynamics

- ◆ Collision process: large velocities

$$v \rightarrow (\alpha v, \beta v)$$



- ◆ Stretching parameters related to impact angle

$$\alpha = (1 - p) \cos \theta \quad \beta = [1 - (1 - p^2) \cos^2 \theta]^{1/2}$$

- ◆ Energy decreases, velocity magnitude increases

$$\alpha^2 + \beta^2 \leq 1 \quad \alpha + \beta \geq 1$$

- ◆ Steady state equation

$$\left\langle \left(\frac{1}{\alpha^{d+\lambda}} f\left(\frac{v}{\alpha}\right) + \frac{1}{\beta^{d+\lambda}} f\left(\frac{v}{\beta}\right) - f(v) \right) \cos^{\lambda} \theta \right\rangle = 0$$

Power-laws are generic

- ◆ Velocity distributions always has power-law tail

$$f(v) \sim v^{-\sigma}$$

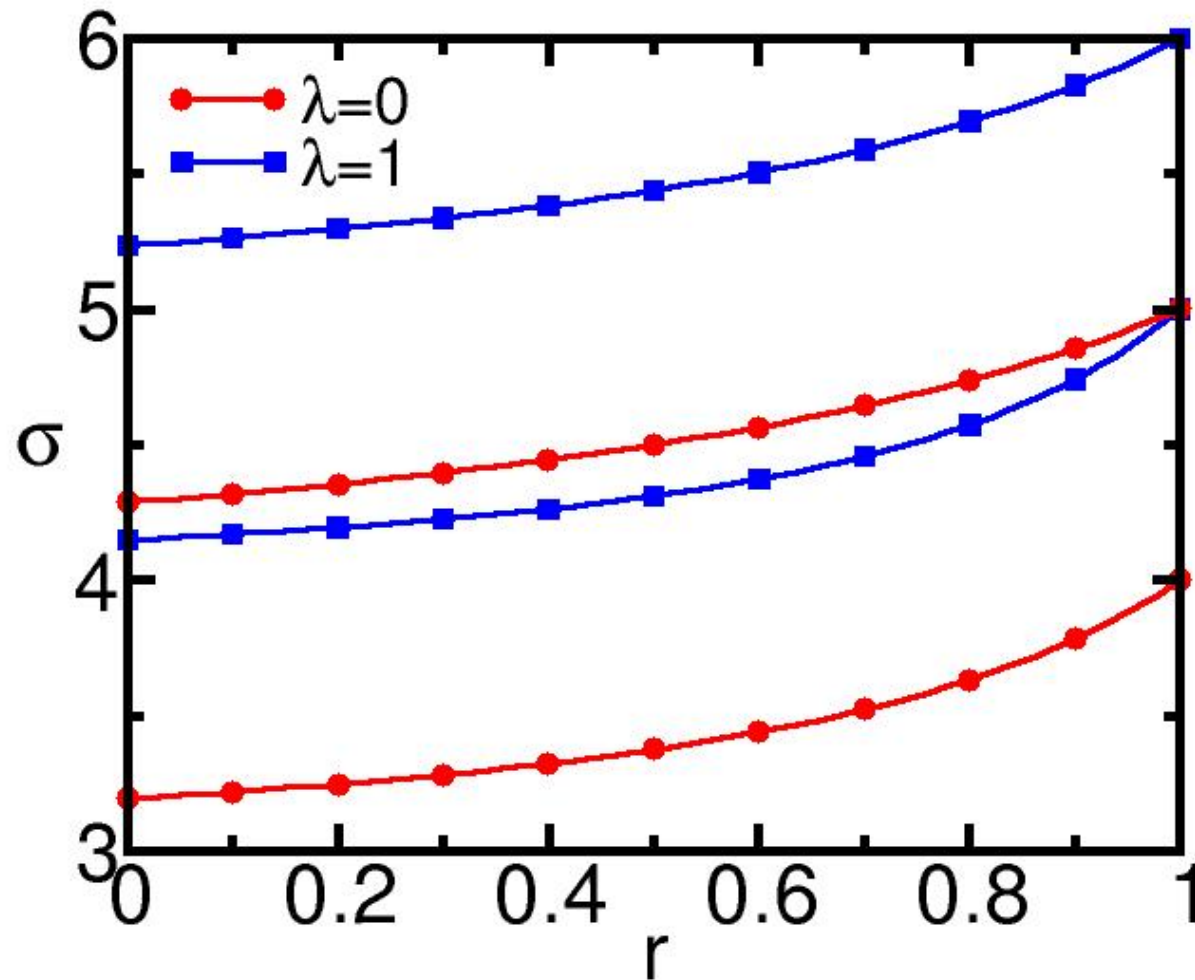
- ◆ Exponent varies with parameters

$$\frac{{}_1F_2\left(\frac{d + \lambda - \sigma}{2}, \frac{\lambda + 1}{2}, \frac{d + \lambda}{2}, 1 - p^2\right)}{(1 - p)^{\sigma - d - \lambda}} = \frac{\Gamma\left(\frac{\sigma - d + 1}{2}\right)\Gamma\left(\frac{d + \lambda}{2}\right)}{\Gamma\left(\frac{\sigma}{2}\right)\Gamma\left(\frac{\lambda + 1}{2}\right)}$$

- ◆ Tight bounds $1 \leq \sigma - d - \lambda \leq 2$
- ◆ Elastic limit is singular $\sigma \rightarrow d + 2 + \lambda$

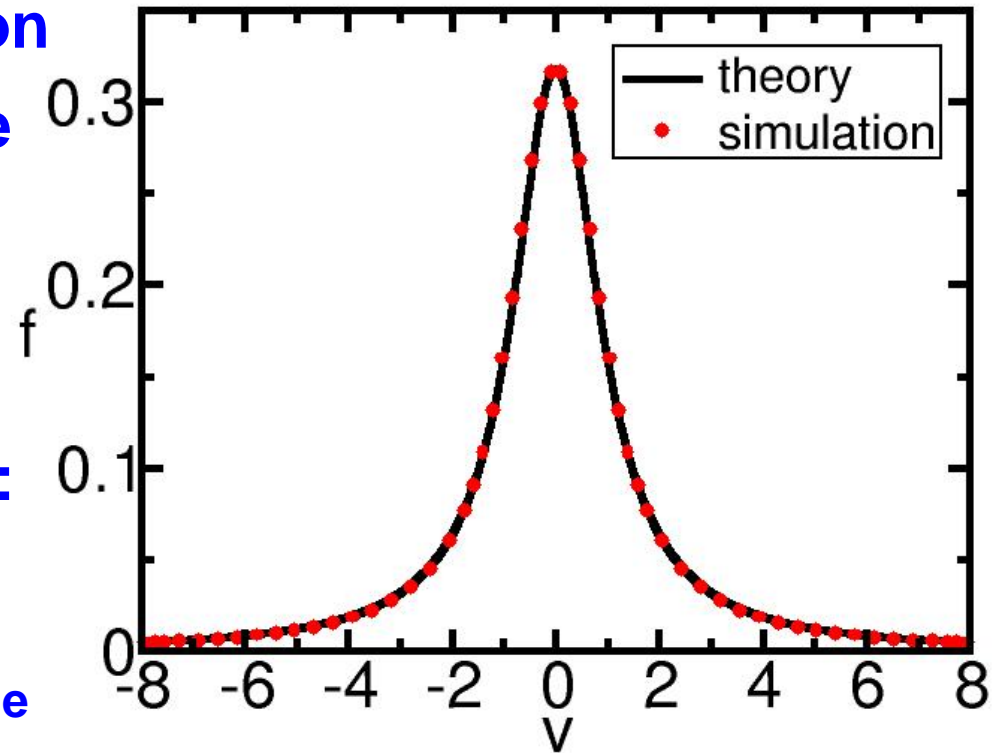
**Dissipation always divergent
Energy finite or infinite**

The Characteristic Exponent



Monte Carlo Simulations

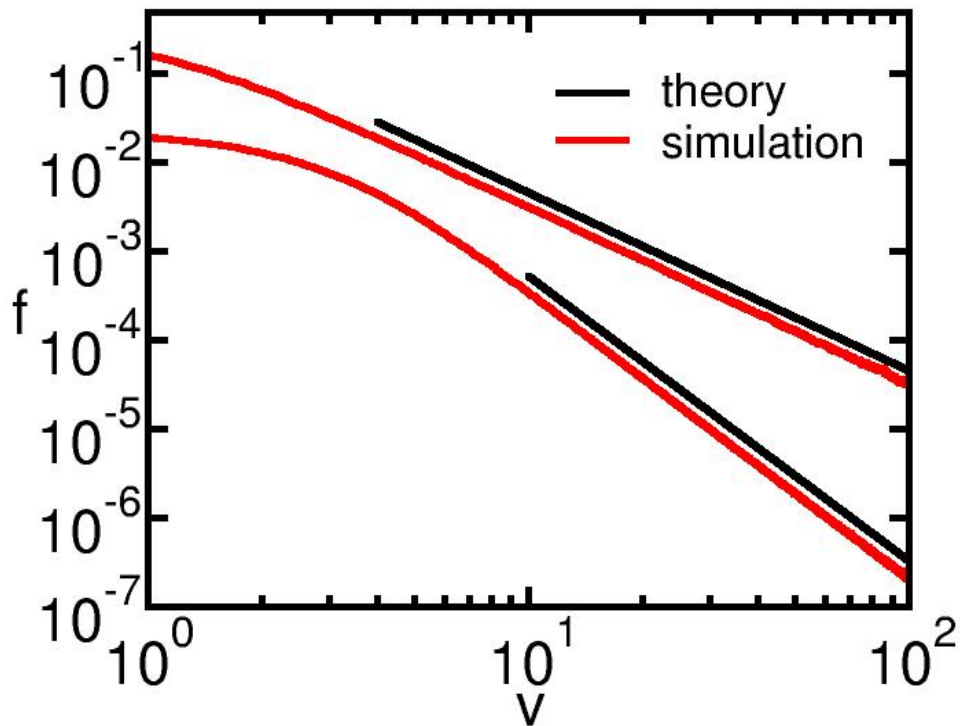
- ◆ Compact initial distribution
- ◆ Inject energy at very large velocity scales only
- ◆ Maintain constant total energy
- ◆ “Lottery” implementation:
 - Keep track of total energy dissipated, E_T
 - With small rate, boost a particle by E_T



Excellent agreement between theory and simulation

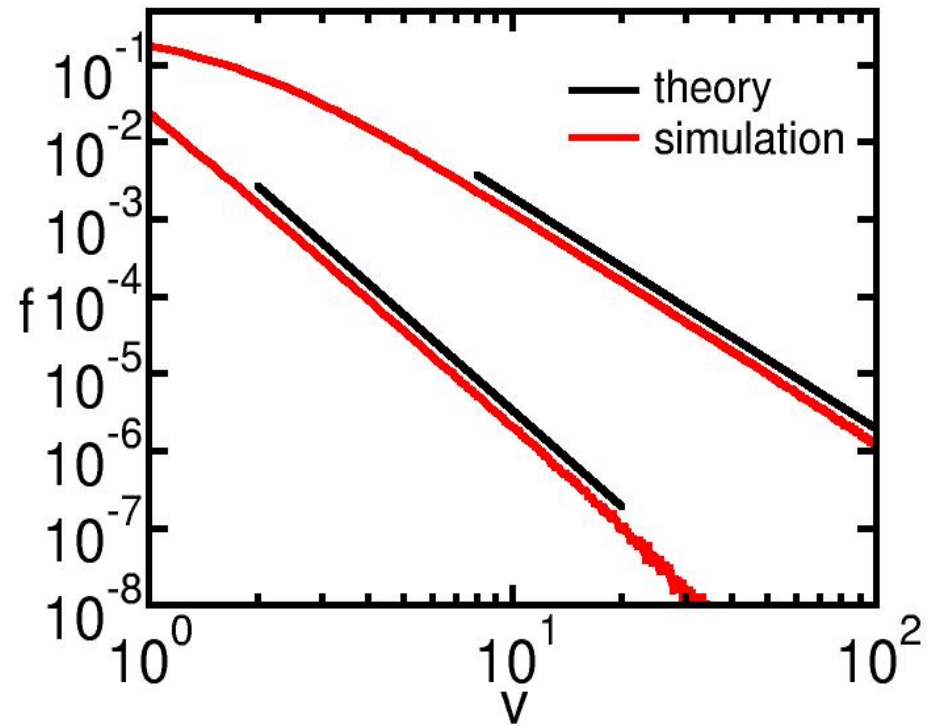
Further confirmation

Maxwell molecules (1D, 2D)



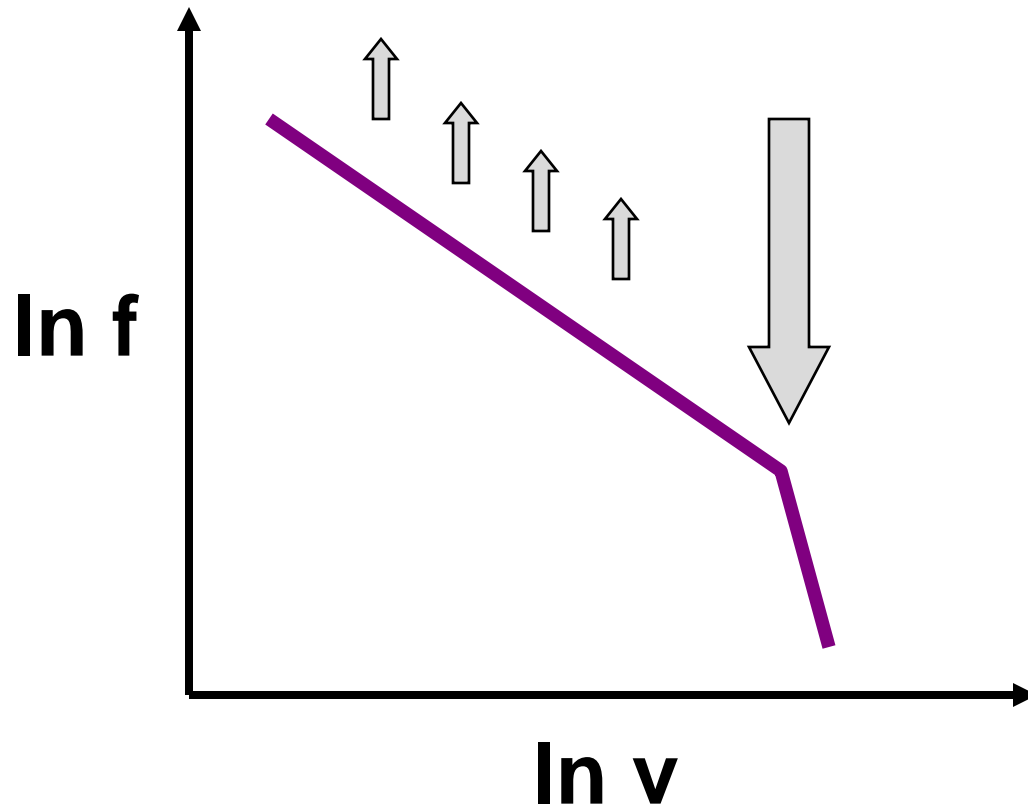
$N=10^7$

Hard spheres (1D, 2D)



$N=10^5$

Injection, cascade, dissipation



- ❖ Energy is injected at large velocity scales
- ❖ Energy cascades from large velocities to small velocities
- ❖ Energy dissipated at small velocity scales

Conclusions

- ◆ **New class of nonequilibrium stationary states**
- ◆ **Energy cascades from large velocities to small velocities**
- ◆ **Power-law high-energy tail**
- ◆ **Energy input at large scales balances dissipation**
- ◆ **Temperature insufficient to characterize velocities**
- ◆ **Experimental realization: requires a different driving mechanism**